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# Control of Separated Flows

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UNIVERSITY OF  
**TORONTO**

Summer School on Sustainable Aviation  
Active Flow Control  
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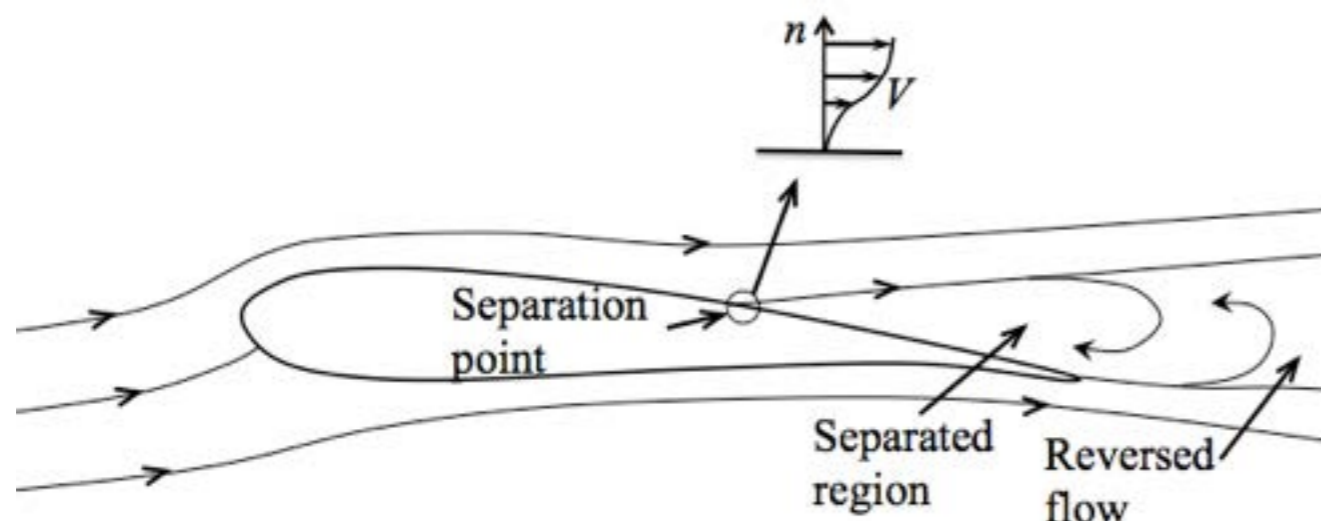
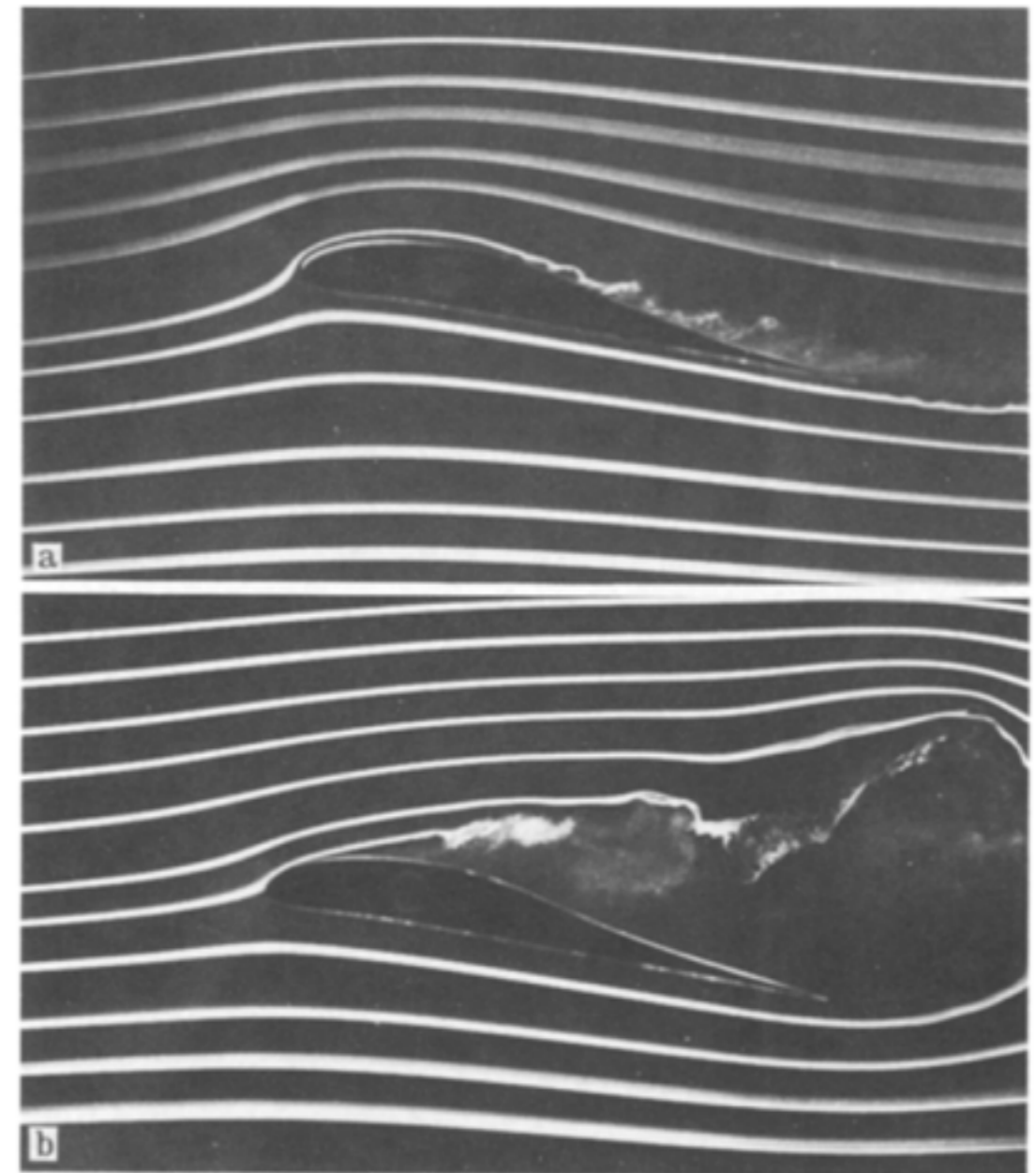
# Outline



- A. What is a separated flow and why does it separate?
- B. Examples of separated flows on an airplane, and what it does to performance
- C. Control of separation on streamlined bodies
- D. Flow State Estimation - Enabler for Closed-loop Control
- E. Summary and Closing Remarks

# Flow Separation

- The streamlines of a flow tends to follow surfaces, due to the Coanda effect.
- If the curvature is “too large”, the flow will “separate” from the surface, creating a separated shear layer.
- More precisely, the adverse pressure gradient, against which the fluid must work, produces a local flow reversal and thus separation of the shear layer.



Mueller et al. (1983)

# Physics Behind Flow Separation



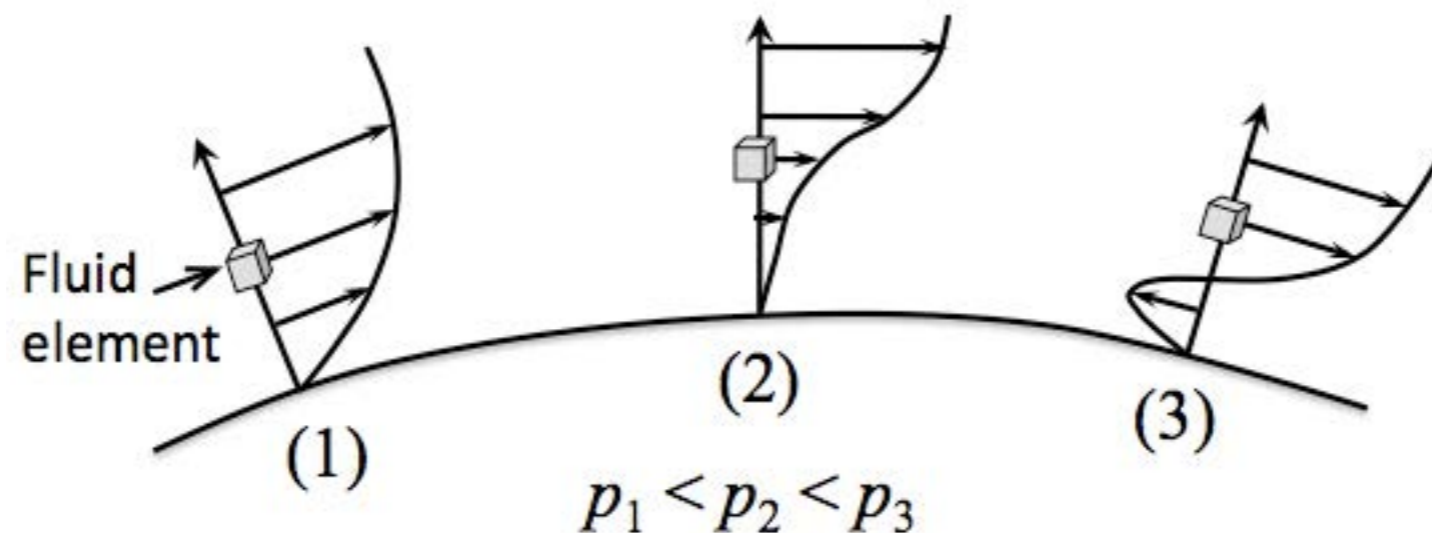
Consider the x-momentum equation for a boundary layer:

$$\rho u \frac{\partial u}{\partial x} + \rho v \frac{\partial u}{\partial y} = -\frac{dp_e}{dx} + \frac{\partial}{\partial y} \left( \mu \frac{\partial u}{\partial y} \right)$$

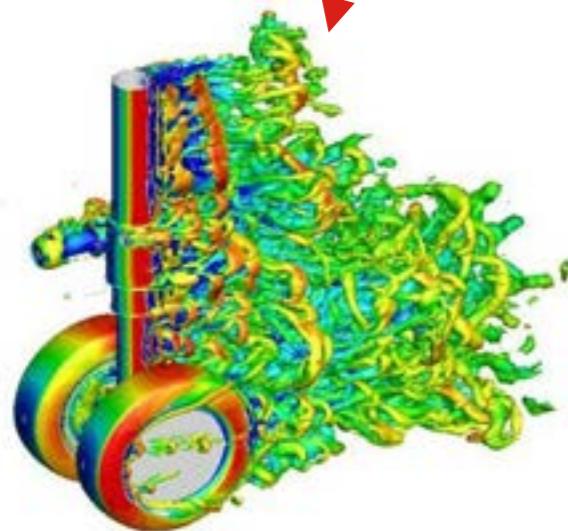
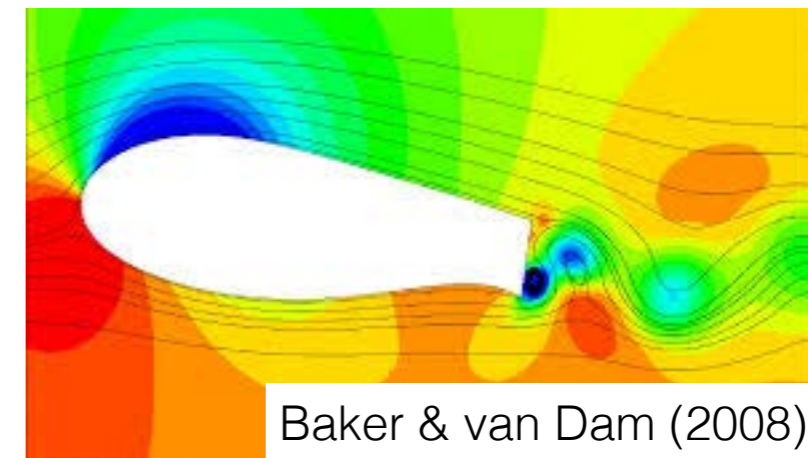
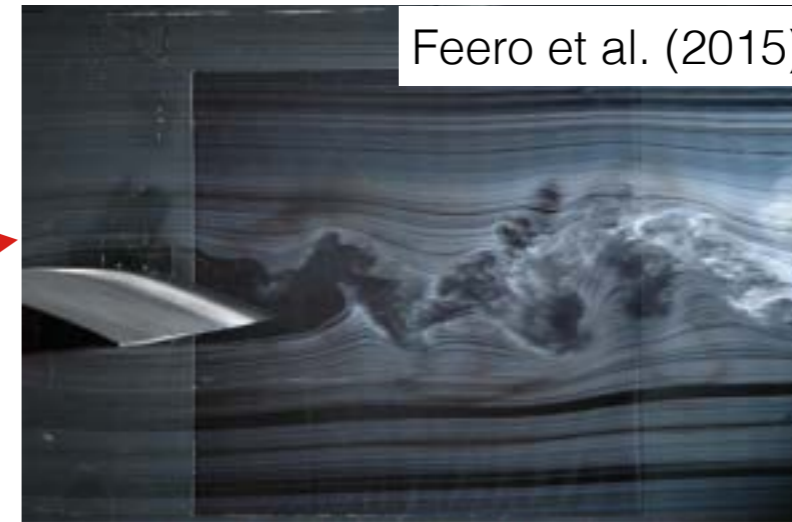
Assuming constant viscosity, the equation at the wall becomes,

$$\frac{dp_e}{dx} = \mu \frac{\partial^2 u}{\partial y^2}$$

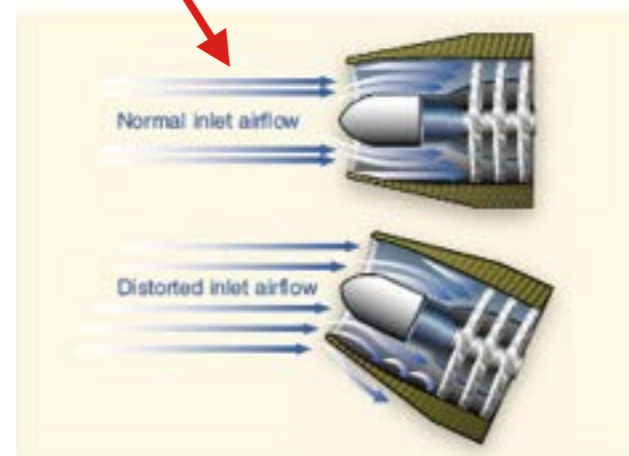
Thus, a positive (adverse) pressure gradient wants to induce a positive curvature to the velocity profile.



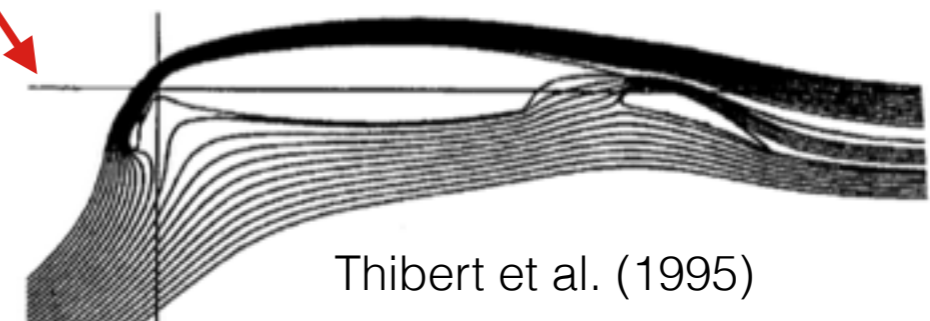
# Flow Separation in Aviation



[www.southampton.ac.uk](http://www.southampton.ac.uk)



[www.flightlearnings.com](http://www.flightlearnings.com)



Thibert et al. (1995)

# Why Control Separation?



If an aircraft is well designed, **flow separation should not be an issue**, except for extreme situations (e.g., one-engine out, large wind gust) or when loss of lift and increase drag is needed (e.g., spoilers, landing gear).

## Why should aviation care then?

A. Many boundaries of the **flight envelop** set by flow separation.

B. Separation control as an **enabler technology**:

- Decrease the size of vertical stabilizer (reduces weight and drag)
- Gapless high-lift (less weight and noise)
- Decrease number of stages in the compressor (less weight)
- Reduce unsteady aerodynamic loads on landing gear and wing (less weight and noise - maybe drag also)

# Technology Enabler (Example 1)



Take-off and landing speed controlled by  $C_{Lmax}$ , which is limited by flow separation

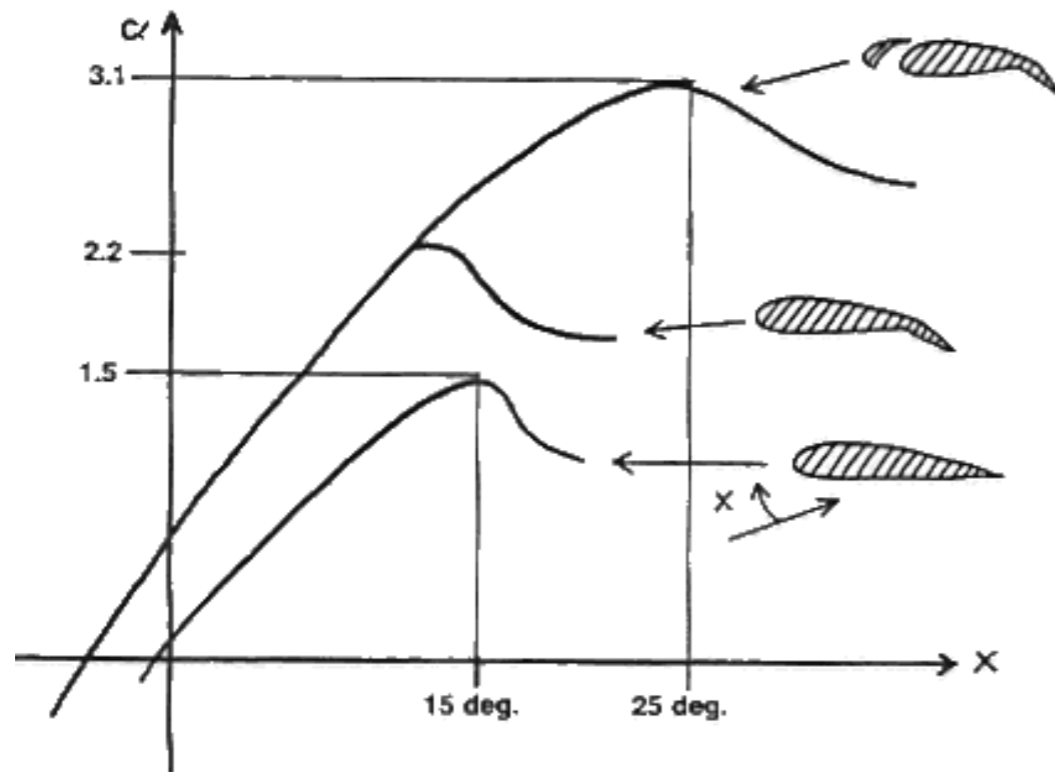


Figure 1



Slotted Fowler flaps of a Boeing 747.

Gap between flap and wing bleeds air from pressure side to energize the boundary layer.

- Very heavy and complex systems

Could separation control replace that heavy system?

# Control on Streamlined Bodies



A key element of separation on streamlined bodies is that the point of separation can move or change based on conditions.

Fundamentally, since separation is caused by the loss of energy near the boundary due to an adverse pressure gradient, there are two possible strategies:

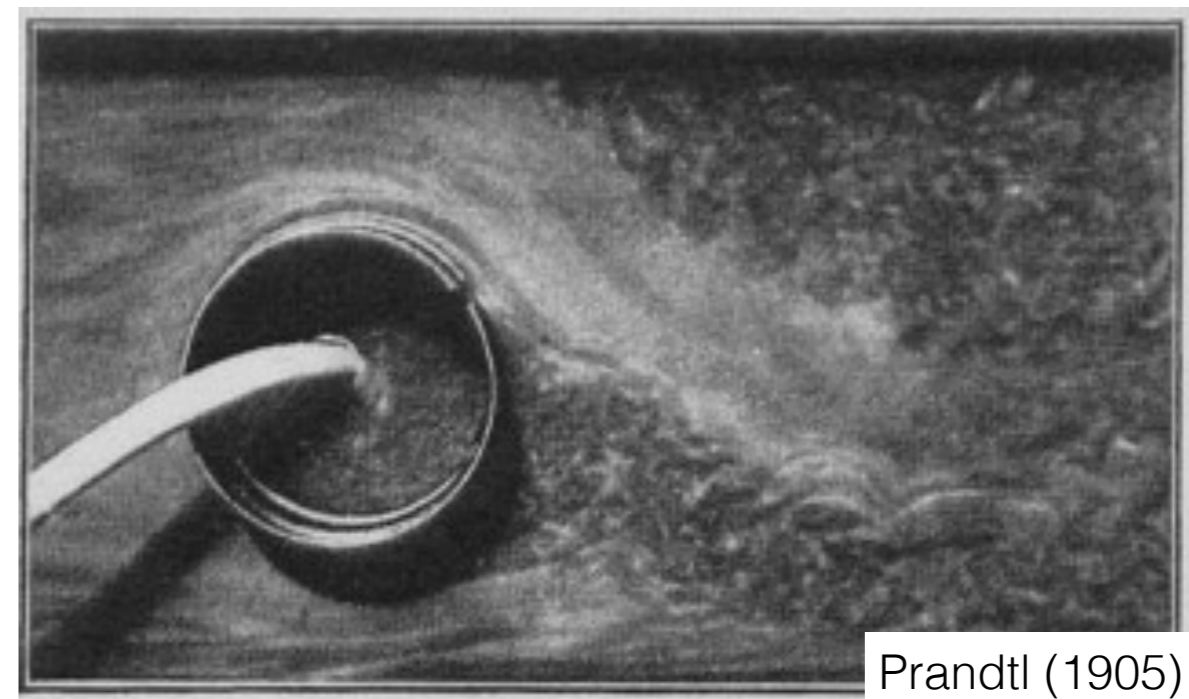
1. Re-energise the boundary layer
  - i. Laminar separation: induce transition to turbulence
  - ii. Turbulent separation: increase turbulence and mixing
2. Virtual aerodynamic shaping



# First Example of Separation Control



Prandtl, in developing his boundary layer theory, devised some of the first systematic flow control experiments using constant suction.



By **removing low momentum fluid** in the boundary layer, separation point moved further back.

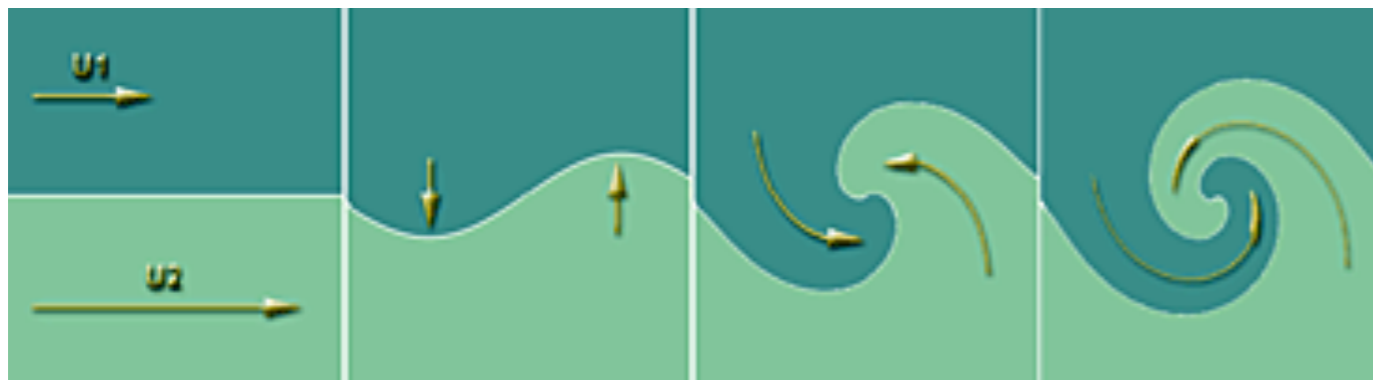
This method not usually used due to **weight** and **energy** requirements.

- Some exceptions, primarily in military applications

# A Smarter Way?

A key element of modern active flow control methodology is to use unsteady forcing, typically exciting an existing instability in the flow.

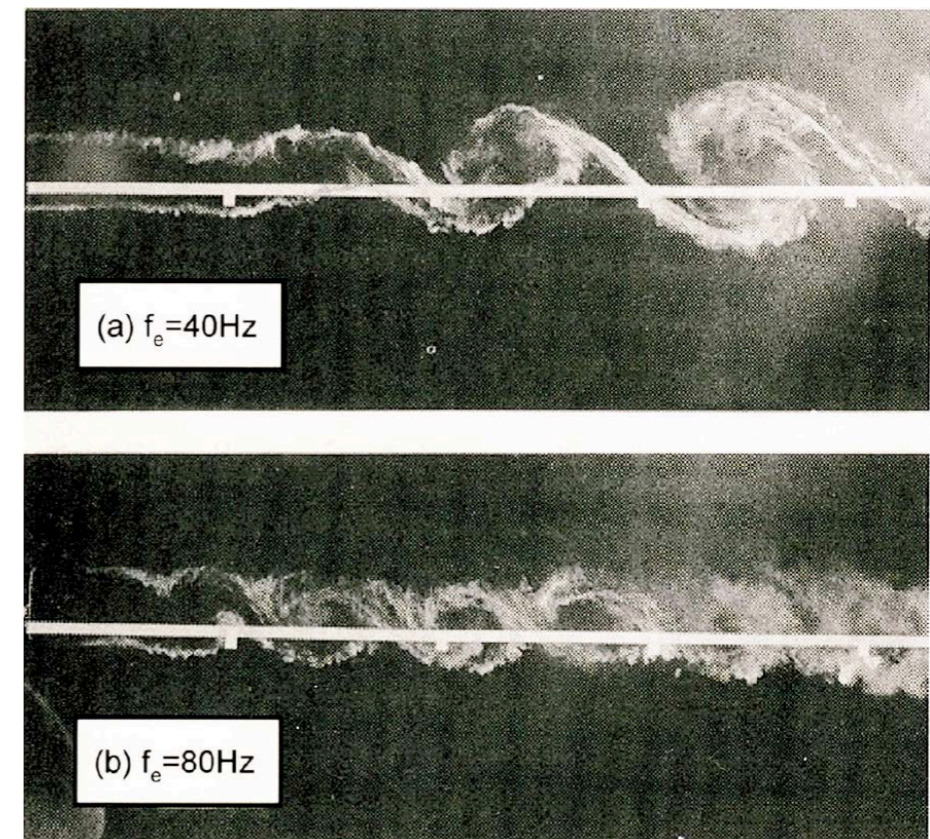
For instance consider a simple free-shear-layer:



Kelvin-Helmholtz instability leads to a roll-up of the shear layer.

Reduced (normalised) frequency:

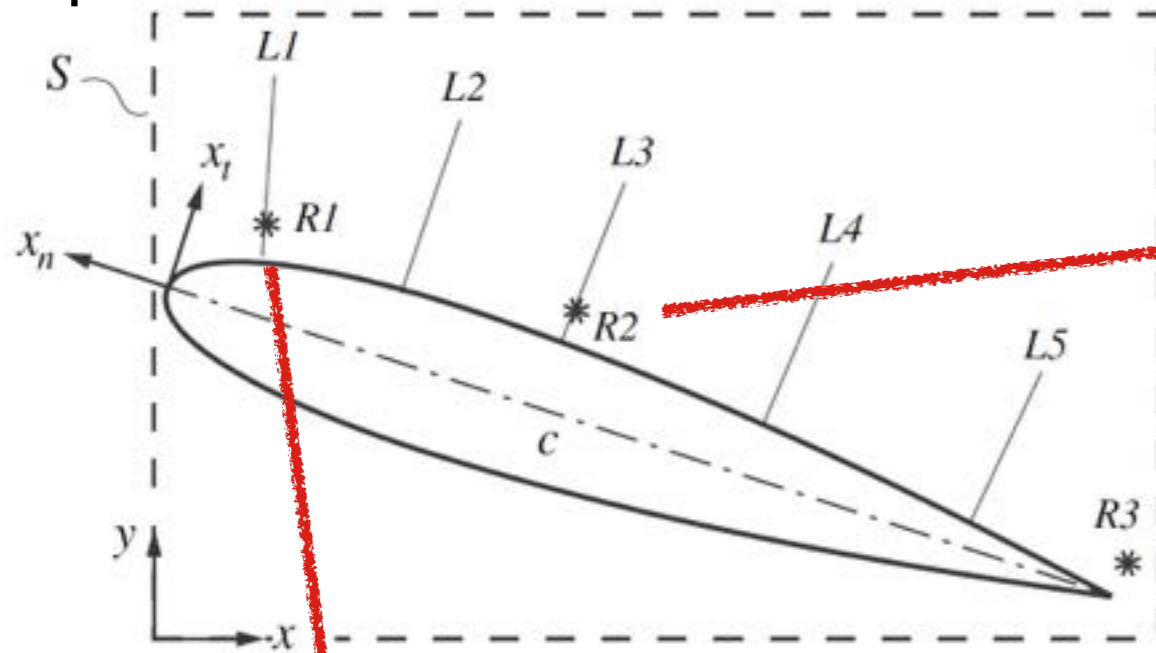
$$f^+ = \frac{f \ell}{U}$$



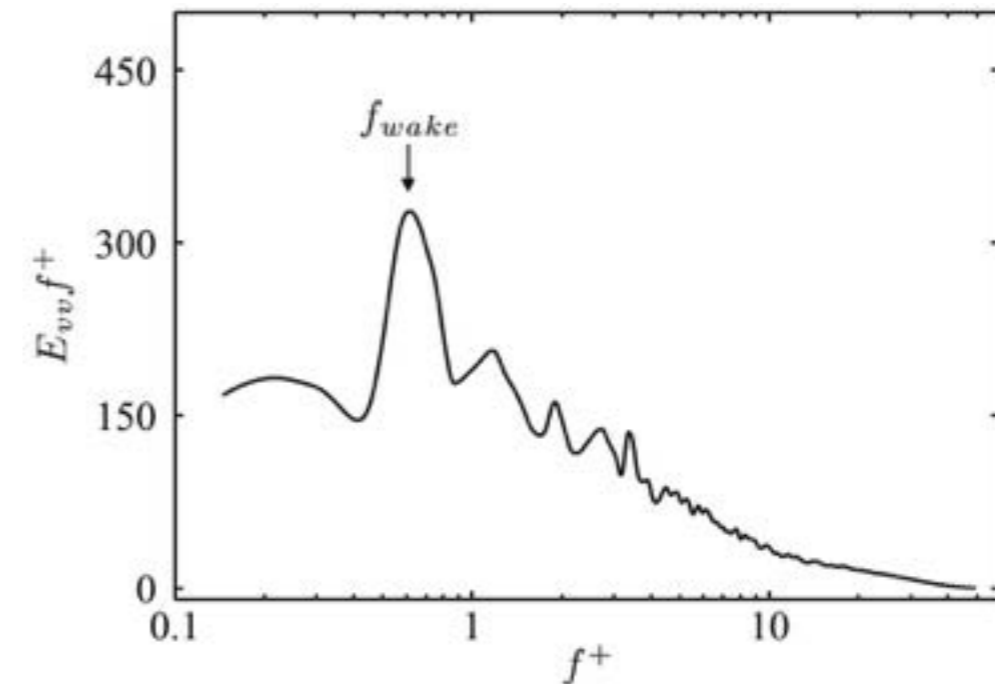
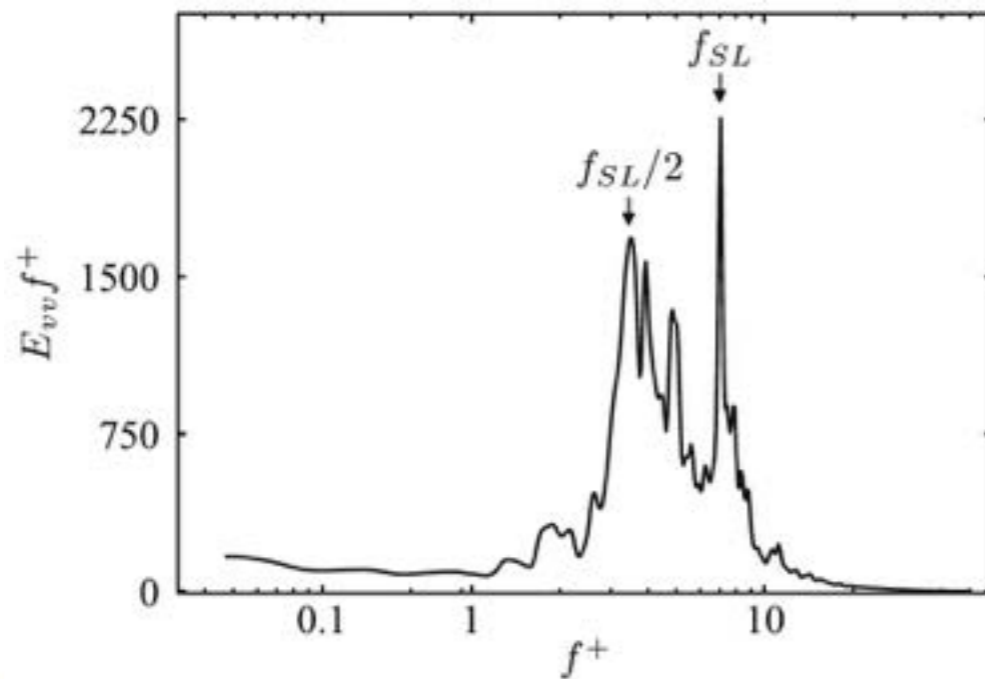
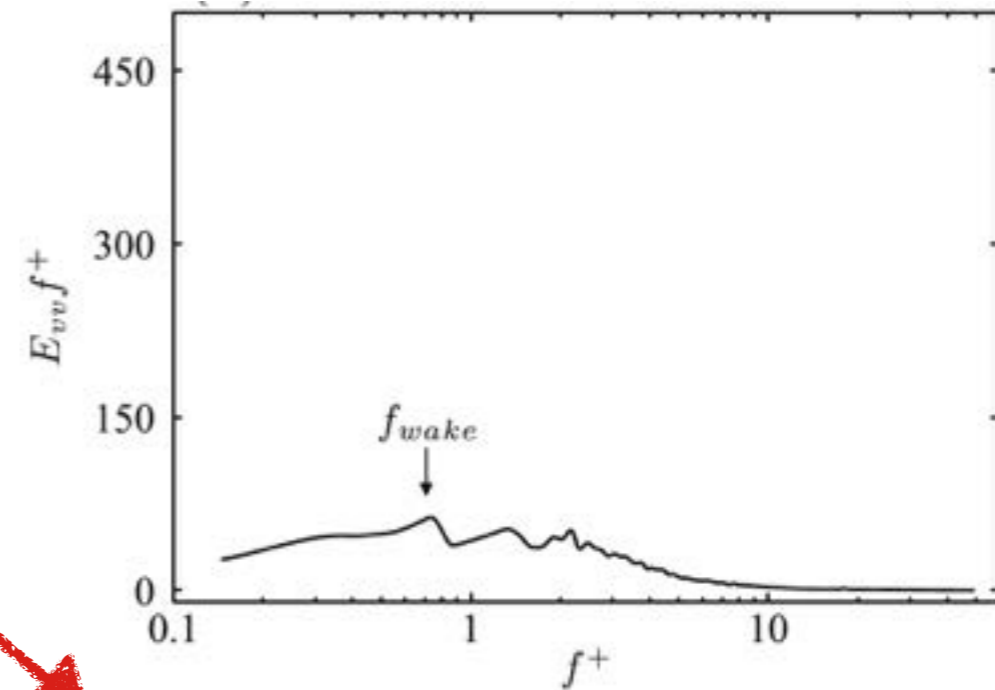
Greenblatt & Wygnanski (2000)

# Instabilities on an Airfoil

Multiple instabilities and flow structures for an airfoil.



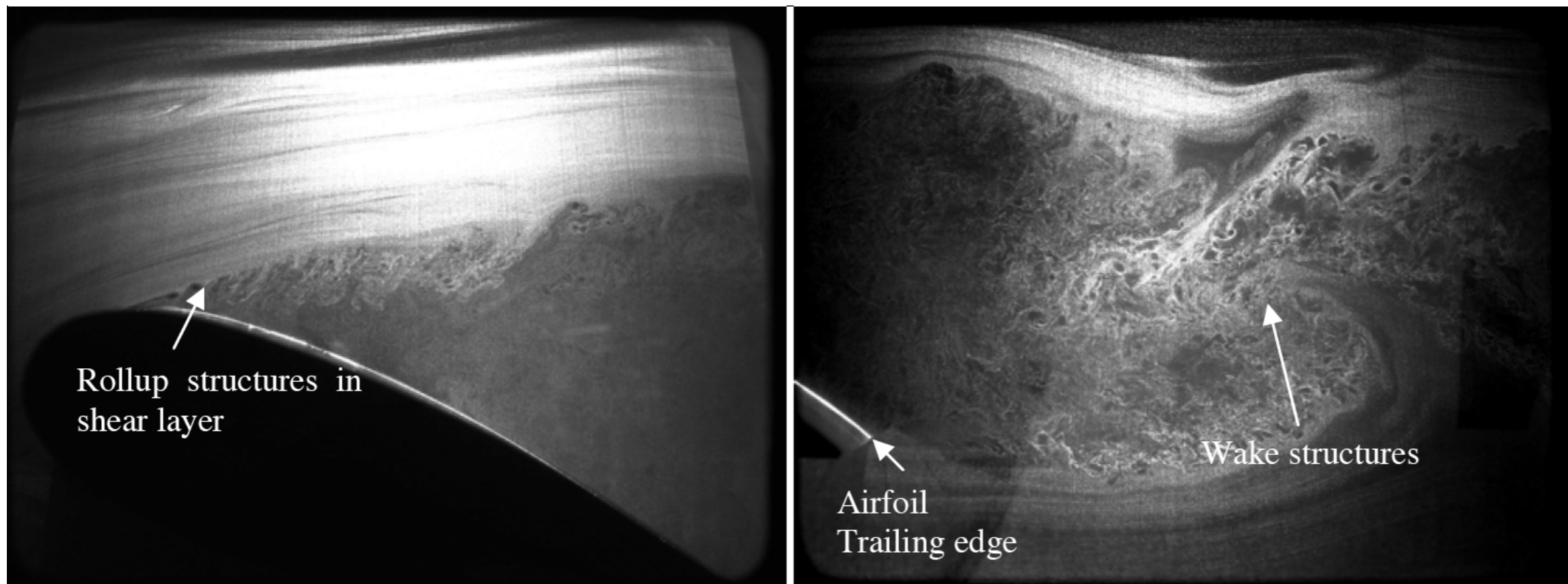
Buchmann et al. (2013)



# Instabilities on an Airfoil



Multiple instabilities and flow structures for an airfoil.



Tian, Cattafesta & Mittal (2006)

Two are of interest today:

- Shear-layer instability - high frequency,  $f^+ = \frac{fc}{U_\infty} = \mathcal{O}(10)$
- Wake instability (von Karman vortices) - low frequency,  $f^+ = \mathcal{O}(1)$

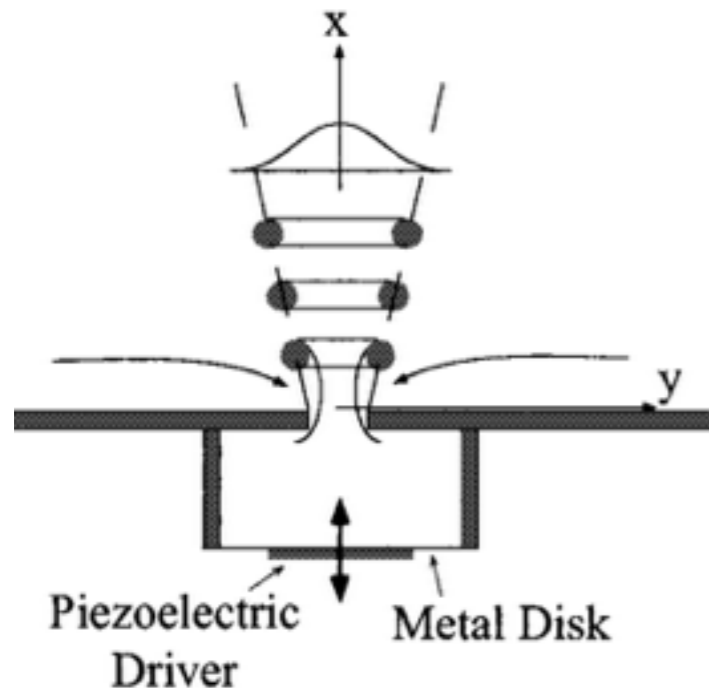
# Separation Control Actuators



There are many types of actuators used for separation control. Will focus here on applications that use two types:

Synthetic jets:

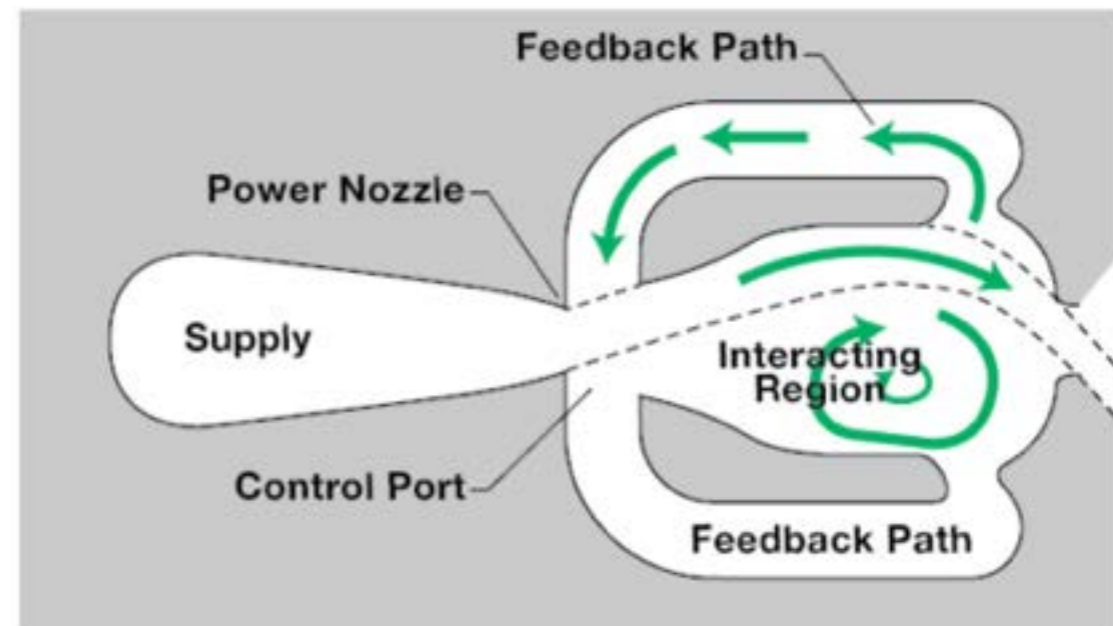
- no fluid source
- high frequency
- moderate velocity
- moving parts



Smith & Glezer (1998)

Fluidic oscillators (or sweeping jets):

- high speed
- no moving parts
- fixed velocity-frequency relation
- external source of fluid



Raman & Raghu (2004)

# Normalised Forcing Amplitude



The forcing amplitude can be expressed either as a momentum or velocity ratio.

Momentum coefficient

$$C_{\mu} = \frac{\text{momentum from actuator}}{\text{momentum of flow}} = \frac{1/2\rho_j U_j A_j}{1/2\rho_{\infty} U_{\infty}^2 A}$$

Velocity ratio

$$V_r = \frac{U_j}{U_{\infty}}$$

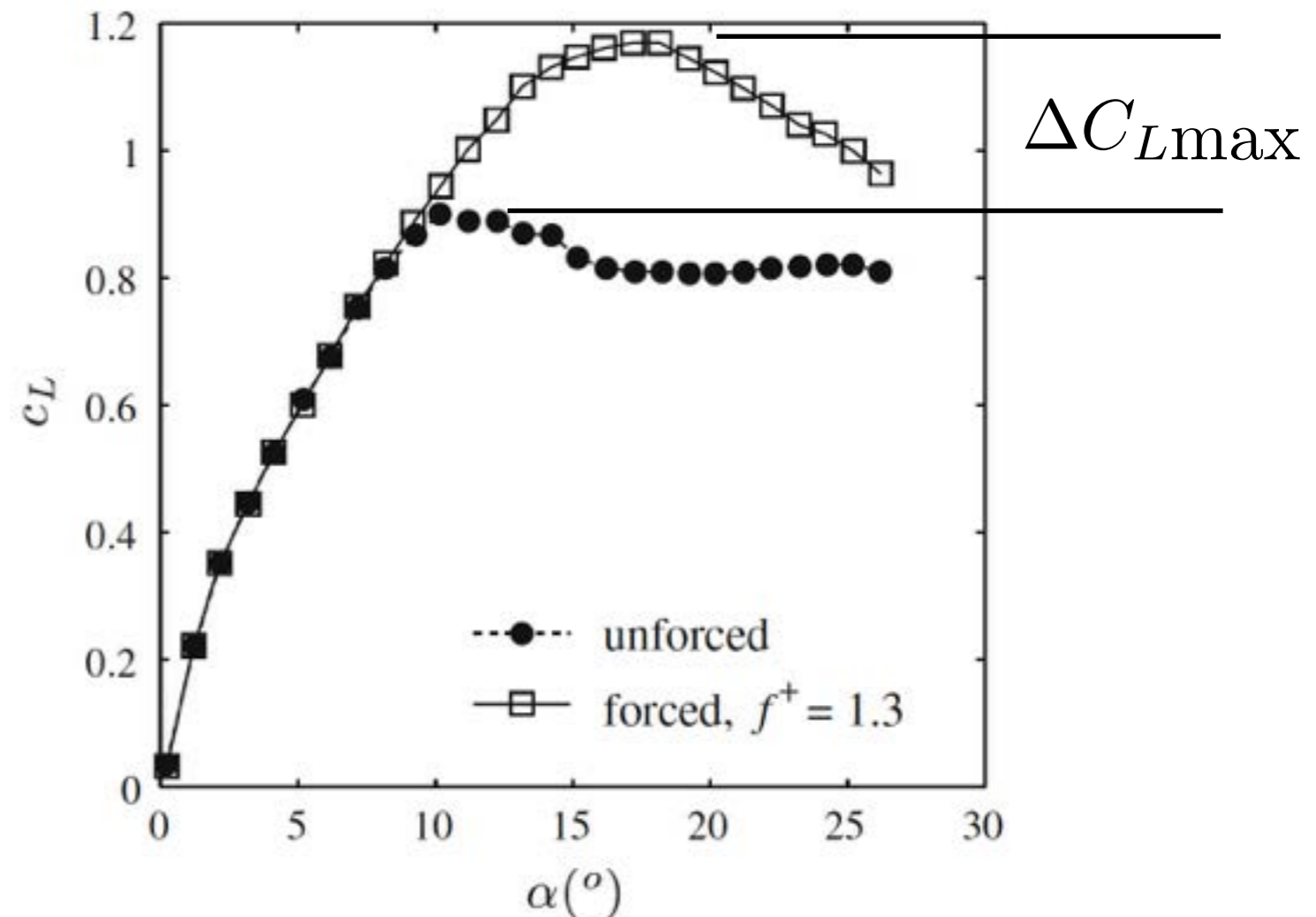
Often not possible to separate which of these is the governing/dominant coefficient to consider. Historically, momentum coefficient is used, but more interest in the effect of velocity ratio in recent years.

# Performance of Separation Control



A few things to expect with control:

- No change at low angles of attack
- Increase in stall angle
- Increase in  $C_{Lmax}$



**Fig. 1** Lift coefficient  $c_L$  of a NACA-0015 airfoil; (*filled circle*) no forcing and (*open square*) ZNMF forcing at the leading edge,  $f^+ = 1.3$ ,  $c_\mu = 0.0014$ . Figure adapted from Tuck and Soria (2008)

# Low Frequency Control

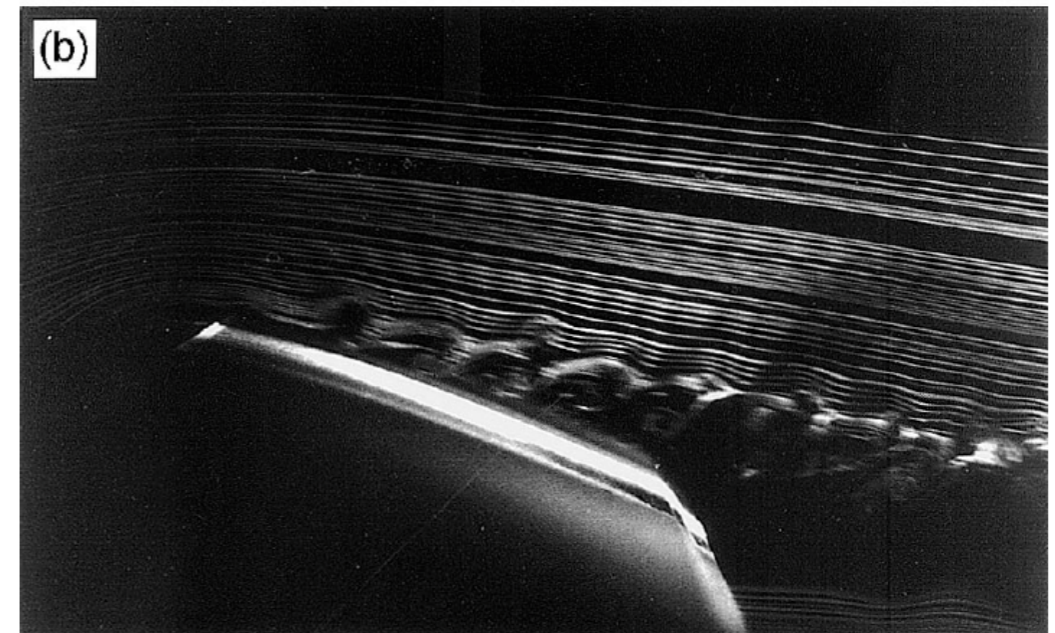
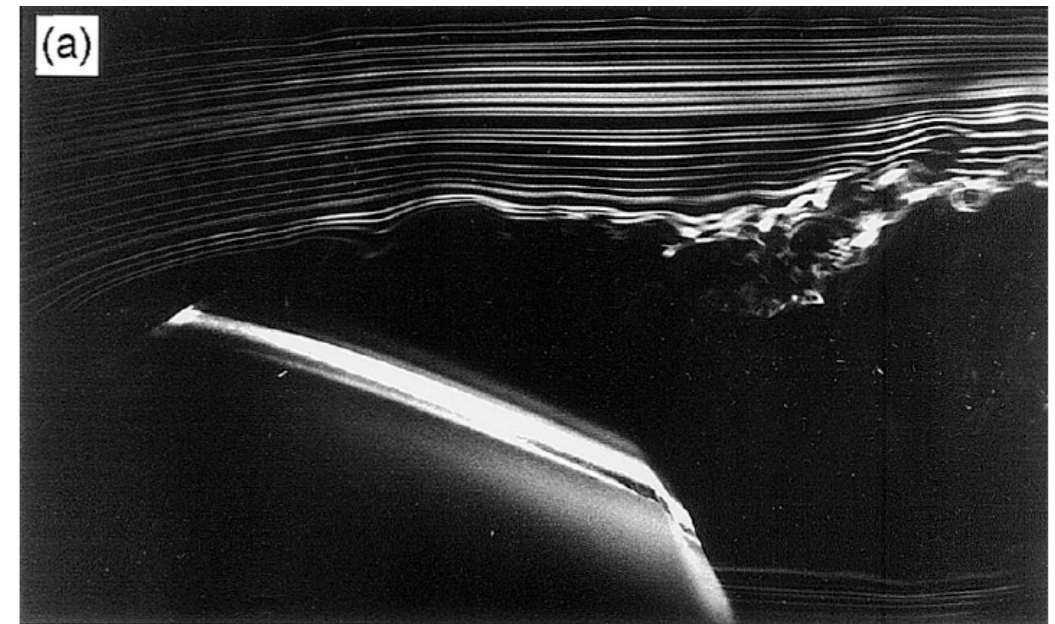


Vibrating ribbon at the leading edge,  
excitation at low frequency  $f^+ = 2.6$ .

Truncated NACA 0015 at  $Re = 40,000$

Flow is reattached. Formation of large  
**coherent structures convecting** over the  
airfoil surface.

Reattachment due to **increase momentum transfer** near the surface due to the  
forcing and resulting coherent structures.



Greenblatt & Wygnanski (2000)



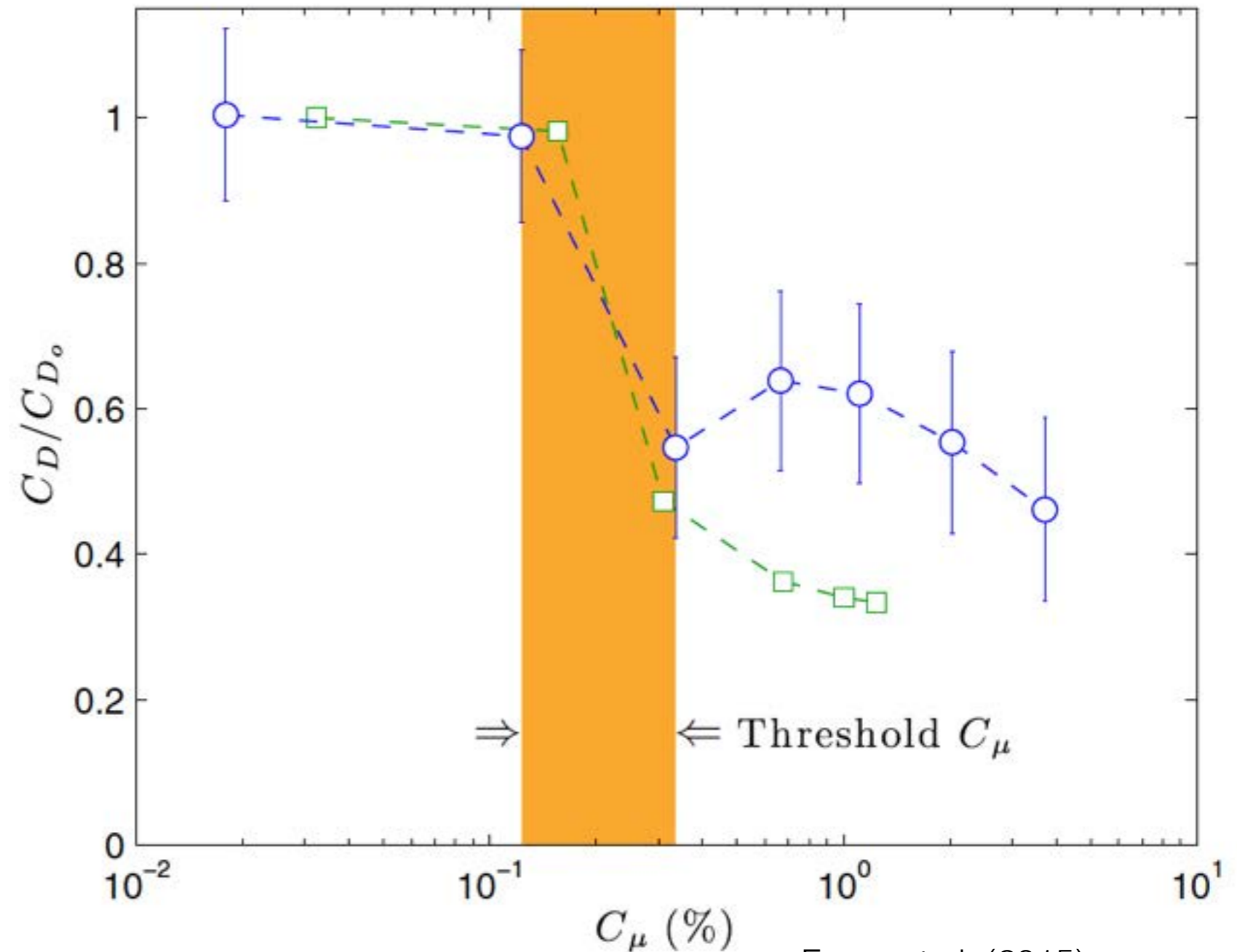
# Effect of Momentum Coefficient



A minimum forcing amplitude is required to reattach the flow.

Beyond threshold, performance remain more or less constant.

Over actuation can lead to small loss of performance.



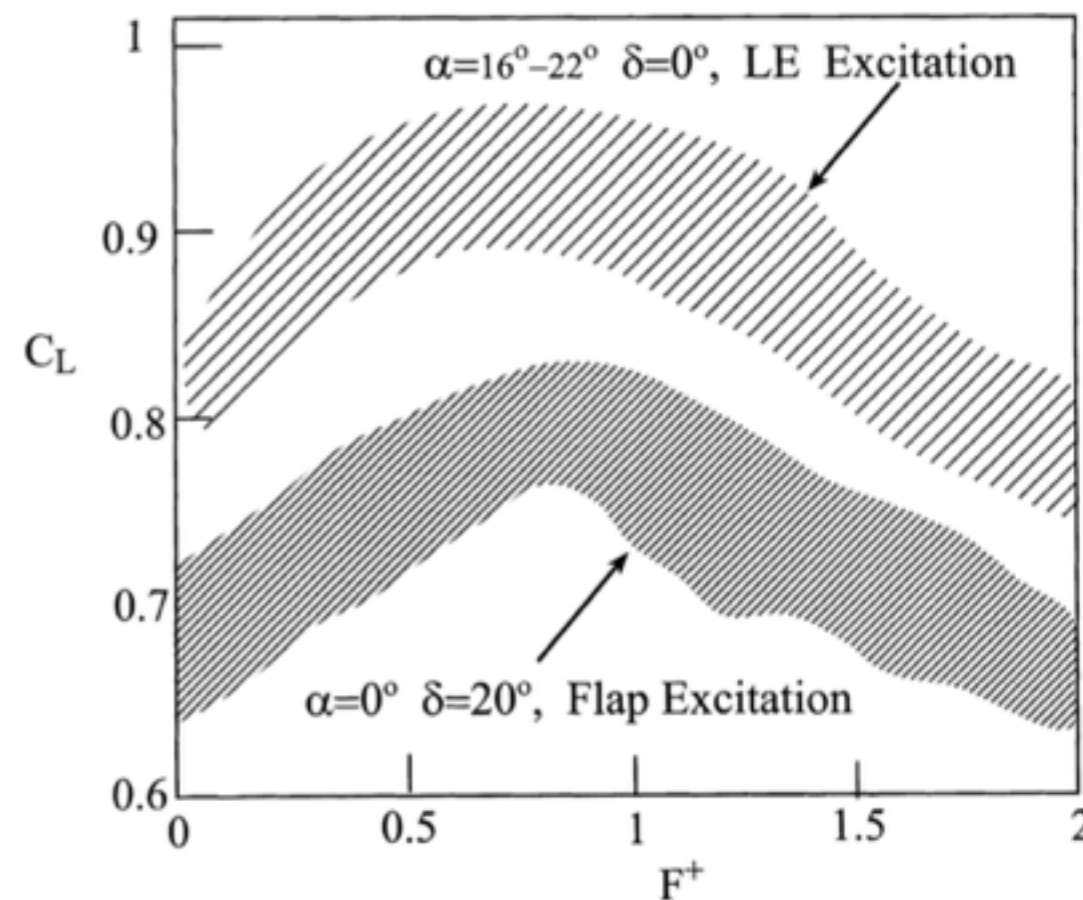
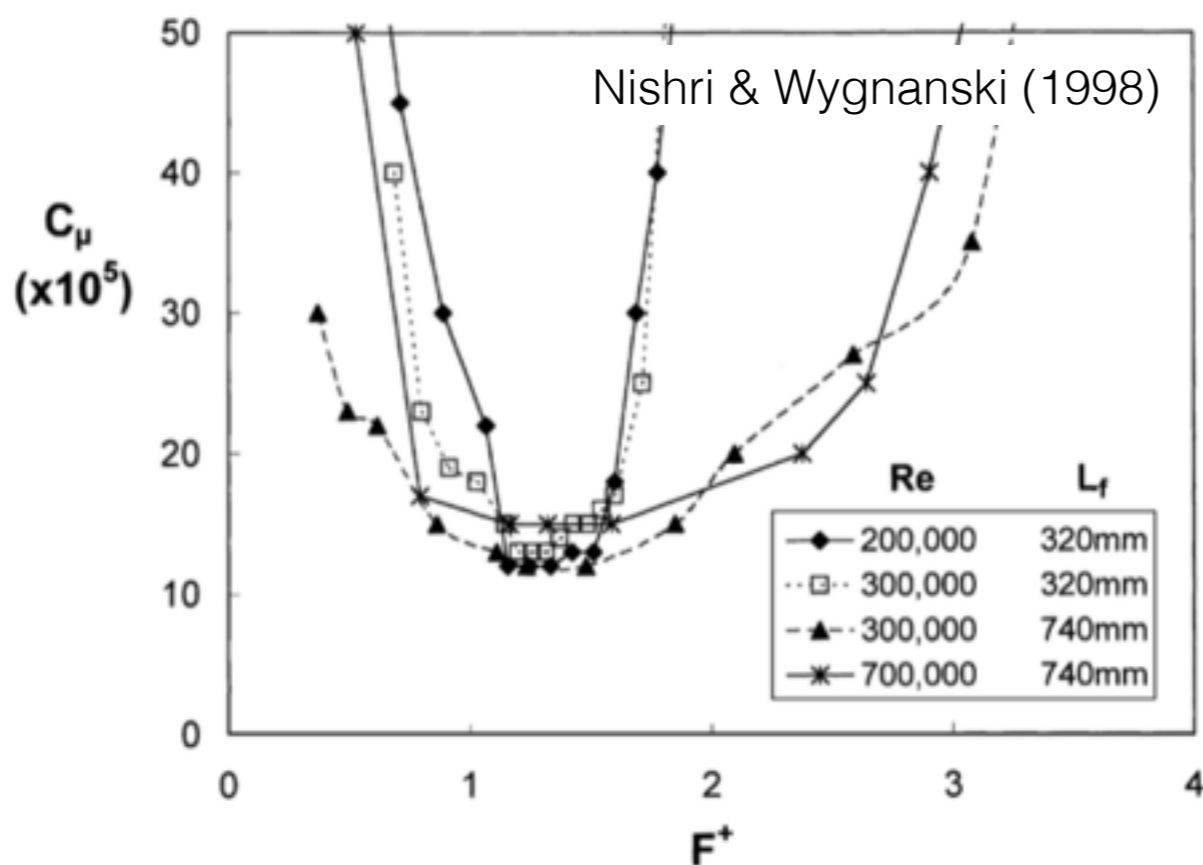
Feero et al. (2015)

# Low Frequency Control



Effect of actuation frequency:

- “narrow-band” efficient - locks onto the receptivity of separated shear layer
- sensitivity reduces with Re



Greenblatt & Wygnanski (2000)

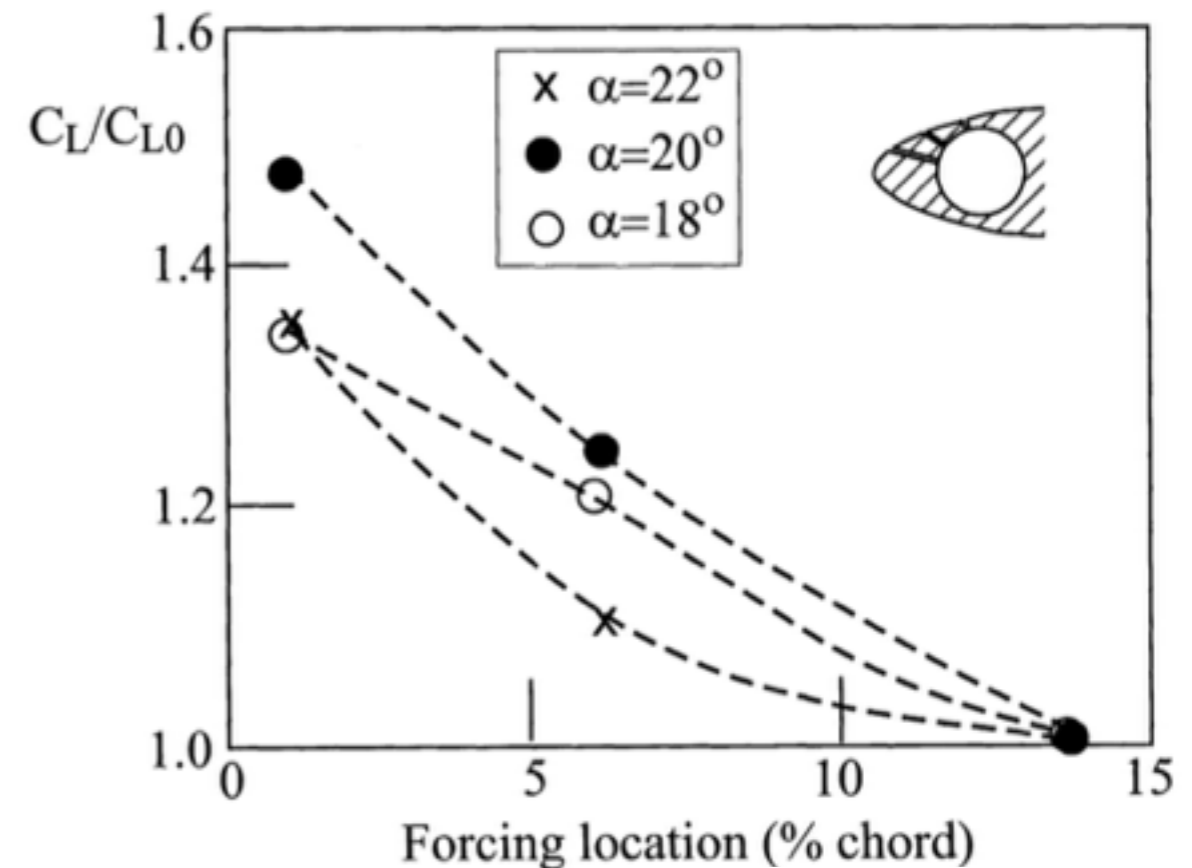
# Effect of Forcing location



Shear layer most receptive at the point of separation

- Forcing upstream will dissipate
- Forcing downstream needs to overpower the separated region to be effective

Unfortunately, some lack of systematic evidence of control effectiveness with forcing location.



Hsiao et al. (1994)

NACA 63<sub>3</sub>-018 with leading edge stall  
Caveat: forcing angle not constant

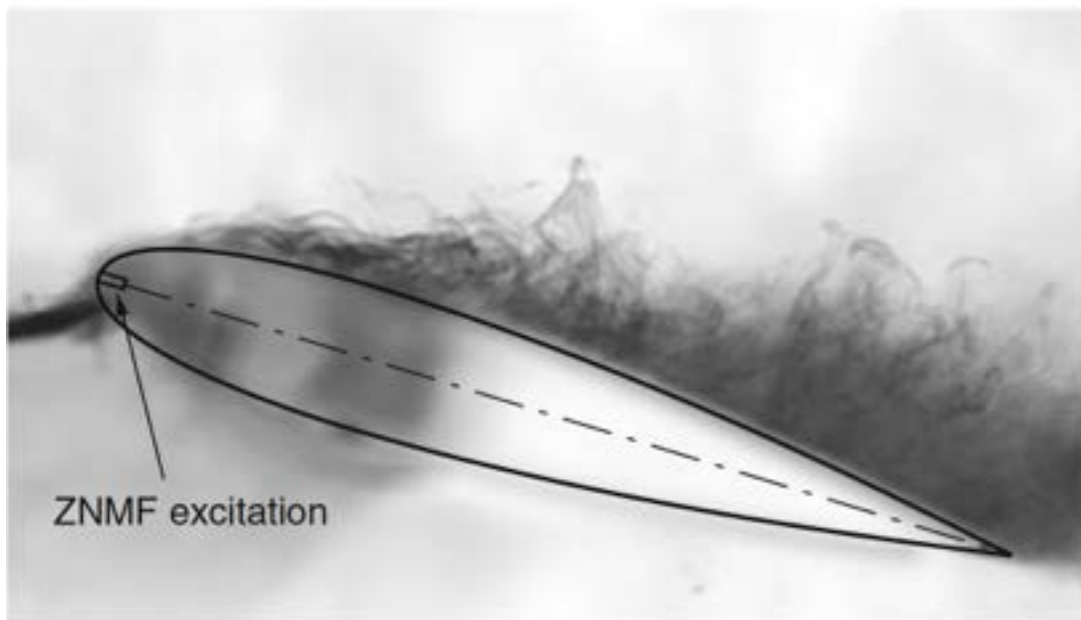
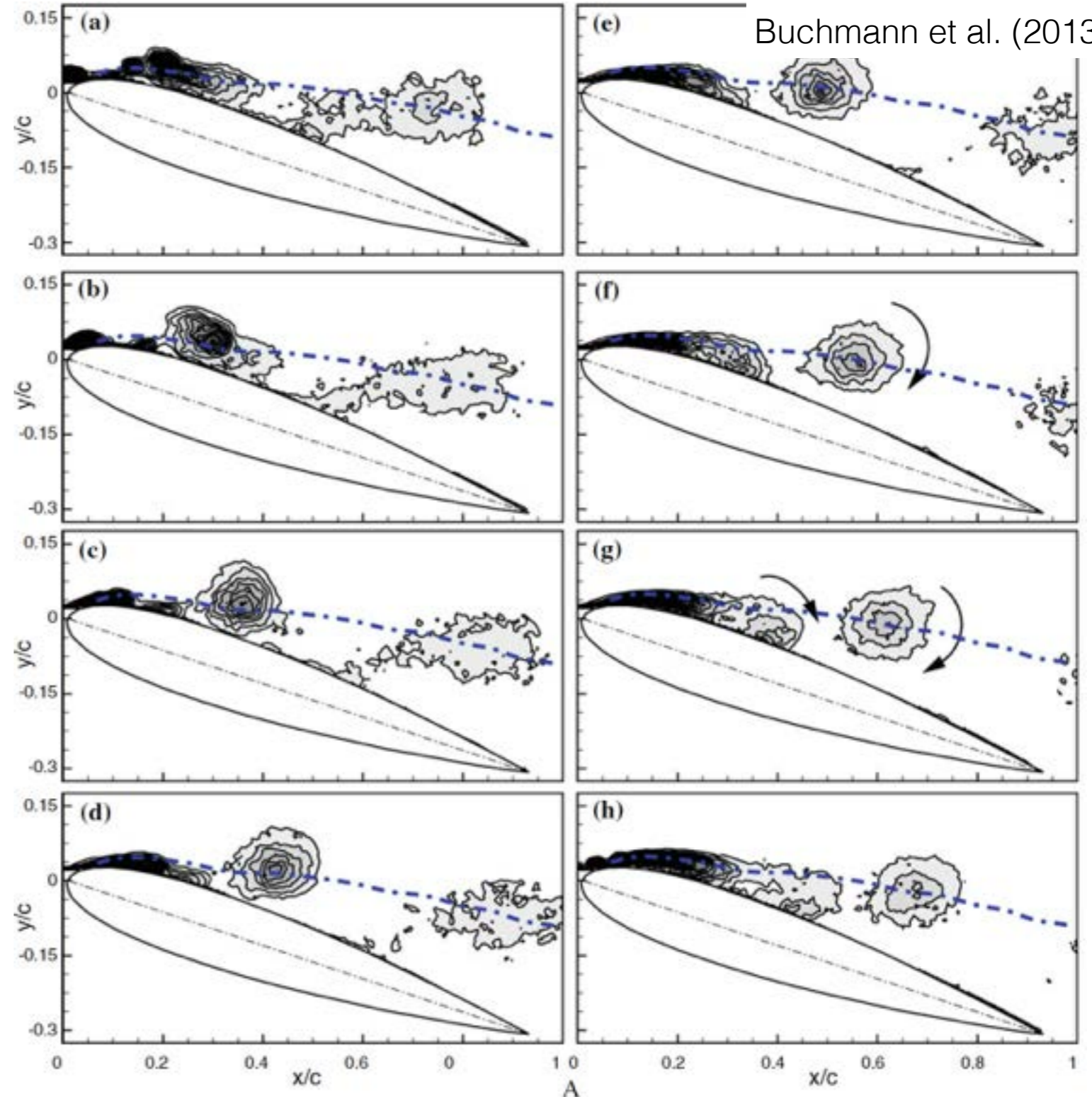
# Unsteadiness of the Control



NACA 0015 at  $Re = 10^4$   
and  $18^\circ$  angle of attack.  
Synthetic jet at the LE

$$f^+ = 1.3 \quad \& \quad C_\mu = 1.4\%$$

Phase-averaged vorticity

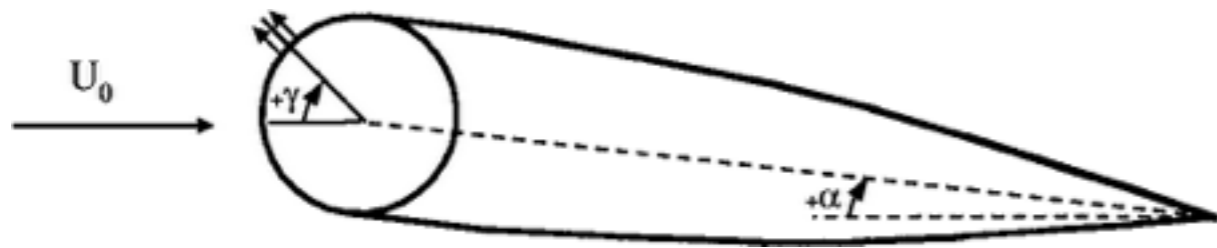


# High Frequency Forcing



Work by Glezer and co-workers using high frequency excitation found interesting results on an unconventional airfoil.

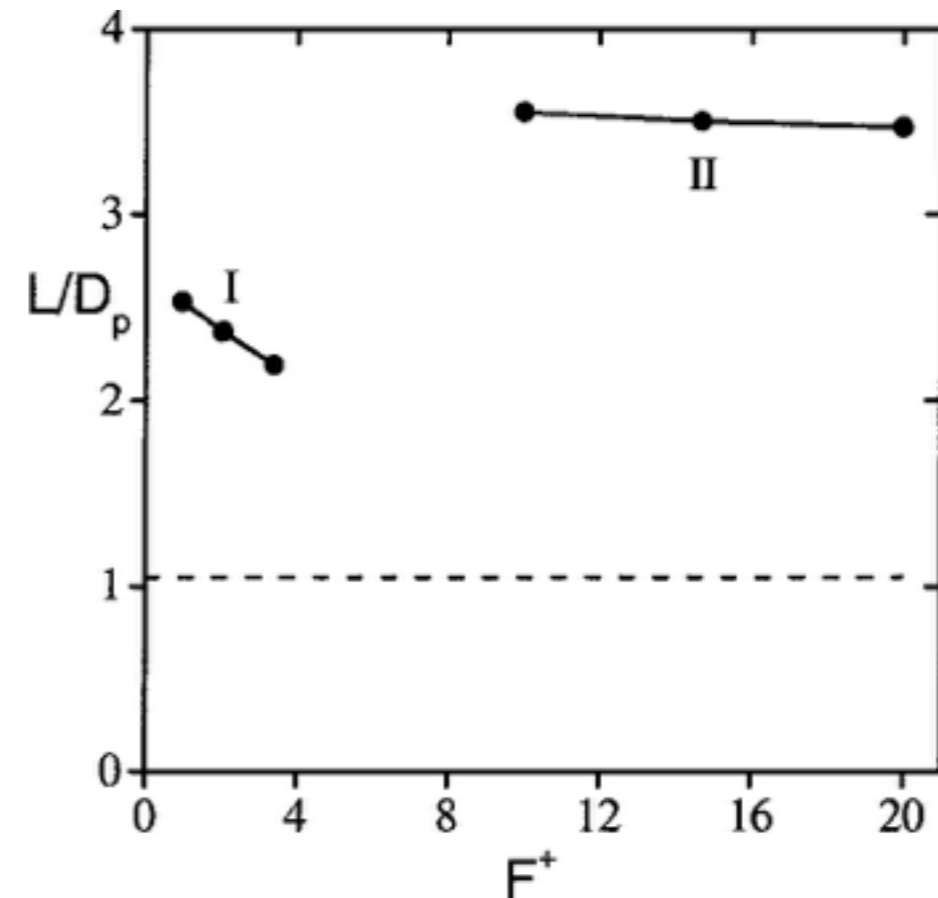
- Streamlined cylinder made to readily change location of excitation.



$$Re_c = 3.1 \times 10^5$$

$$C_\mu = 3.5 \times 10^{-3}$$

High frequency excitation, uncoupled from the natural frequency of the separated shear layer.



Amitay et al. (2001)

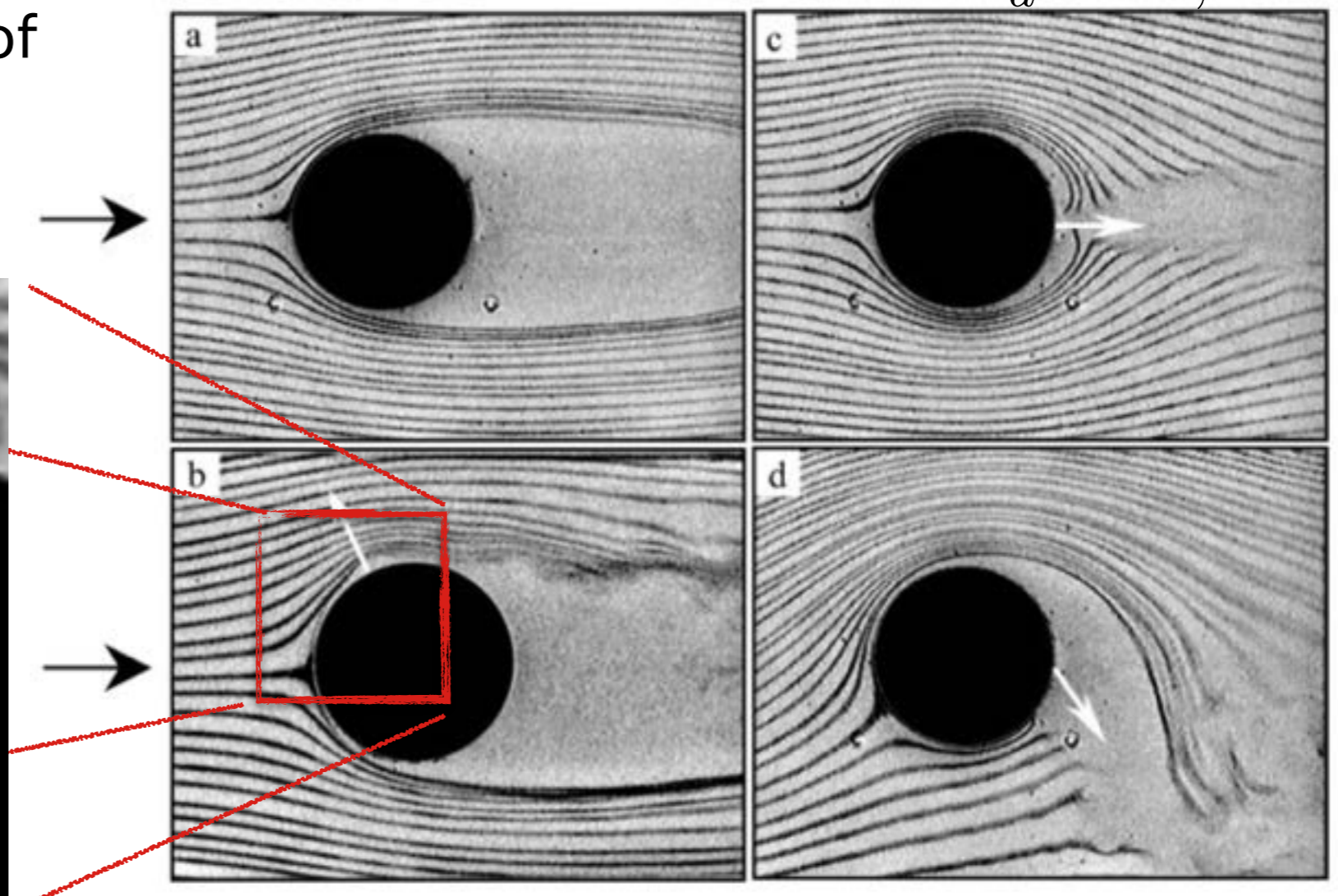
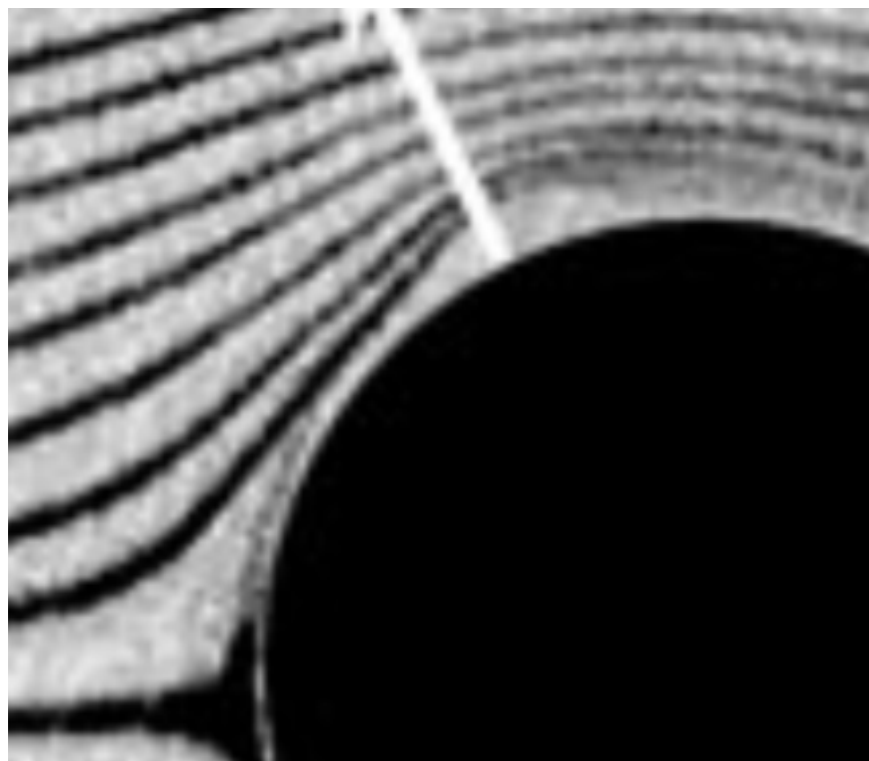
# Virtual Aerodynamic Shaping



At high frequency, the forcing can appear quasi-steady to the flow if the highest times scale of the flow is much smaller than that of the forcing.

$$Re_d = 40,000$$

Notice the small deviation of the streamlines in Fig. 7(b)



**Figure 7** Smoke of the flow around a circular cylinder visualization: (a) baseline; and (b) actuated:  $\phi = 0$ ,  $\gamma = 60^\circ$  and (c)  $180^\circ$ , and (d)  $\phi = 120^\circ$ ,  $\gamma = 180^\circ$ .

# Virtual Aerodynamic Shaping



Flow near the exit scales  
momentum flux ratios and jet  
width.

For low momentum ratio and  
Strouhal number, discrete  
vortices shed.

For high momentum ratio and  
Strouhal number, a closed  
recirculation zone forms.

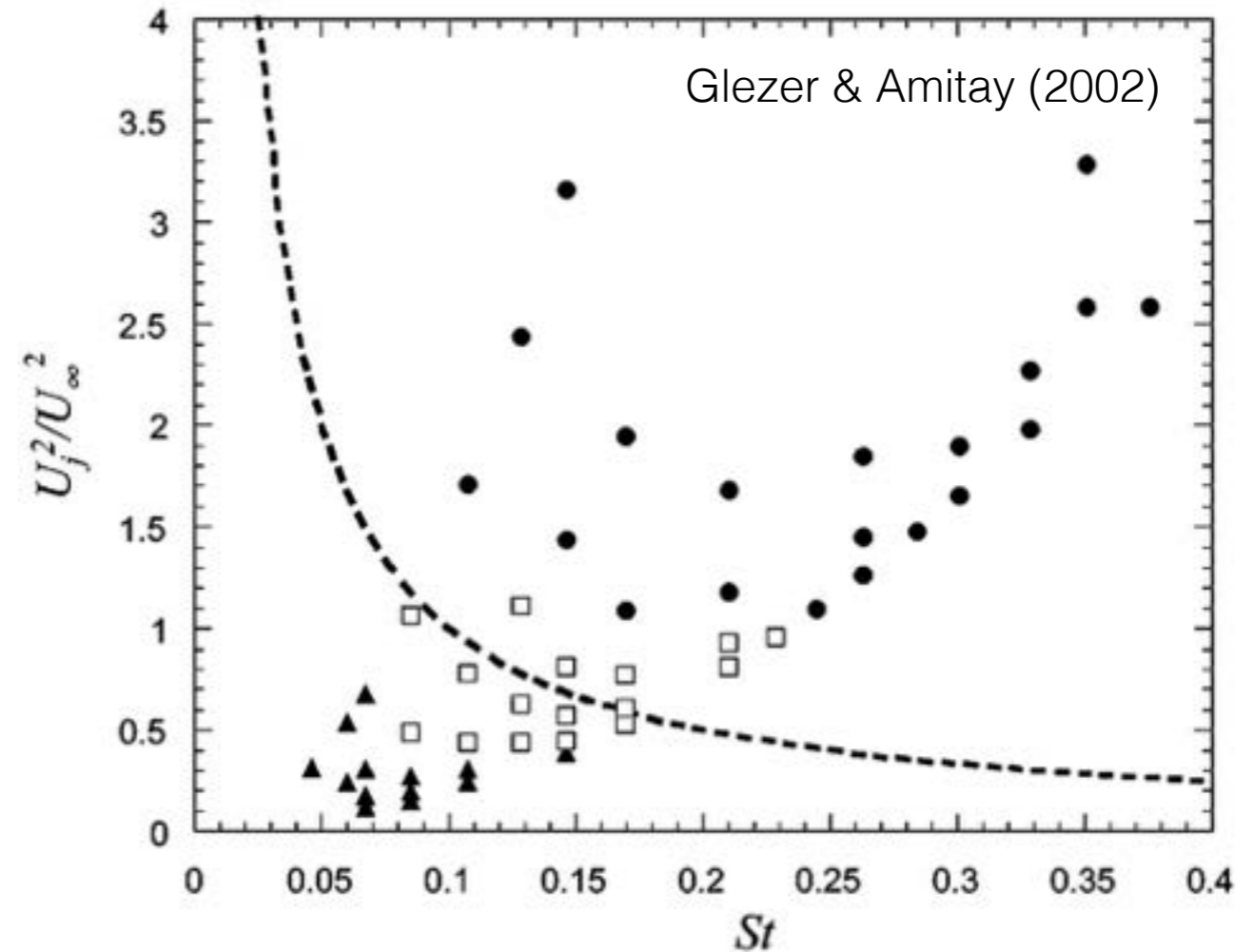
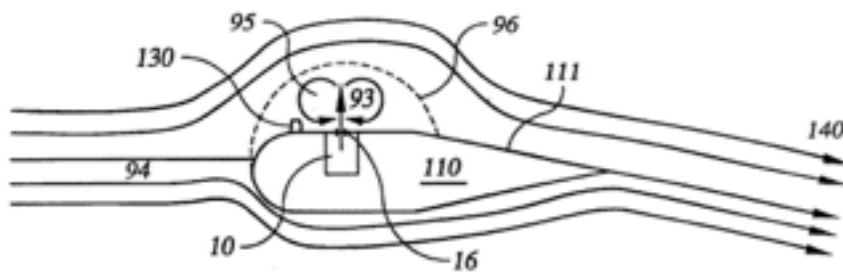


Figure 12 Interaction domain map. ( $\blacktriangle$ ) Discrete vortices, ( $\square$ ) transitory, ( $\bullet$ ) closed recirculation. Dashed line corresponds to  $\hat{f} = 0.1$ .



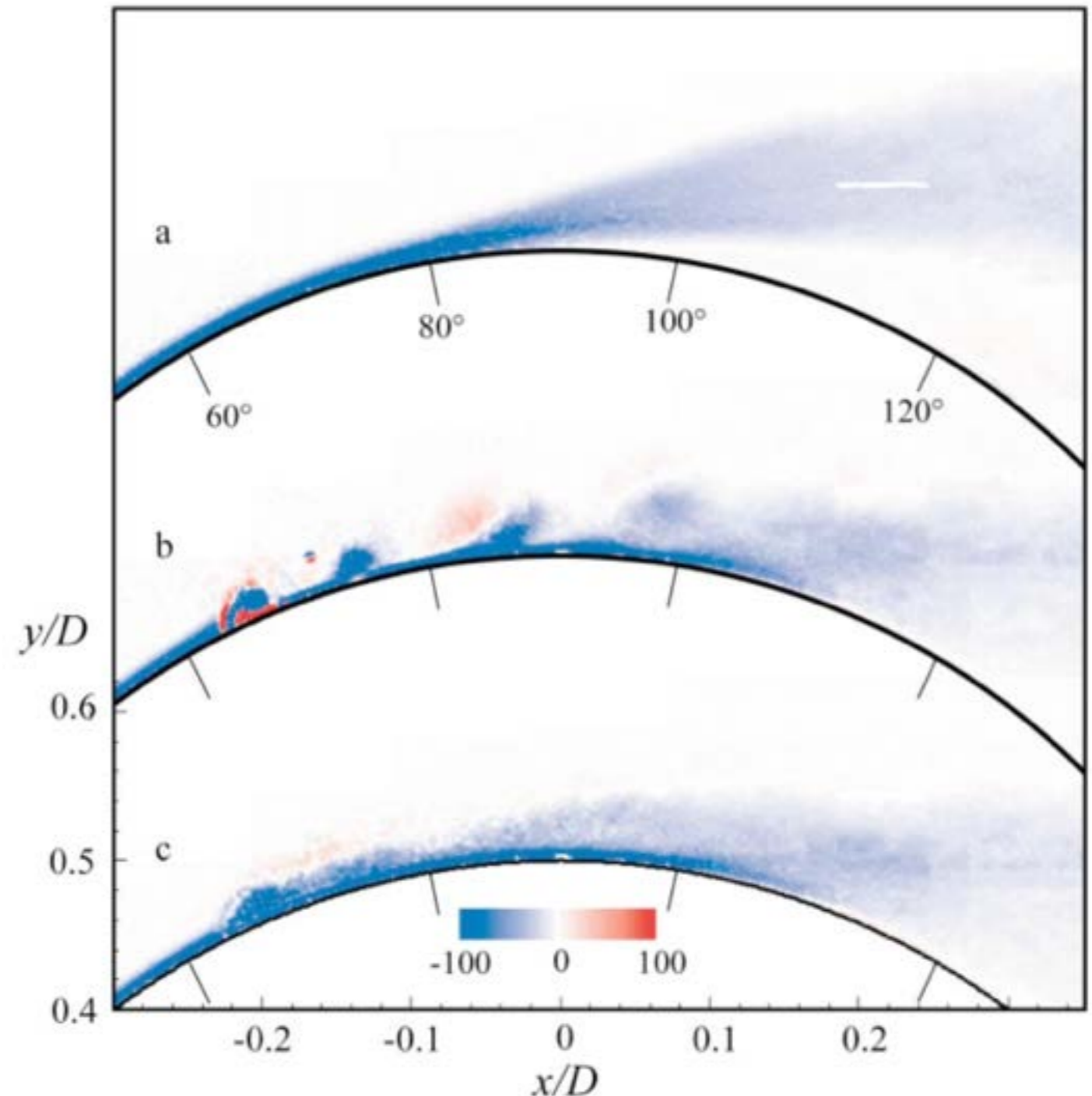
# Virtual Aerodynamic Shaping



Baseline flow on top. Separates at about  $83^\circ$

With excitation, separation moves to  $110^\circ$

CW vortices convect along the surface, transporting high momentum fluid near the wall



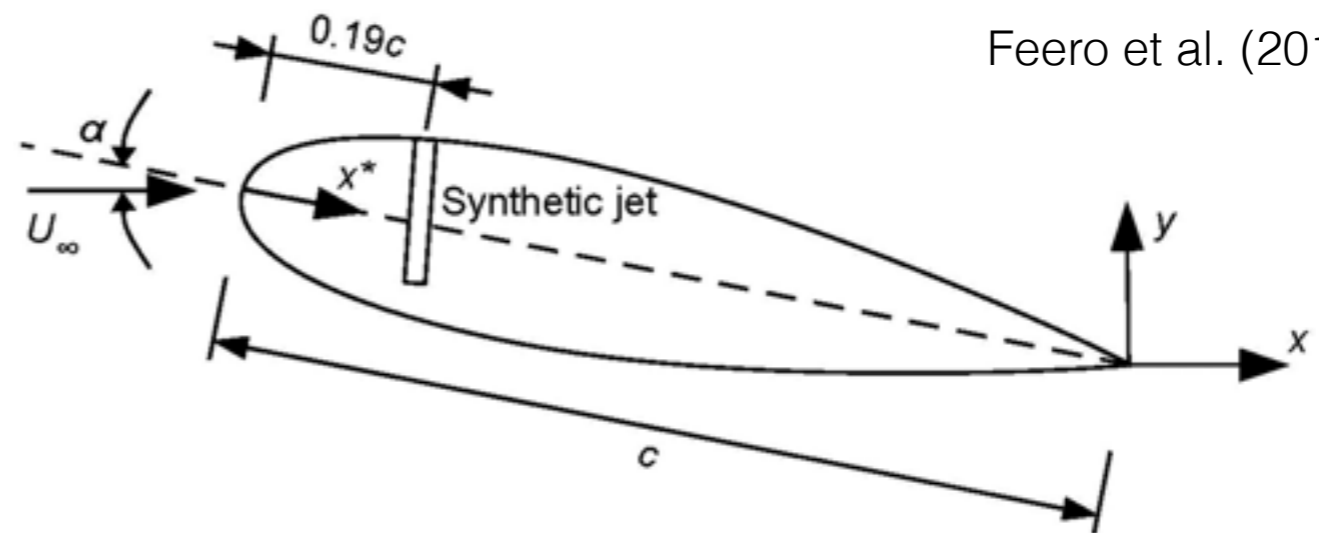
**Figure 14** Normalized vorticity.  $Re_D = 21,500$ ,  $\gamma = 63^\circ$ ,  $C_\mu = 5.1 \times 10^{-2}$ ,  $\hat{f} = 0.035$ .  
(a) baseline; actuated: (b) phase locked and (c) time averaged.



# Low vs High Frequency Forcing

NACA 0025 at  $Re = 10^5$   
and  $10^\circ$  angle of attack  
Forcing at separation

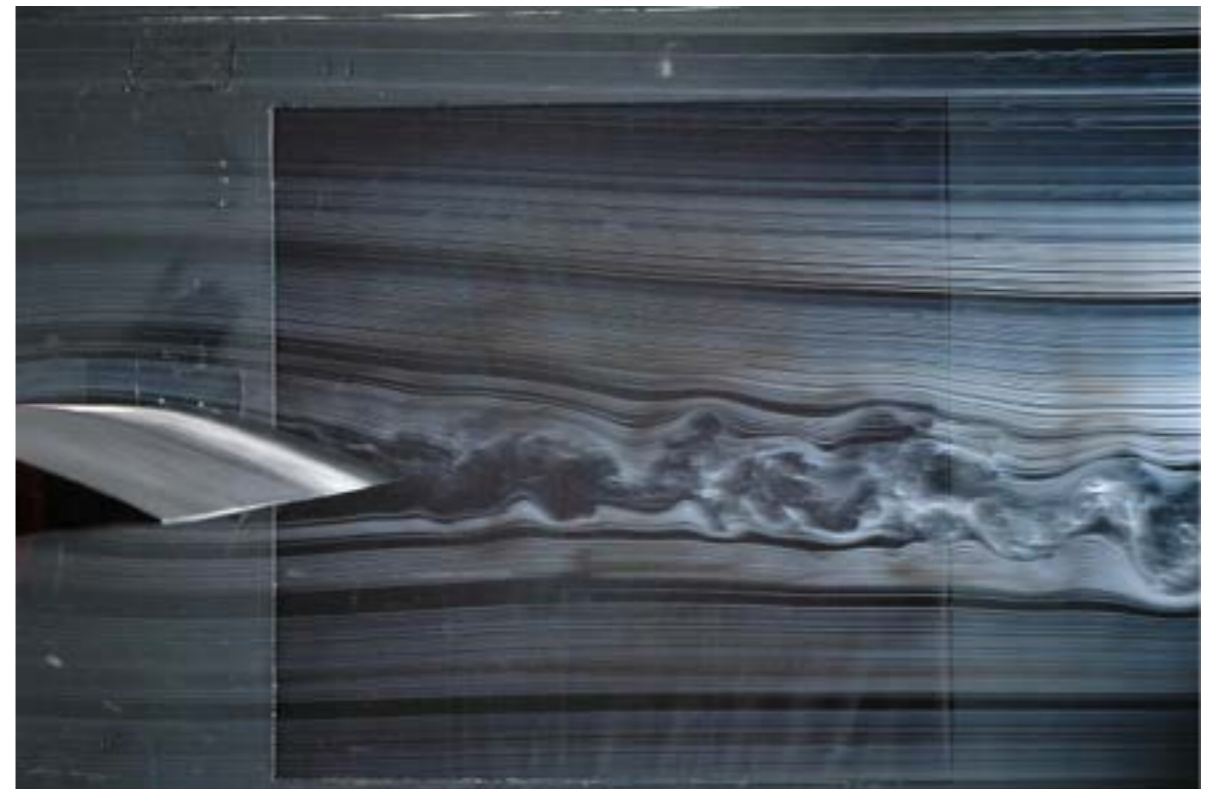
$$St_e (= f^+) = 58$$



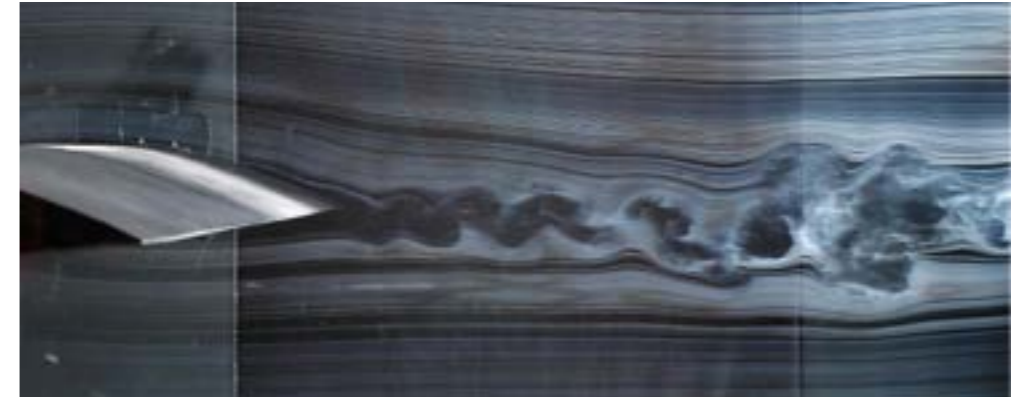
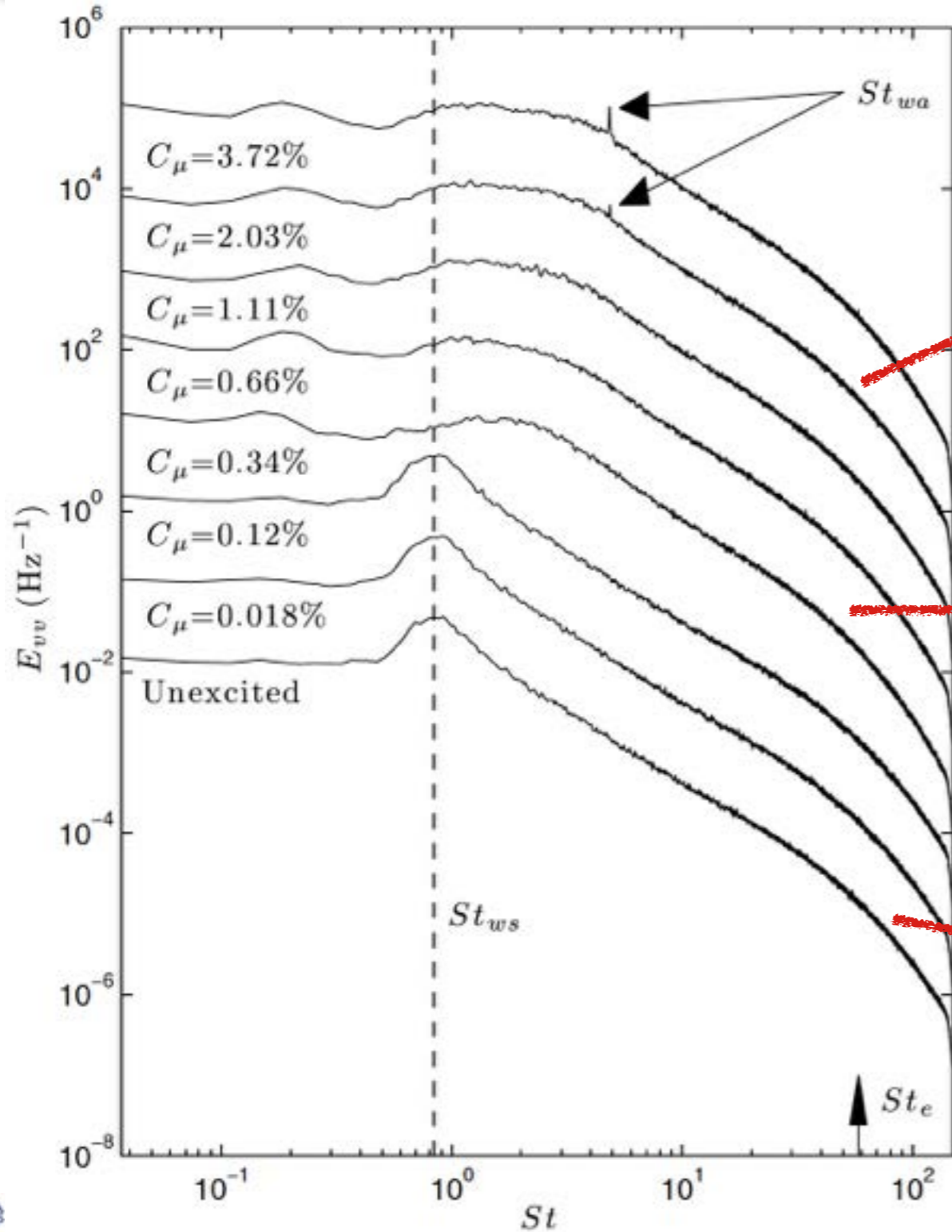
Uncontrolled



Forced

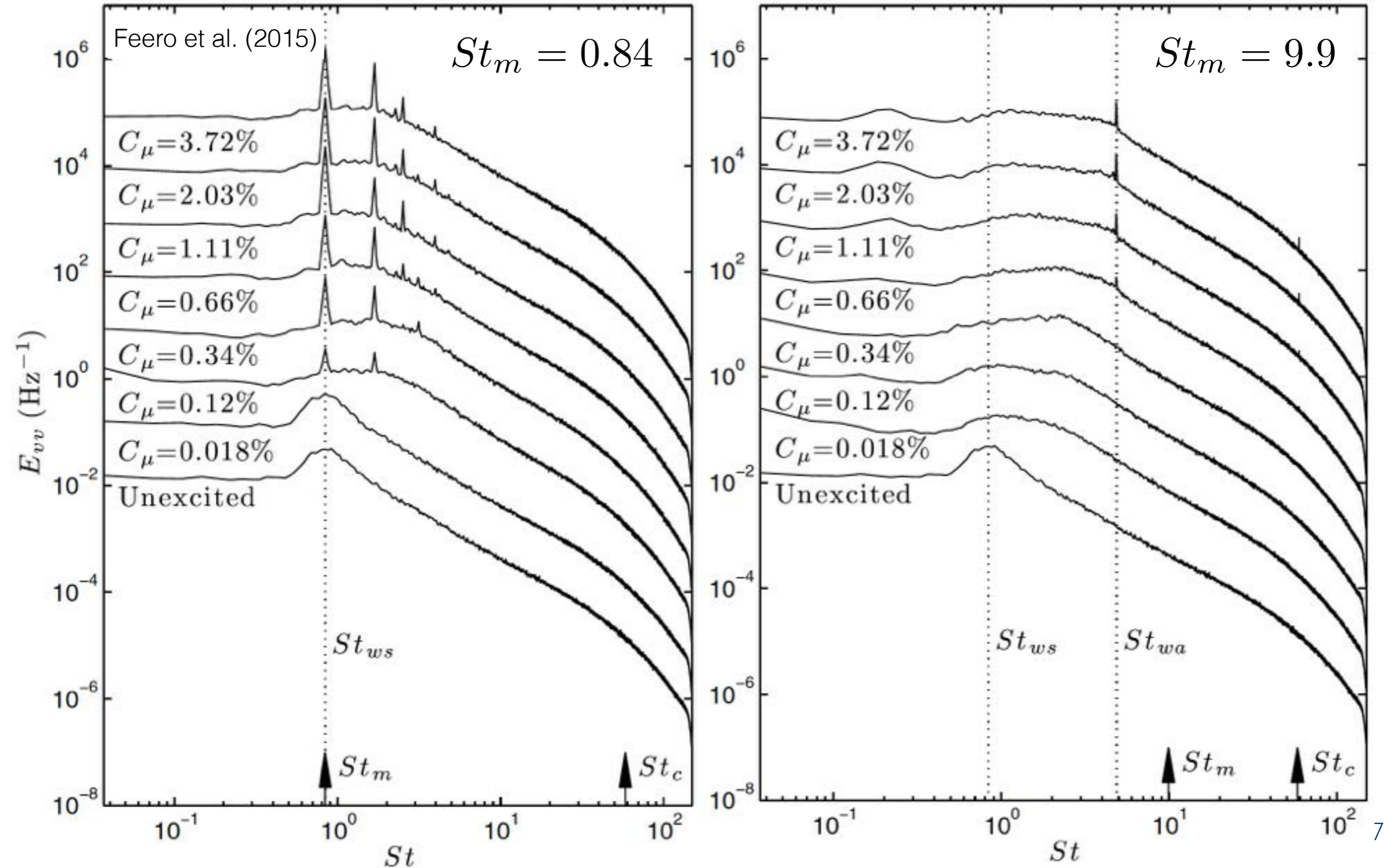


# Low vs High Frequency Forcing



Feero et al. (2015)

# Pulse Modulated Forcing



# Low vs High Frequency Forcing



- Low frequency forcing ( $f^+ = \mathcal{O}(1)$ ) produces a stronger periodic component in the flow
  - the train of vortices observed previously
  - can produce more “unsteady” aerodynamic forces
  - flow reattached at lower forcing than  $f^+ = \mathcal{O}(100)$
- High frequency forcing ( $f^+ = \mathcal{O}(10)$ ) produces more steady reattachment
  - No significant periodicity in the wake except with over actuation
  - Reattachment at lower forcing amplitude as other frequency cases
- Non-linear interactions between structures of different scales play an important and complex role in the reattachment of the flow.

# NASA/Boeing ERA Control



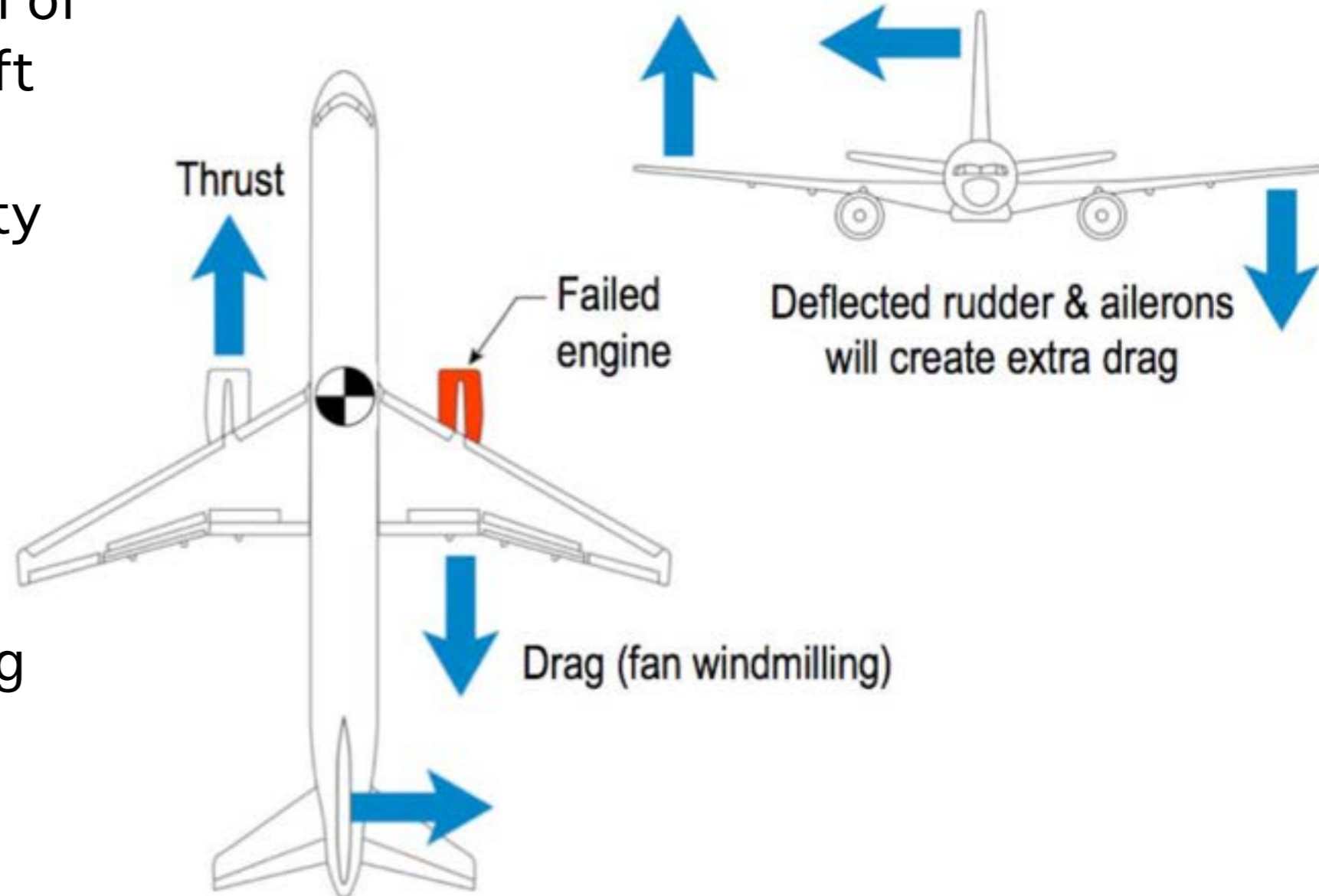
Active flow control used on vertical tail.

Lin et al. (2016)

Good point for introduction of technology on a real aircraft

- non-critical application
- can demonstrate reliability

Potential for about 1% drag reduction.

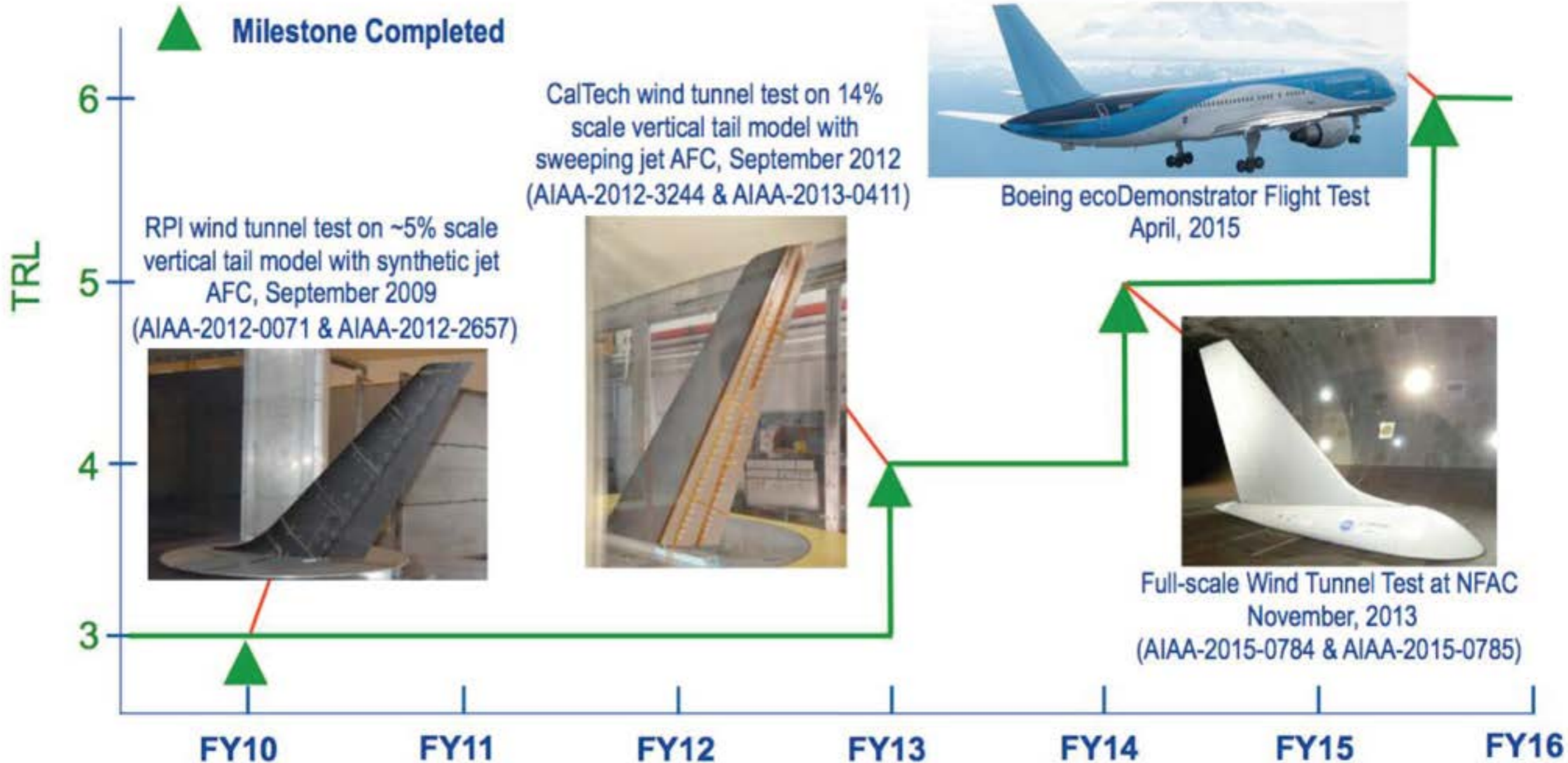


# NASA/Boeing ERA Control



Multi-scale, multi-stage effort to move technology from the lab to flight!

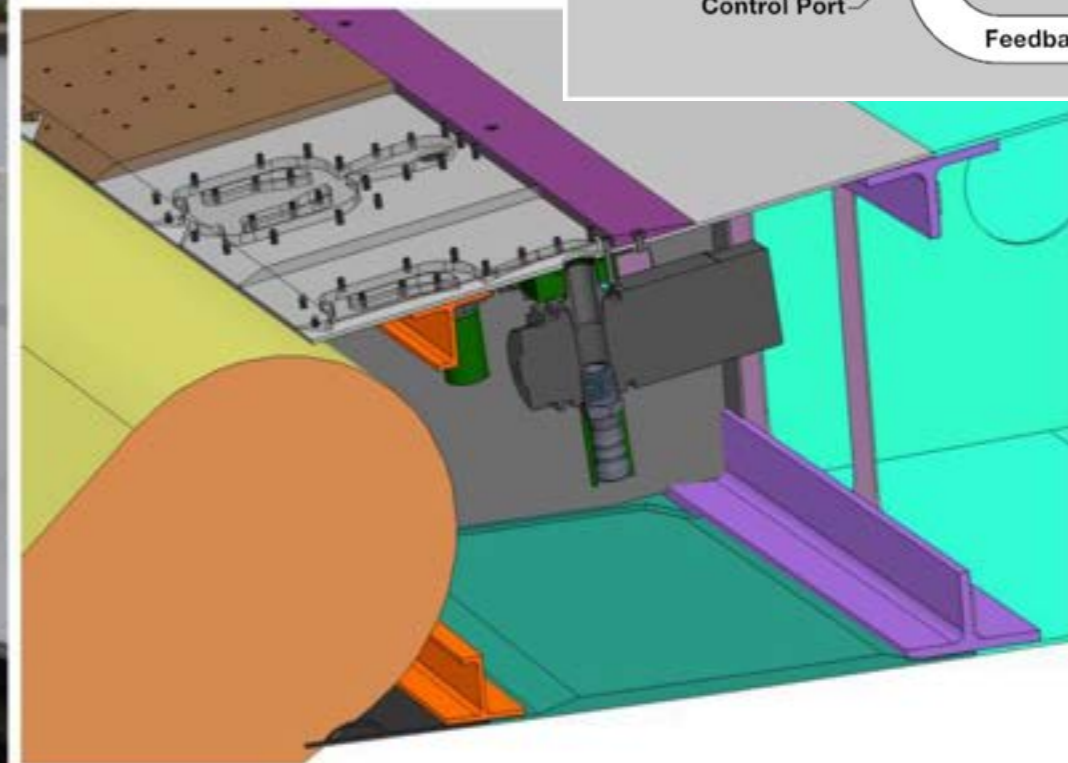
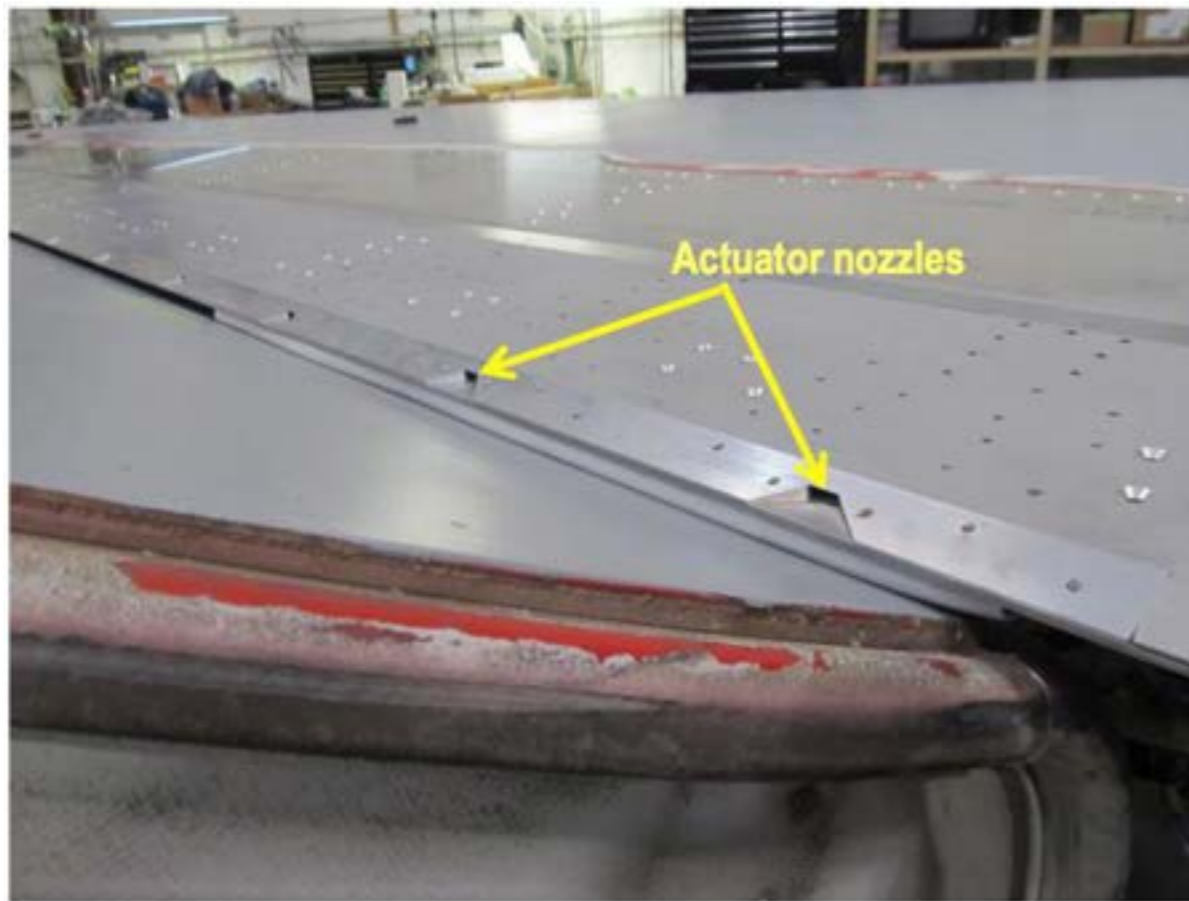
Lin et al. (2016)



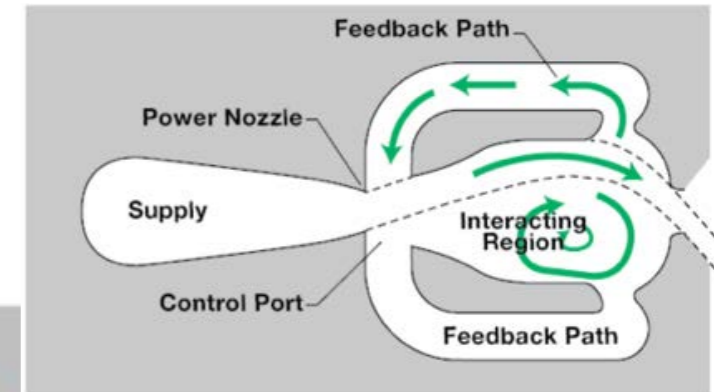
# NASA/Boeing ERA Control



For full-scale wind tunnel testing, 37 fluidic oscillators are applied upstream of the rudder hinge.



Close-up view of actuators



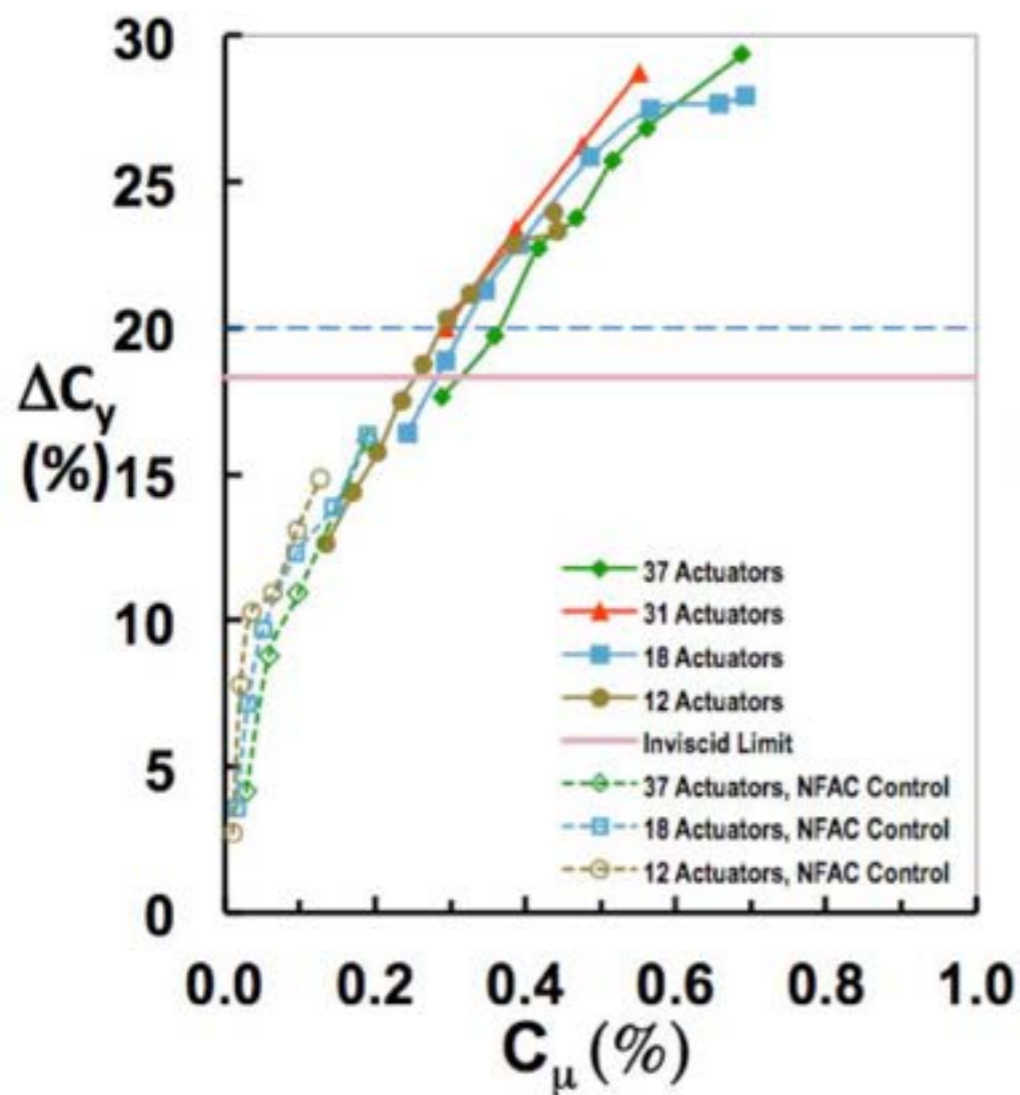
Whalen et al. (2015)

# NASA/Boeing ERA Control

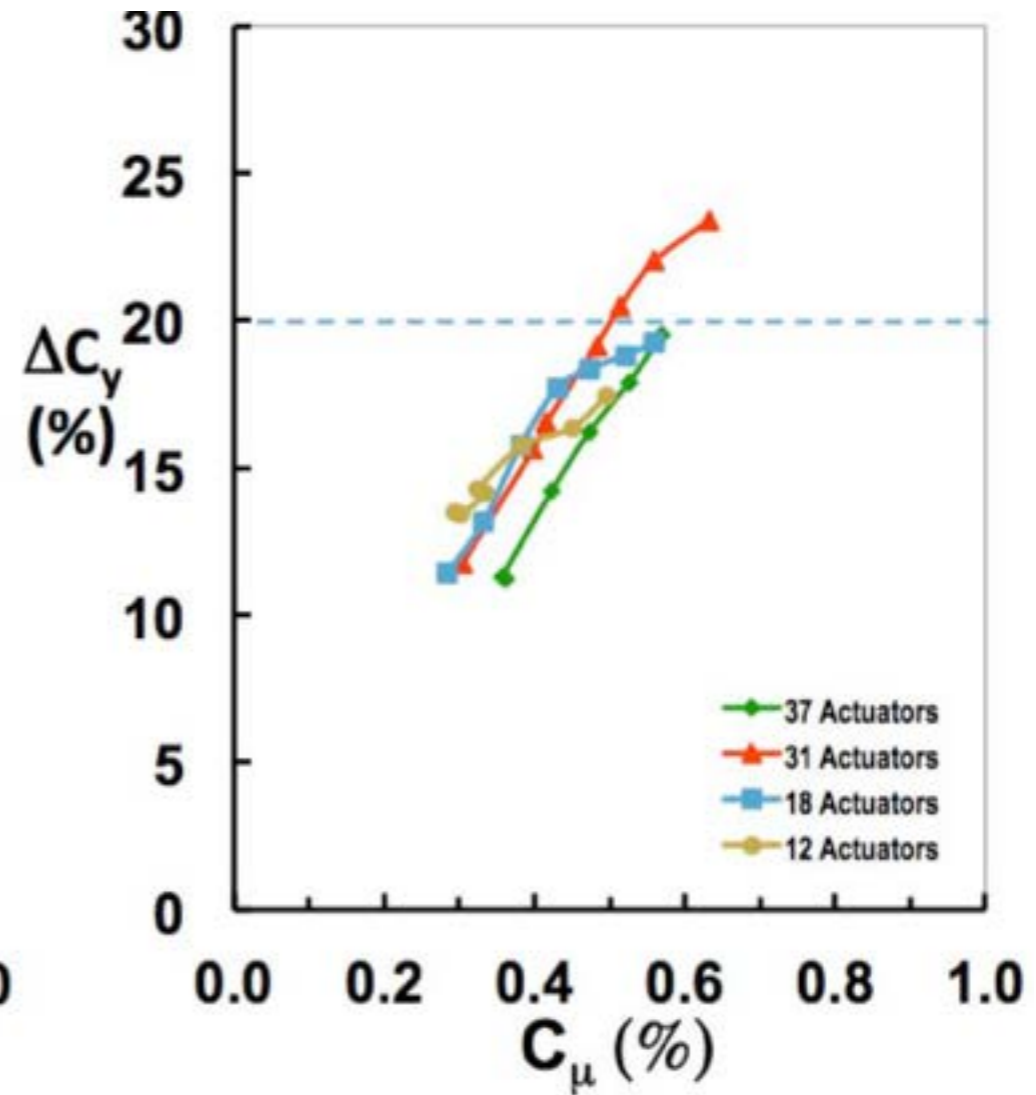


For full-scale wind tunnel testing, 37 fluidic oscillators are applied upstream of the rudder hinge.

20% increase in side force (more than twice what vortex generators give)



(a)  $\beta = 0^\circ$



(b)  $\beta = -7.5^\circ$

Whalen et al. (2015)

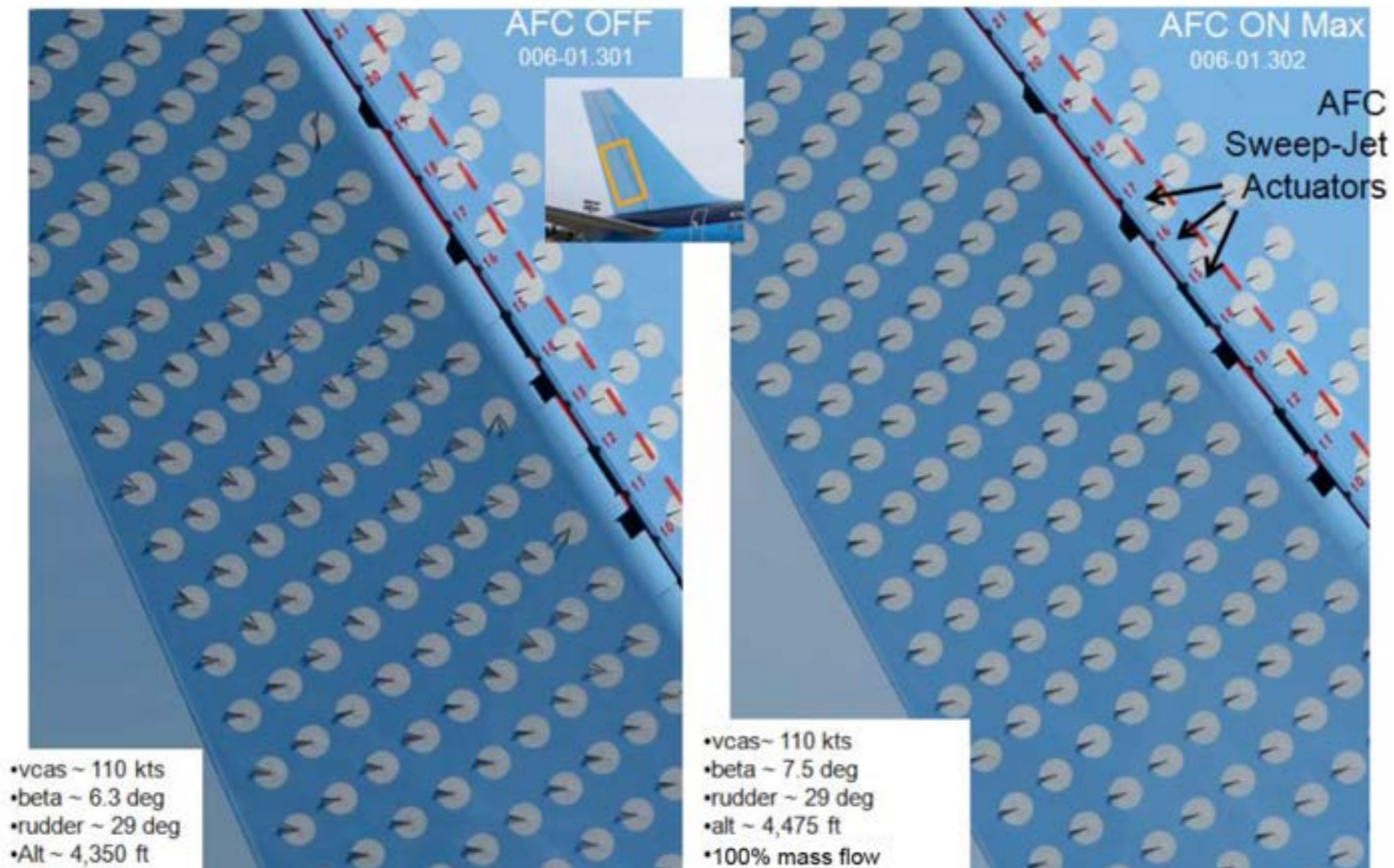


# NASA/Boeing ERA Control



For flight tests, 31 actuators used on the 757 ecoDemonstrator

13-16% increase in side force.



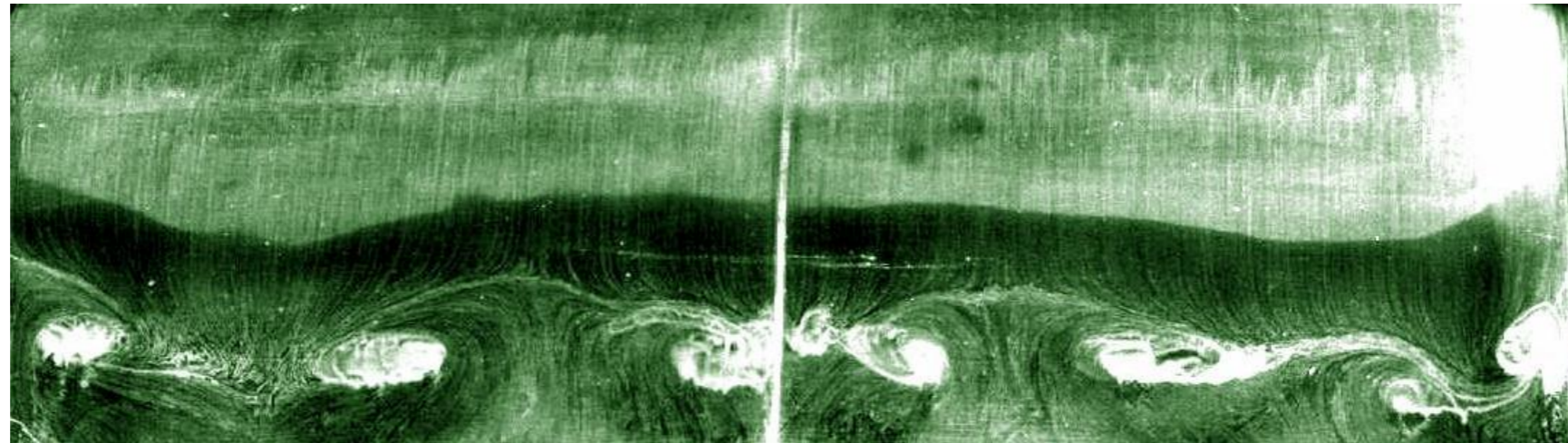
Lin et al. (2016)



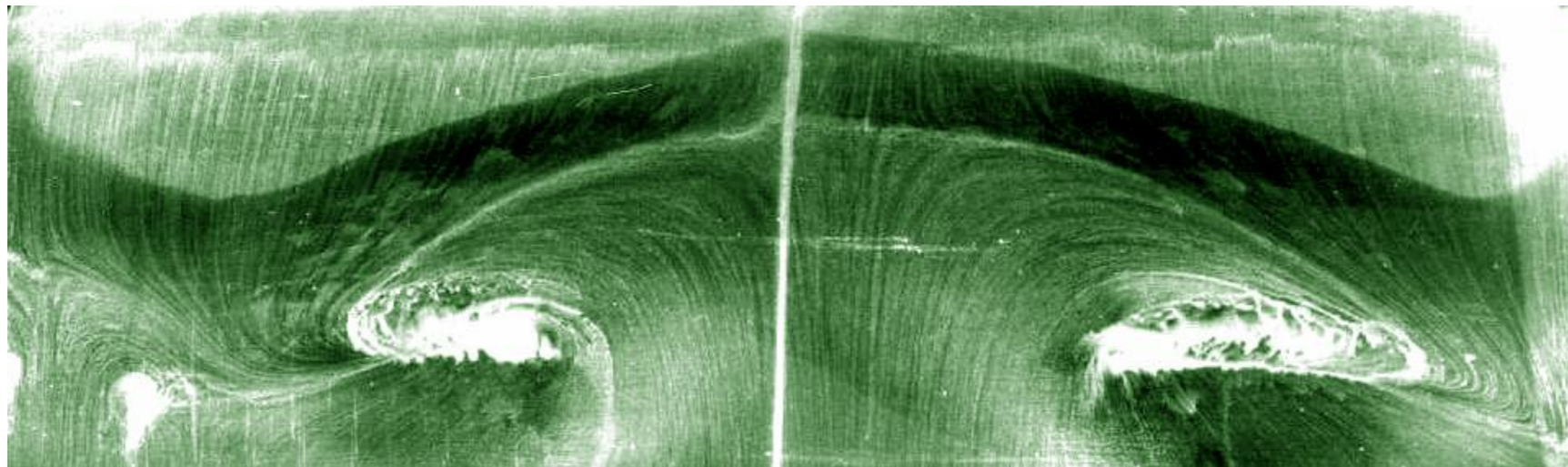
# 3D Nature of Separation

So far, separation is more or less discussed as a 2D phenomena, but it is not!

$$Re_c = 10^6$$
$$\alpha = 16^\circ$$

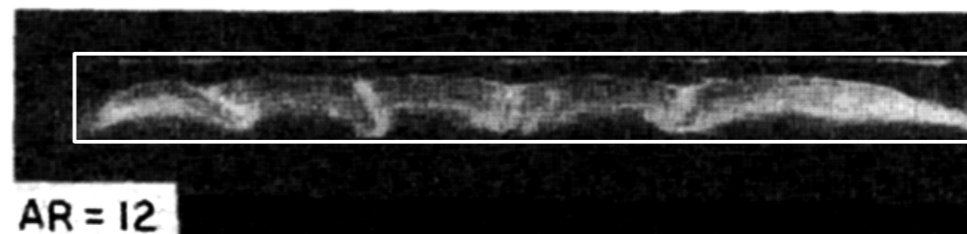
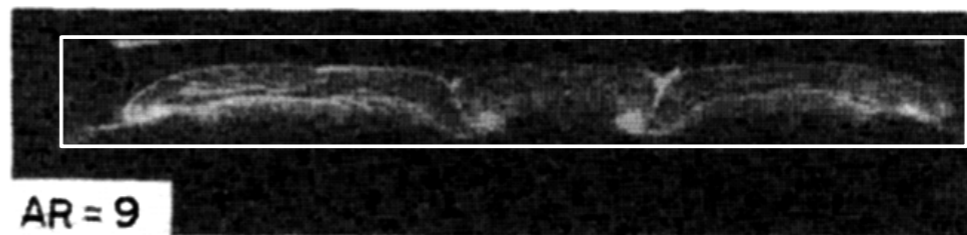
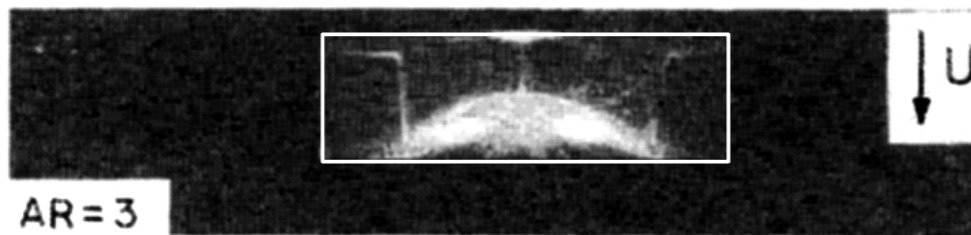
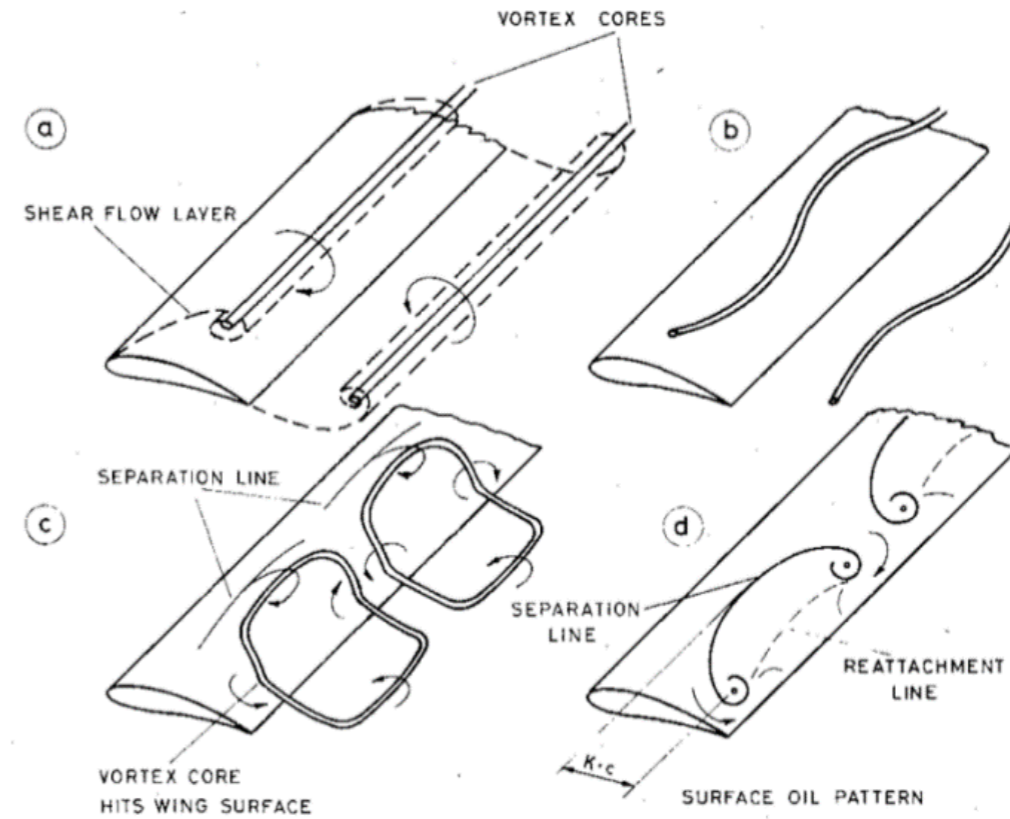
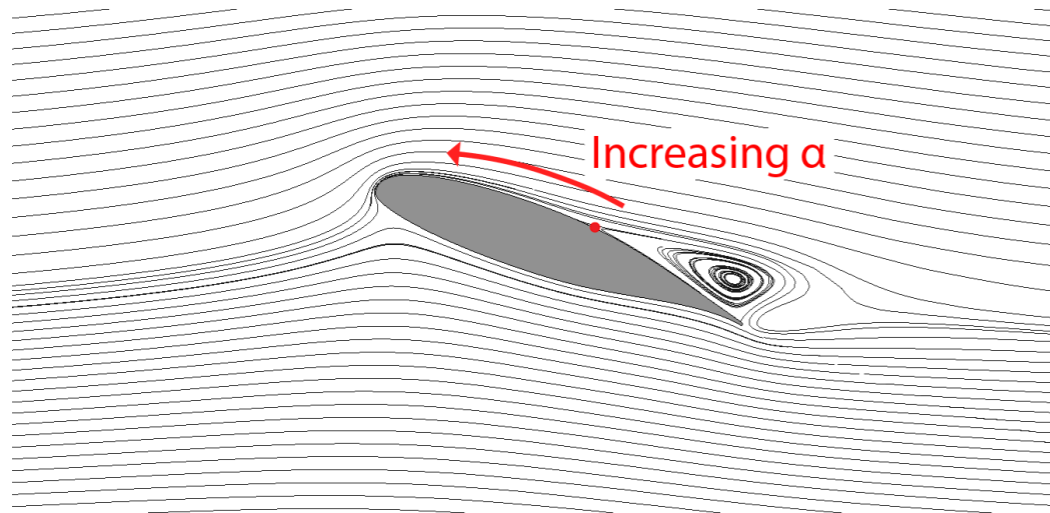


$$Re_c = 10^6$$
$$\alpha = 18^\circ$$



Images courtesy of JF Morrison (Imperial College)

# 3D Nature of Separation



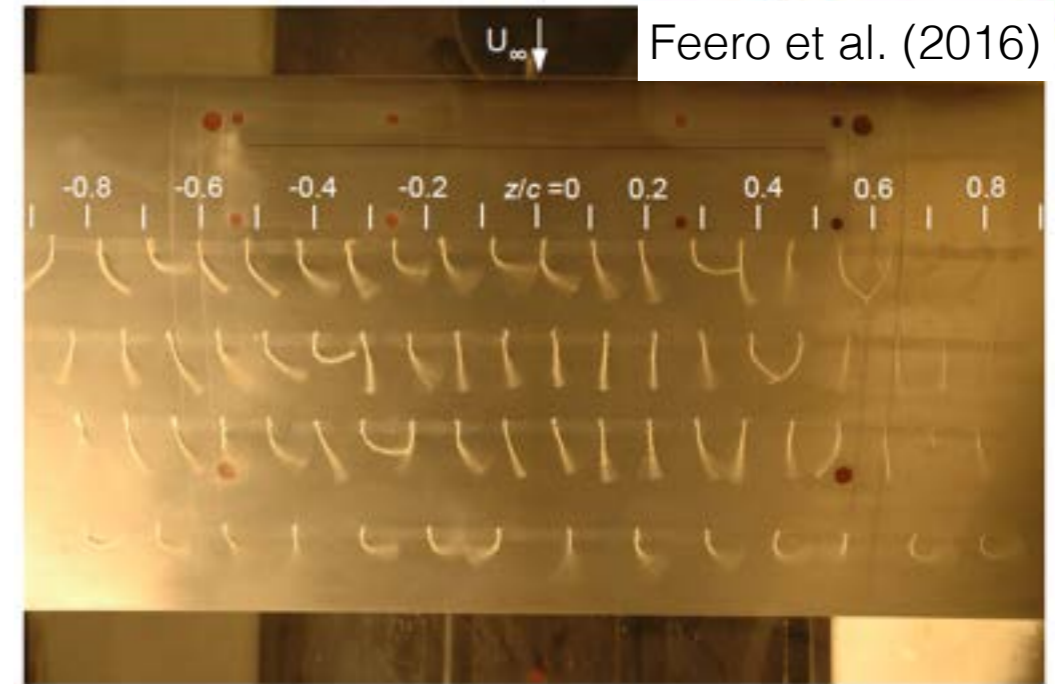
Winkelman & Barlow (1980)

# 3D Nature of Separation Control

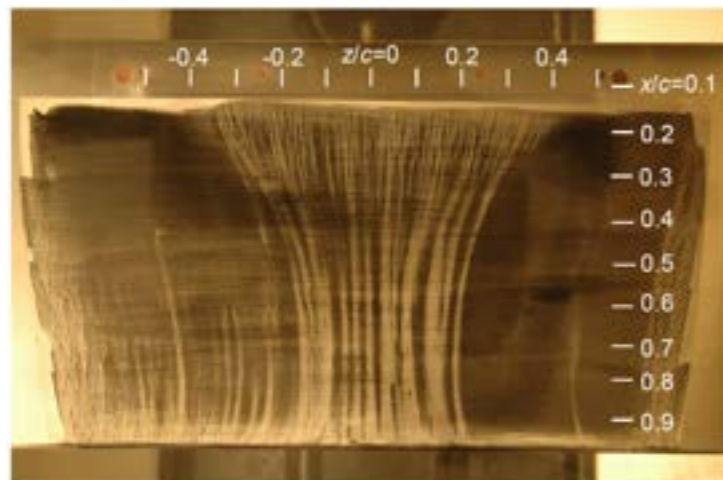


Nominally 2D actuator spanning 1/3 of model span.

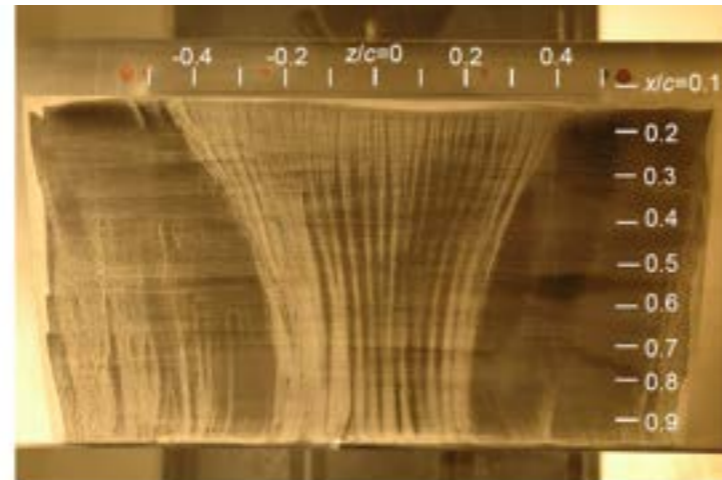
Significant 3D effect and rapid reduction of the reattached region downstream of the actuators.



(a) Baseline

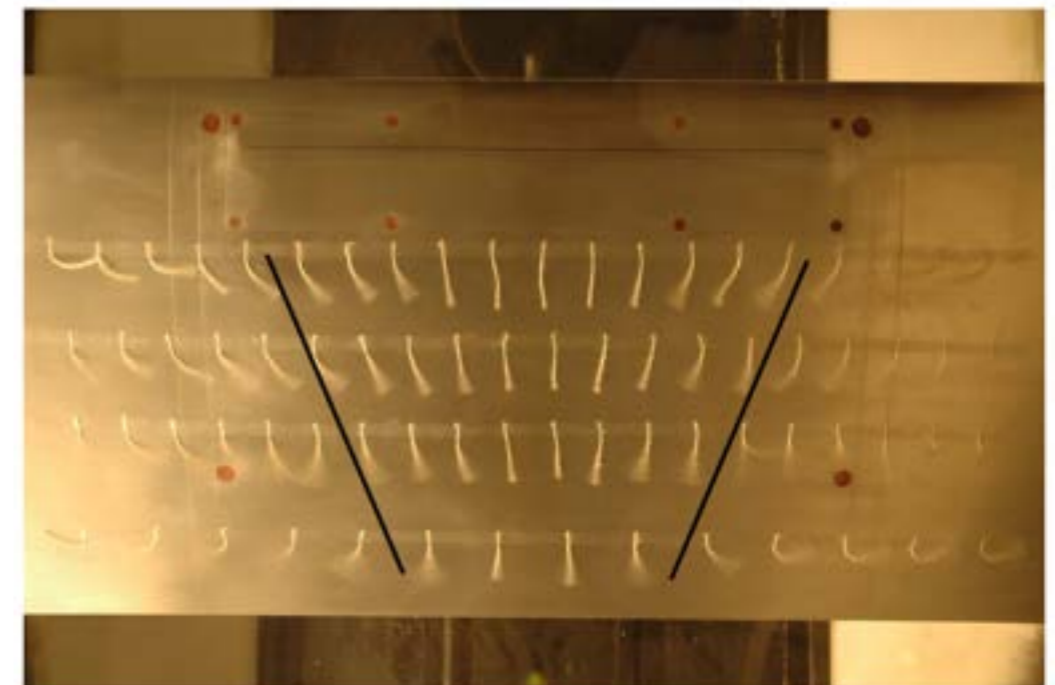


(a)  $C_B = 1$



(b)  $C_B = 2$

Region of effect increases with actuation amplitude.

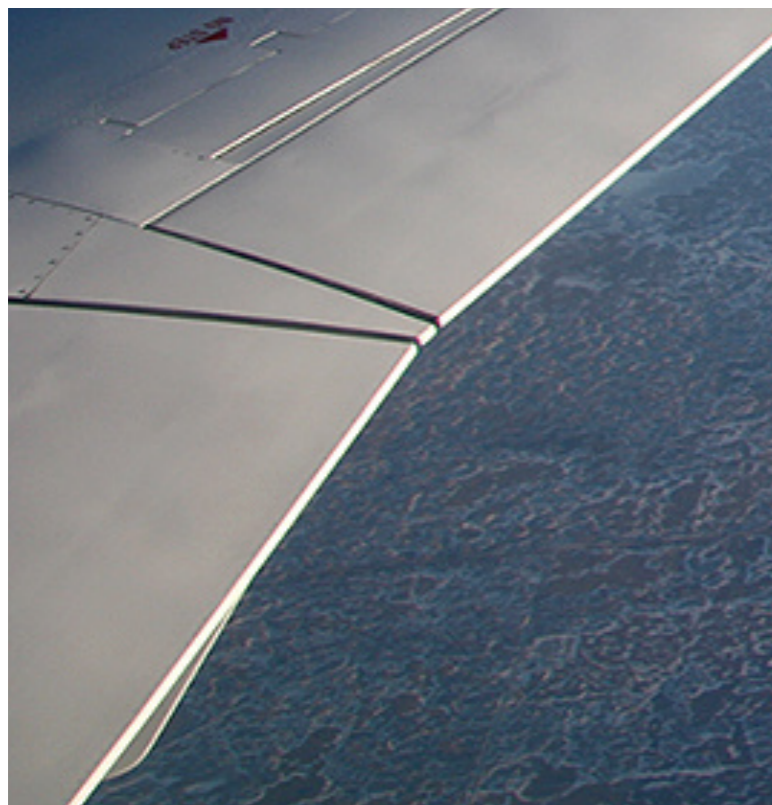


(c)  $F^+ = 1$

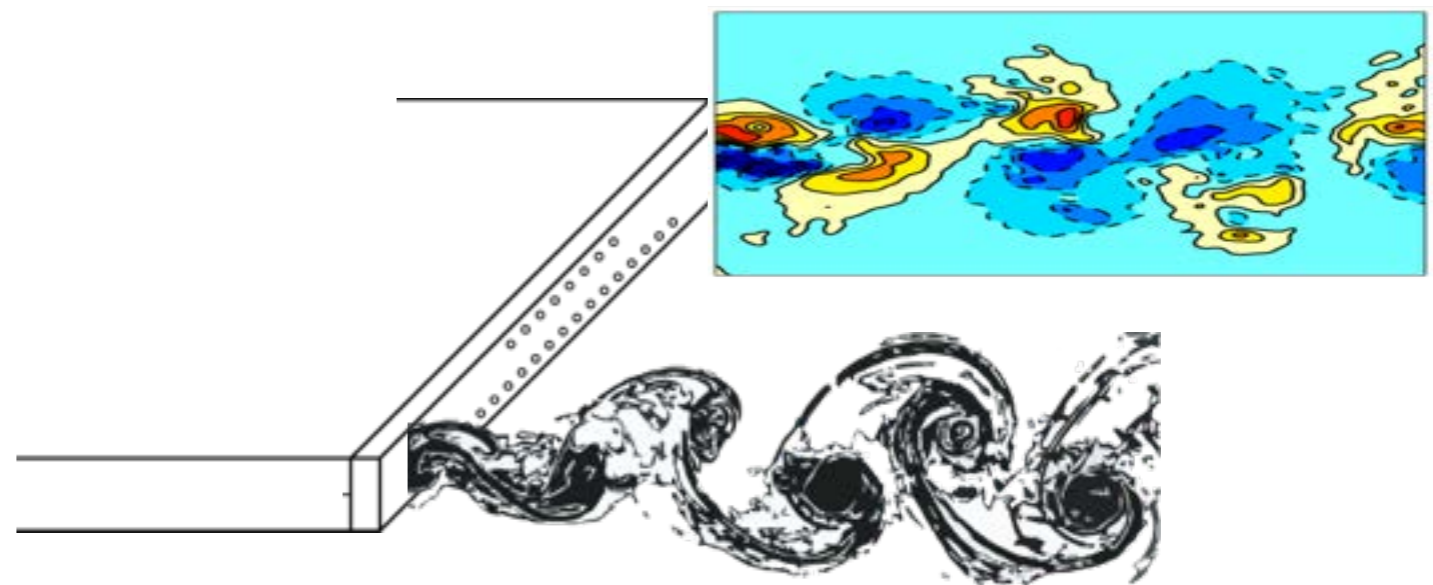
# Control of Bluff Bodies

For bluff bodies, the control strategy will be fundamentally different since the point of separation is either fixed, or very difficult to move.

We will focus primarily on the blunt trailing edge airfoil here, but will bring in some other canonical flow examples.



Bombardier CRJ trailing edge



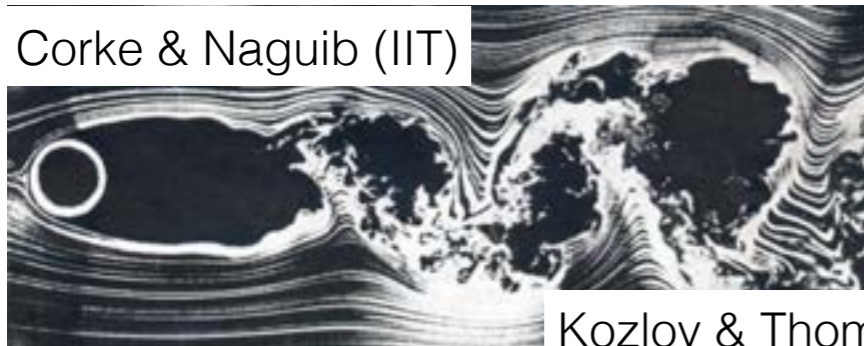
# Other Bluff Bodies

A few other bluff bodies of interest.

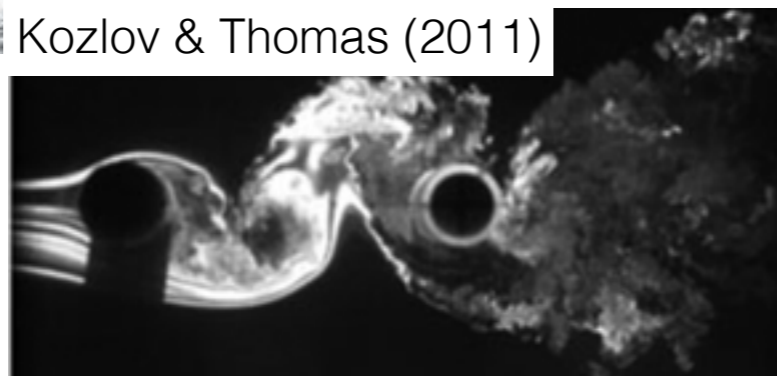
Landing gears involve a number of them:

- Struts (cylinders)
- Wheels (blocks and cavities)
- Landing gear bay (cavity)
- Wires and hoses (roughness)

Corke & Naguib (IIT)



Kozlov & Thomas (2011)



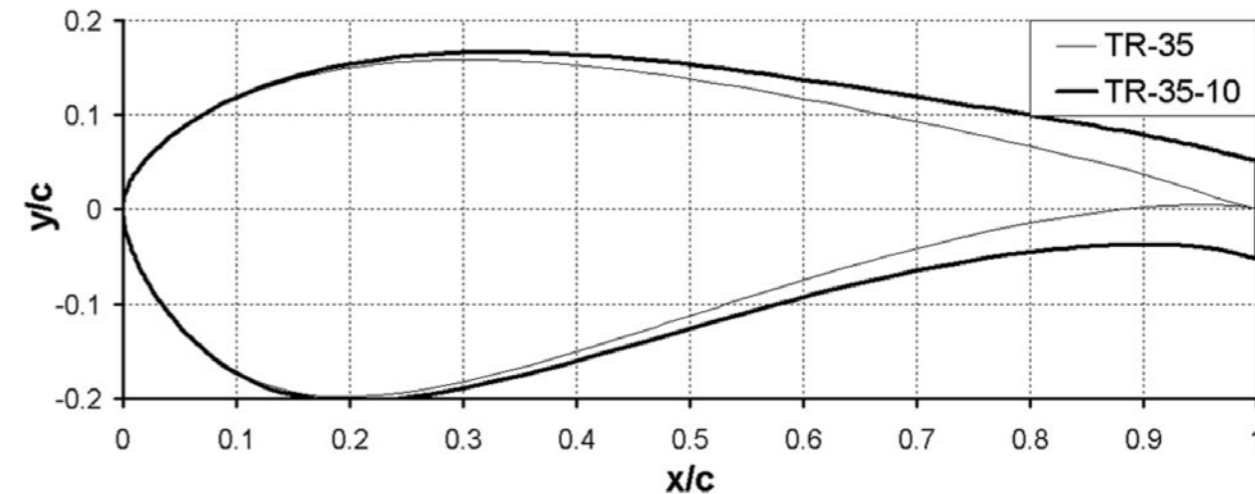
# Blunt Trailing Edge Airfoils



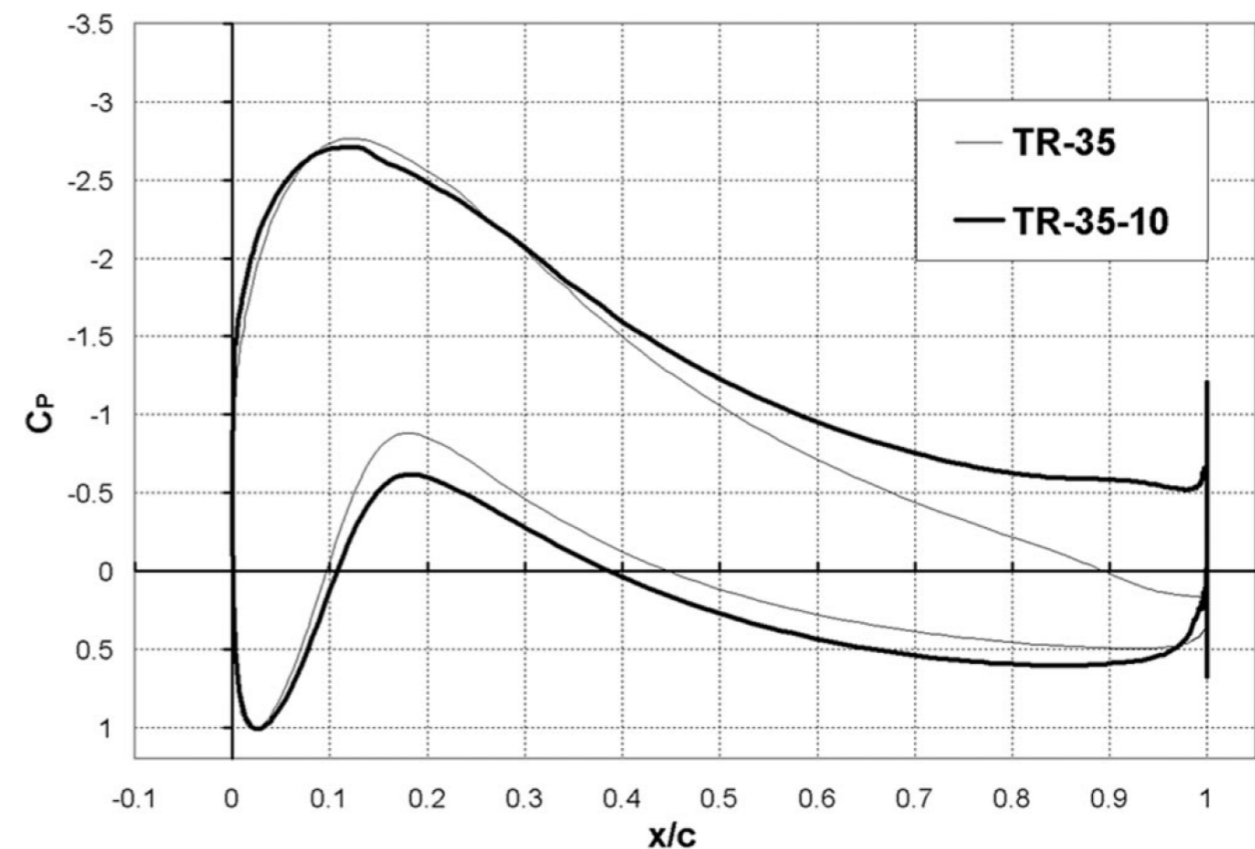
## Why blunt trailing edge?

- Increased sectional area and sectional moment of inertia for fixed airfoil maximum thickness
- Reduction of adverse pressure gradient on suction side
- Increased  $C_{L\alpha}$  and  $C_{Lmax}$ , decreased sensitivity to transition location
- Weaker shock - reduce wave drag

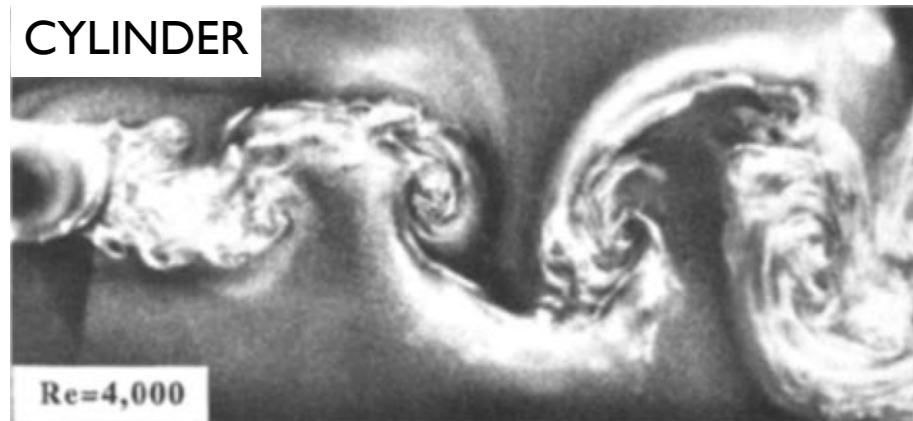
But, **increase in drag** and **unsteady loads** associated with **vortex shedding**, and **decrease in base pressure**



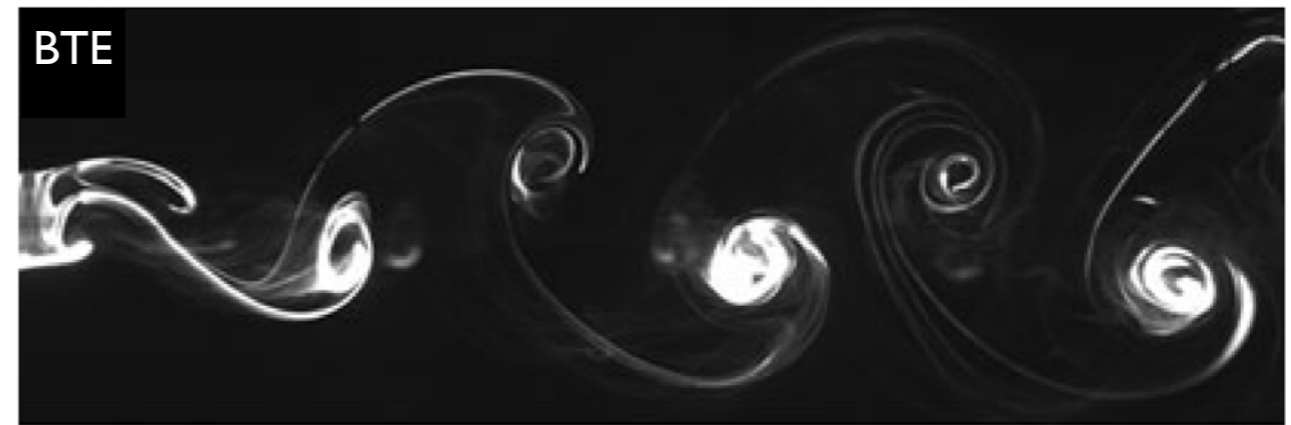
Figures from Standish & van Dam (2003)



# Bluff Body Flow (2D)

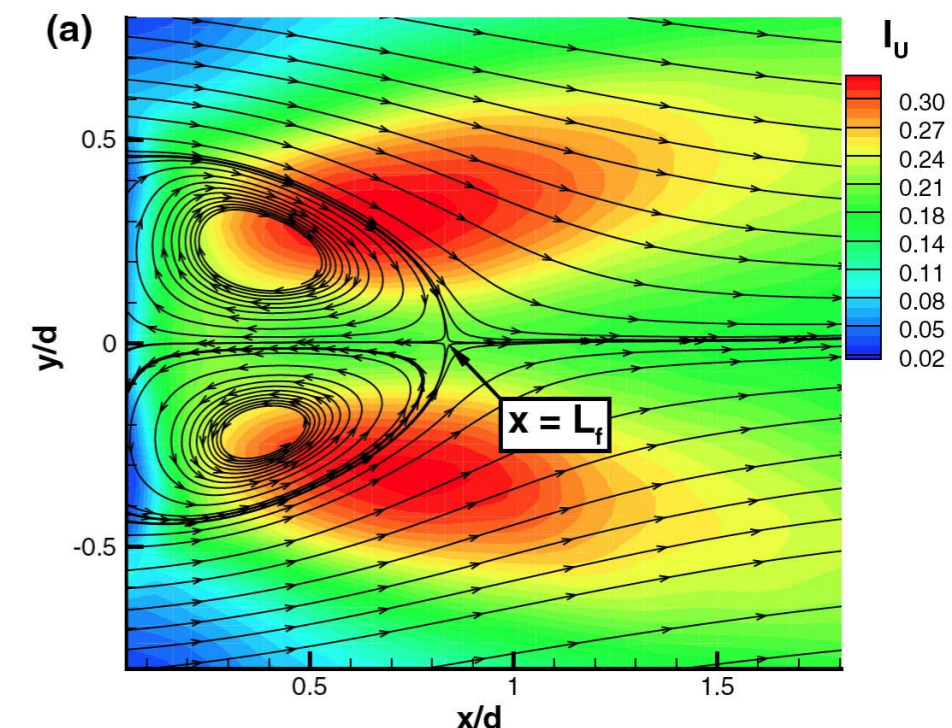


Williamson (1996)



Naghieb-Lahouti et al. (2012)

- **Vortex shedding** occurs above a threshold Reynolds number in wake of nominally two-dimensional bodies
  - Interacting shear layers: vortex grows until fluid with oppositely signed vorticity is entrained from across the wake (Gerrard, 1966)
- Vortex shedding leads to **unsteady aerodynamic forces, vibrations and noise**
- A convenient and simple way to characterize the strength of the vortex shedding is the **vortex formation length,  $L_f$** :
  - streamwise distance at which rms velocity fluctuations are maximum



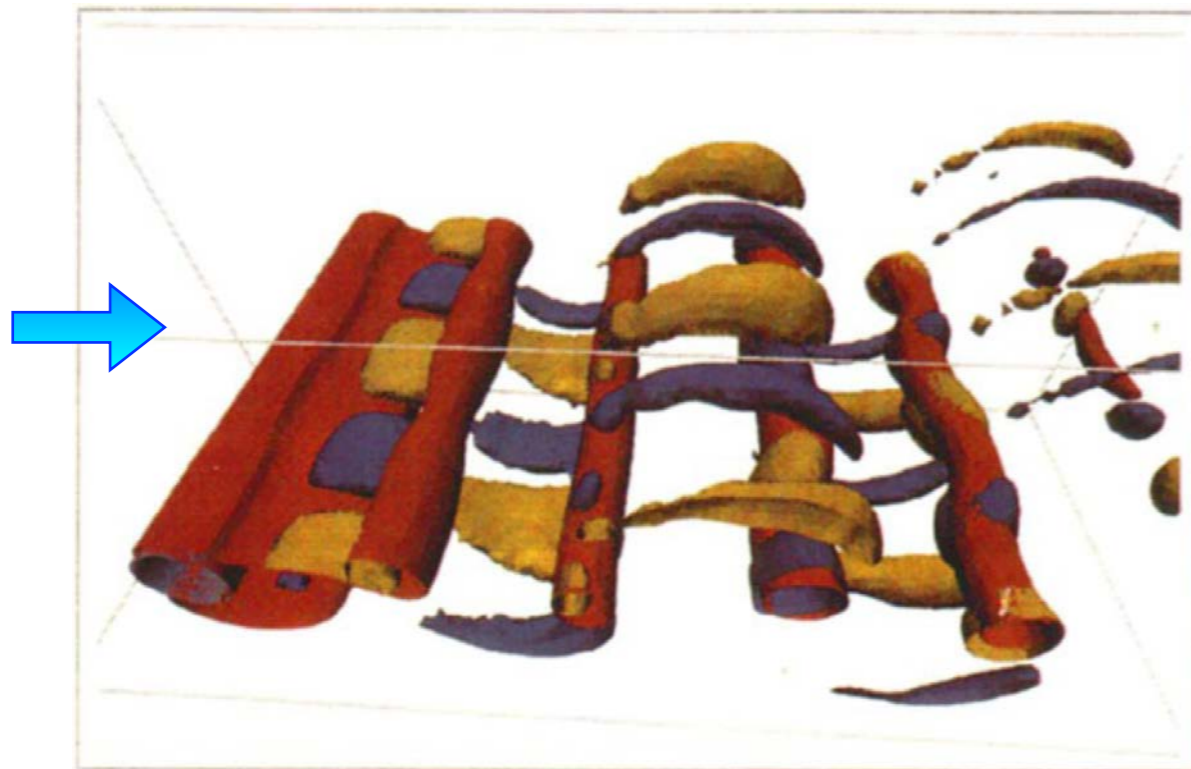
Figures from Naguib-Lahouti, Lavoie & Hangan (2014)



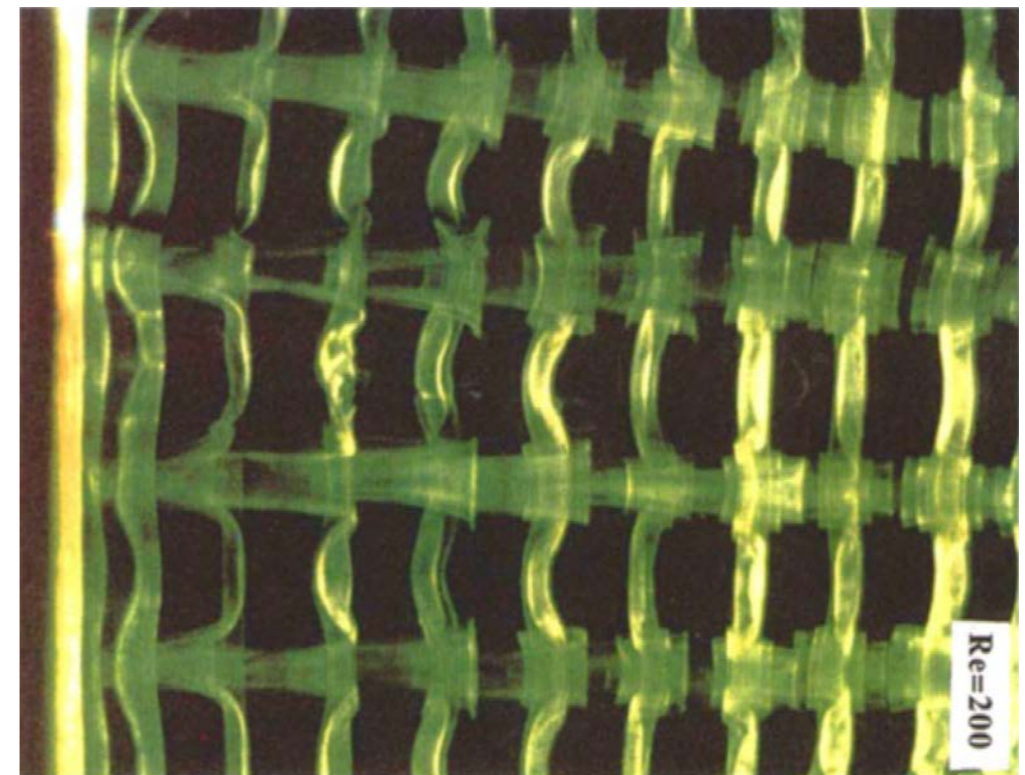
# Secondary Instabilities

- Primary vortices are connected by **pairs of streamwise vortices** with characteristic spacing,  $\lambda_z$
- Unstable mode topologies **depend on body geometry and Reynolds number**
- Cylinder wake
  - Williamson (1996): critical  $Re \approx 180-190$ ,  $\lambda_z \approx 3-5d$  (Mode B)

numerical



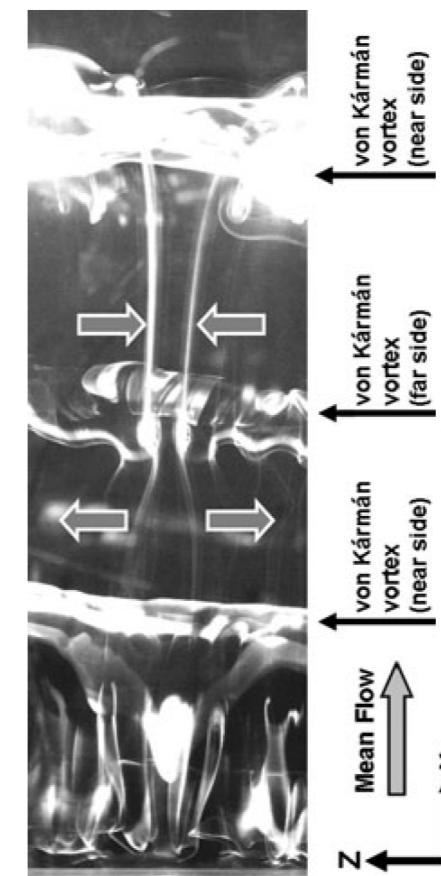
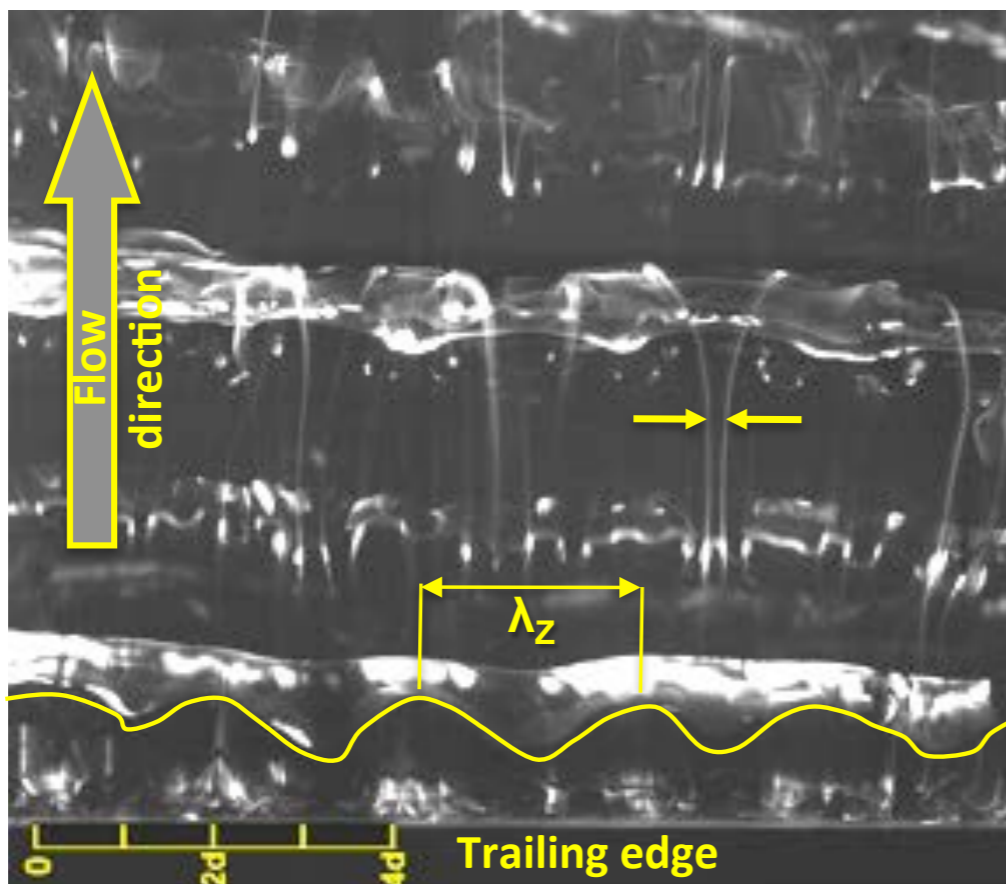
experiment



# Secondary Instabilities



- Primary vortices are connected by **pairs of streamwise vortices** with characteristic spacing,  $\lambda_z$
- Unstable mode topologies **depend on body geometry and Reynolds number**
- For BTE geometry with  $c/d = 12.5$ , Mode B'
  - Ryan et al. (2005), numerical stability analysis: critical  $Re_d \approx 410$ ,  $\lambda_z \approx 2.2d$
  - Naghib-Lahouti et al. (2012, 2014), experiments:  $\lambda_z \approx 2.3 - 2.5$  for  $Re_d = 2,000 - 30,000$



Figures from Naghib-Lahouti et al. (2012)

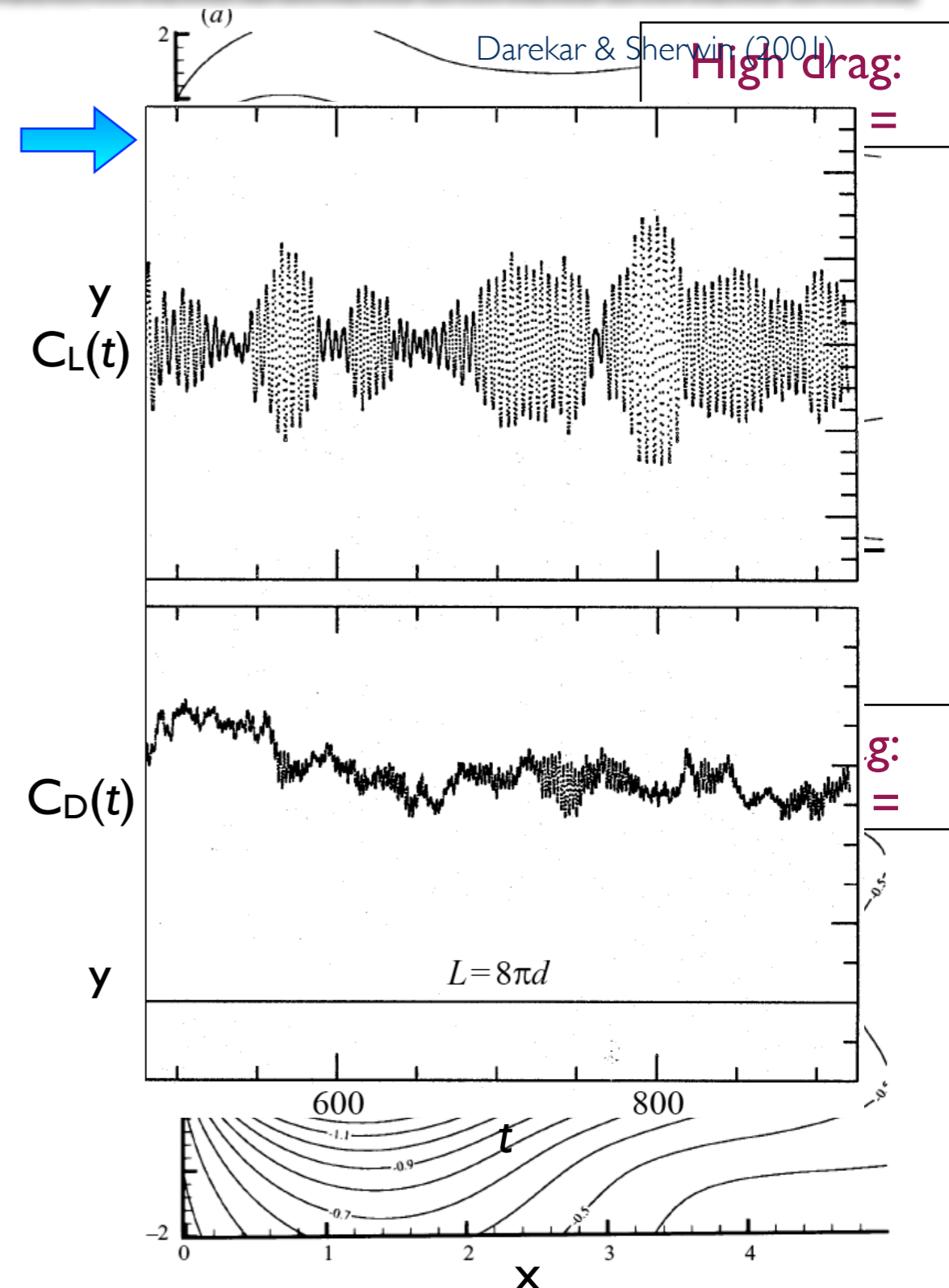
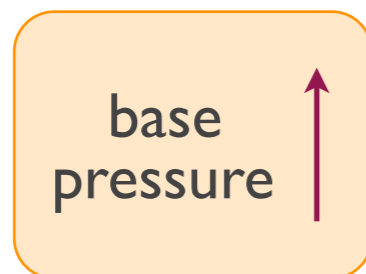
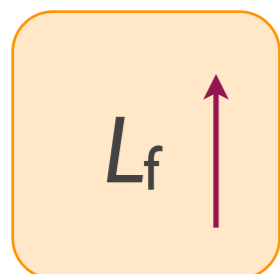


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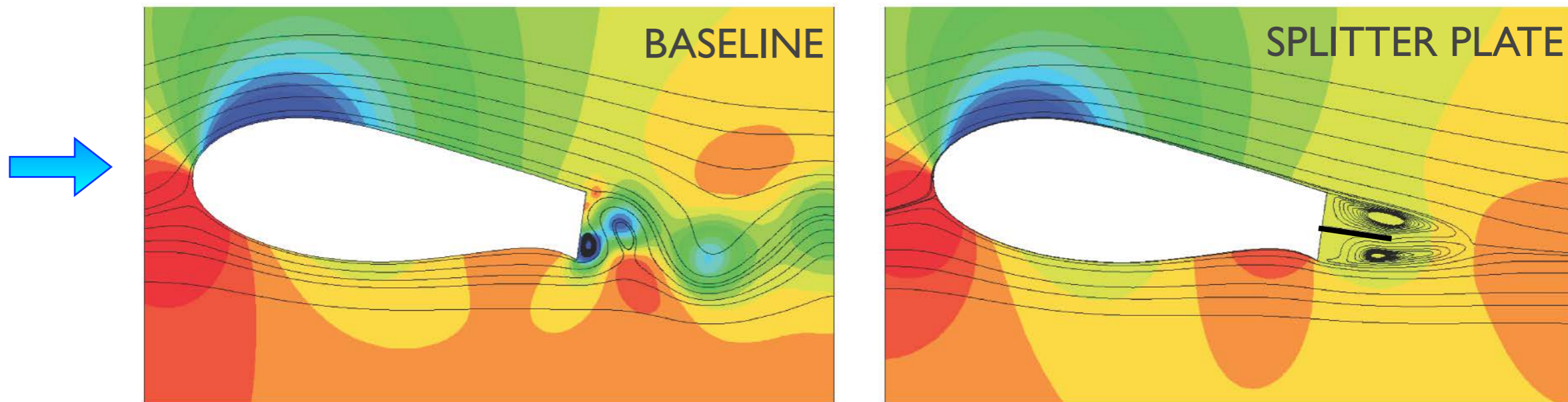
# Three-dimensional Effects

- Low-frequency modulation of local measurements and global properties
- Henderson (1997), cylinder flow: modulation of lift from “phase dislocations,” occurs in “bursts at irregular intervals”
- Darekar & Sherwin (2001), wavy square cylinder: **flow switches between two states**, mildly or highly 3D
- Najjar & Balachandar (1998), normal flat plate: two states correspond to **different  $L_f$  and  $C_p$**
- **Increasing wake 3-dimensionality** ↑



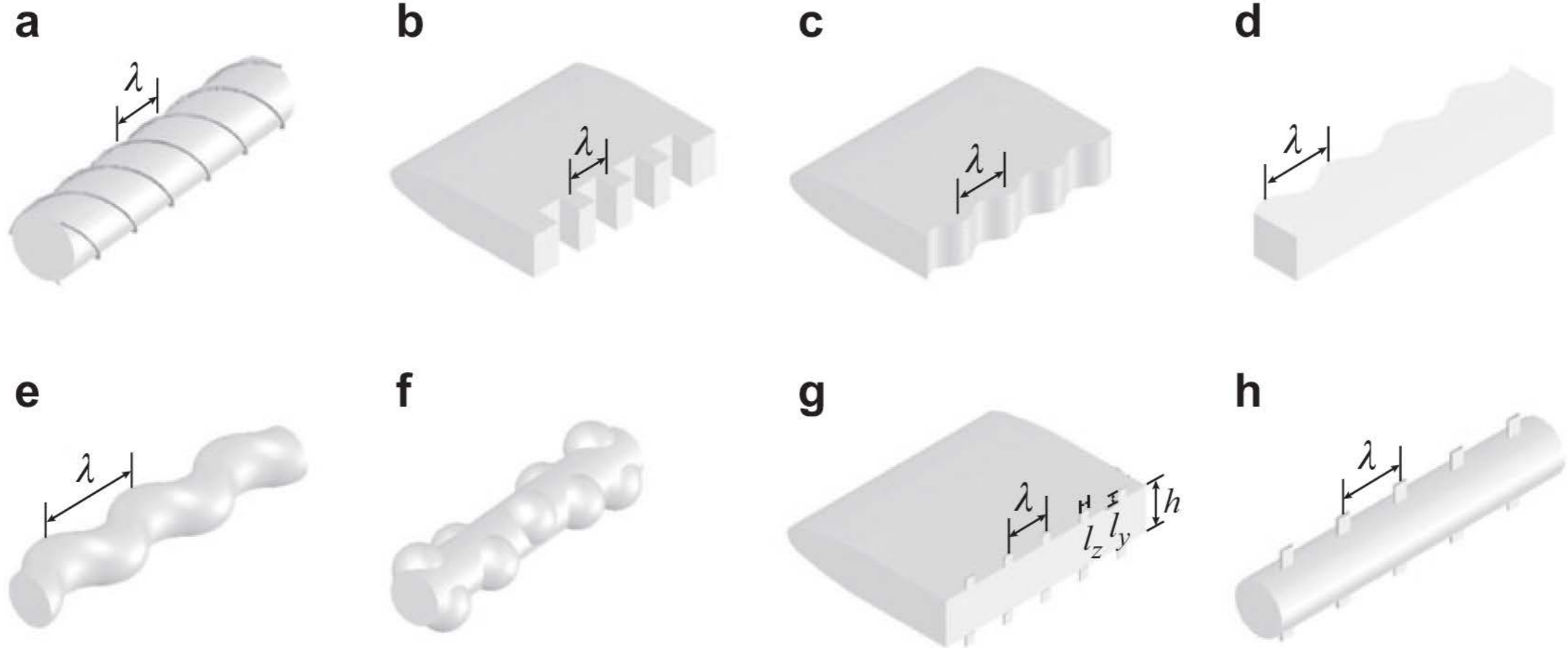
# Control of Bluff Body Wake

- Considerations for wake modification: drag reduction, modify acoustic signature, reduction of unsteady structural loading
- Baker & van Dam (2008): experiments and simulation of BTE airfoil
  - Passive, 2-dimensional forcing with splitter plate to increase  $L_f$

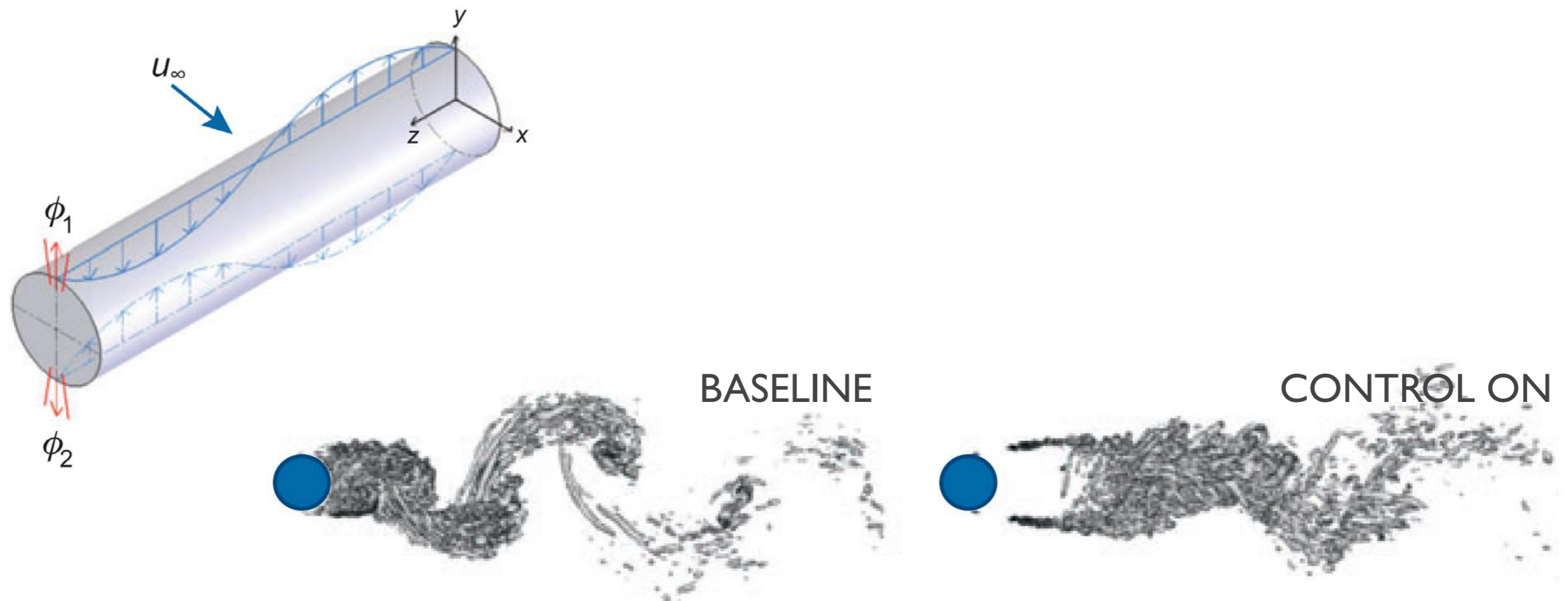


# Control of Bluff Body Wake

- Considerations for wake modification: drag reduction, modify acoustic signature, reduction of unsteady structural loading
- Choi et al. (2008): reviewed control methods for bluff bodies
  - Passive, 3-dimensional forcing to leverage intrinsic secondary instability

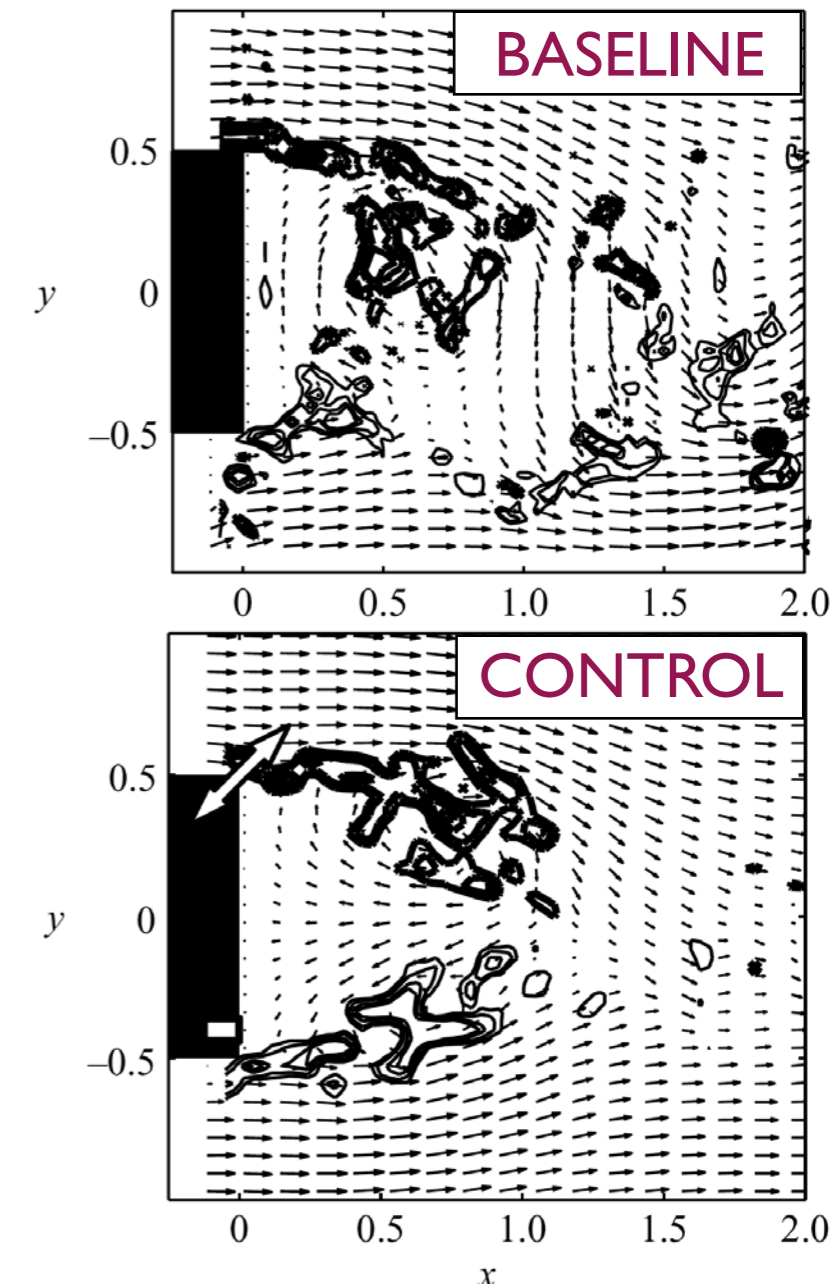
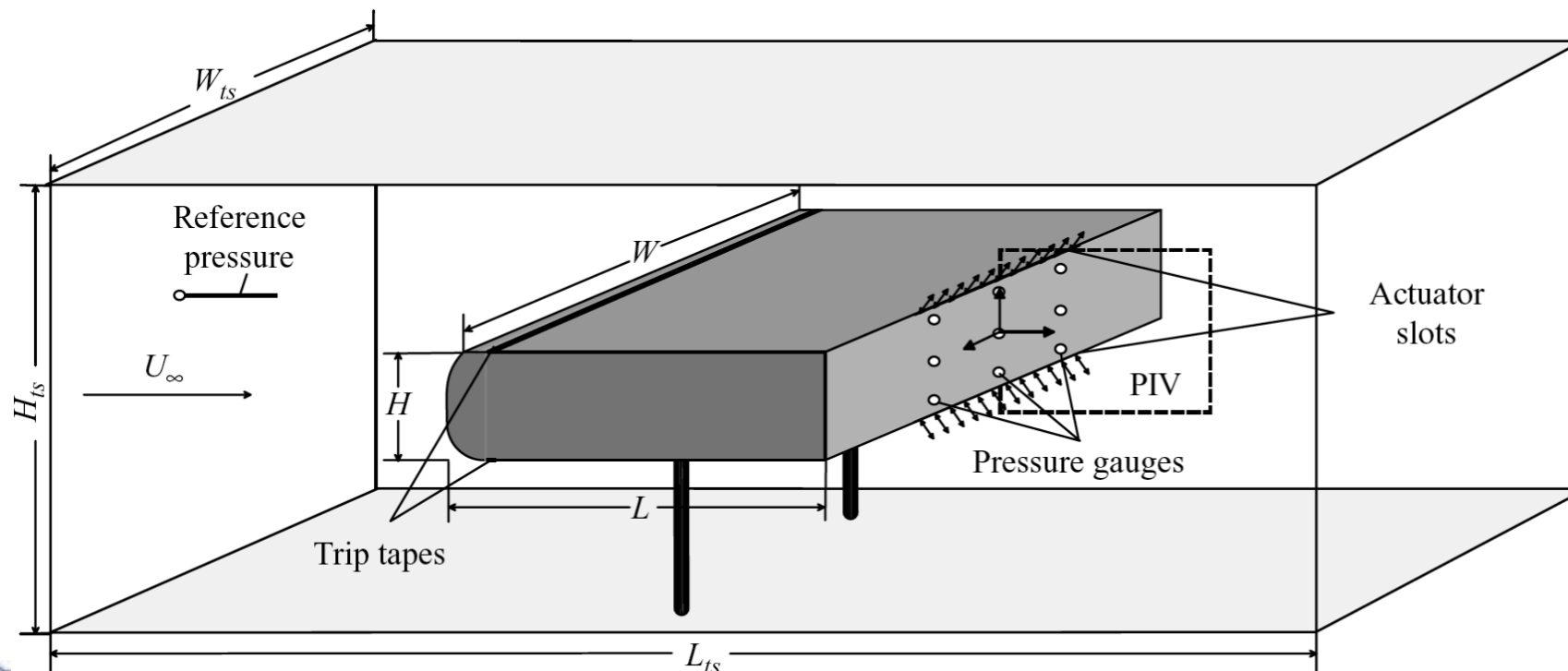


- Considerations for wake modification: drag reduction, modify acoustic signature, reduction of unsteady structural loading
- Kim & Choi (2005): open loop, spatially sinusoidal suction and blowing
  - Attenuation or annihilation of vortices reduced mean and fluctuating drag



# Closed-loop Control

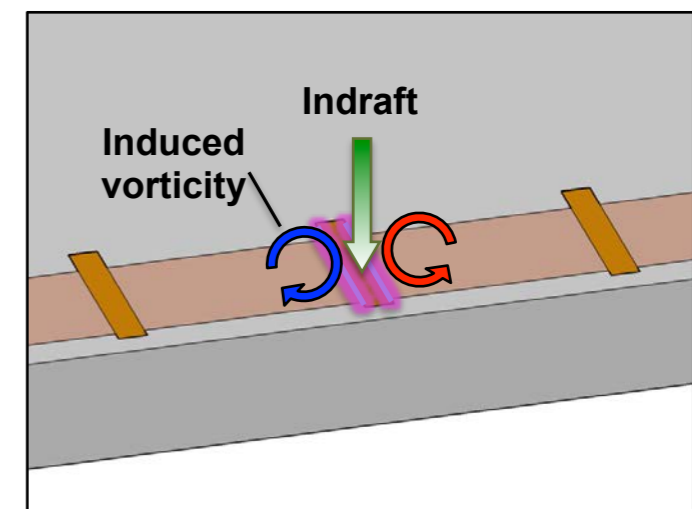
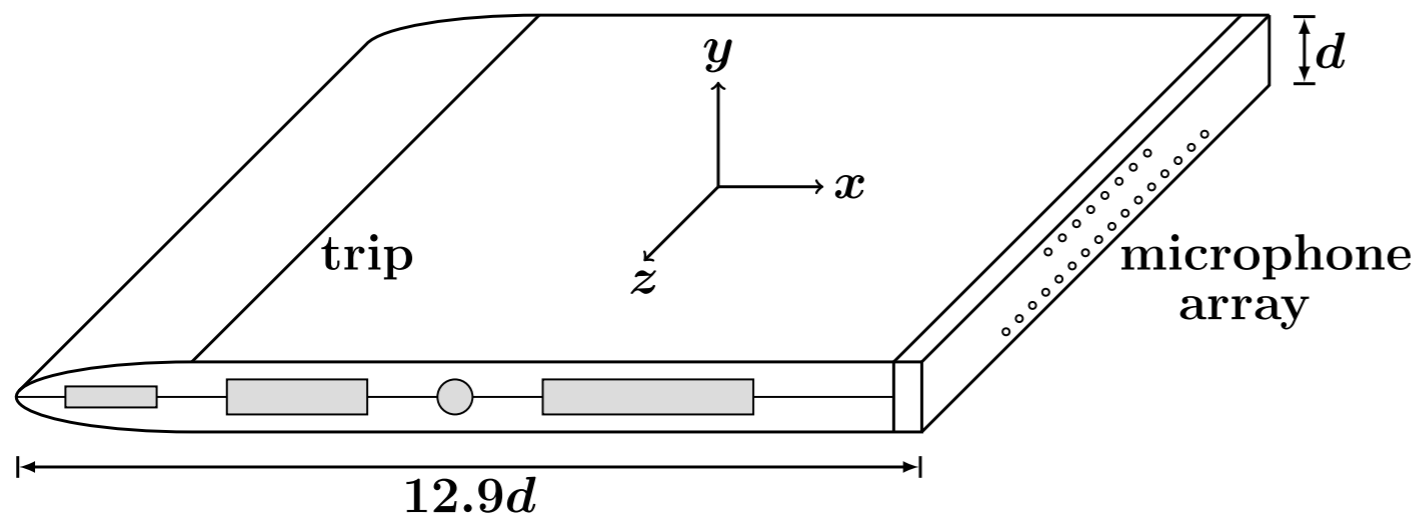
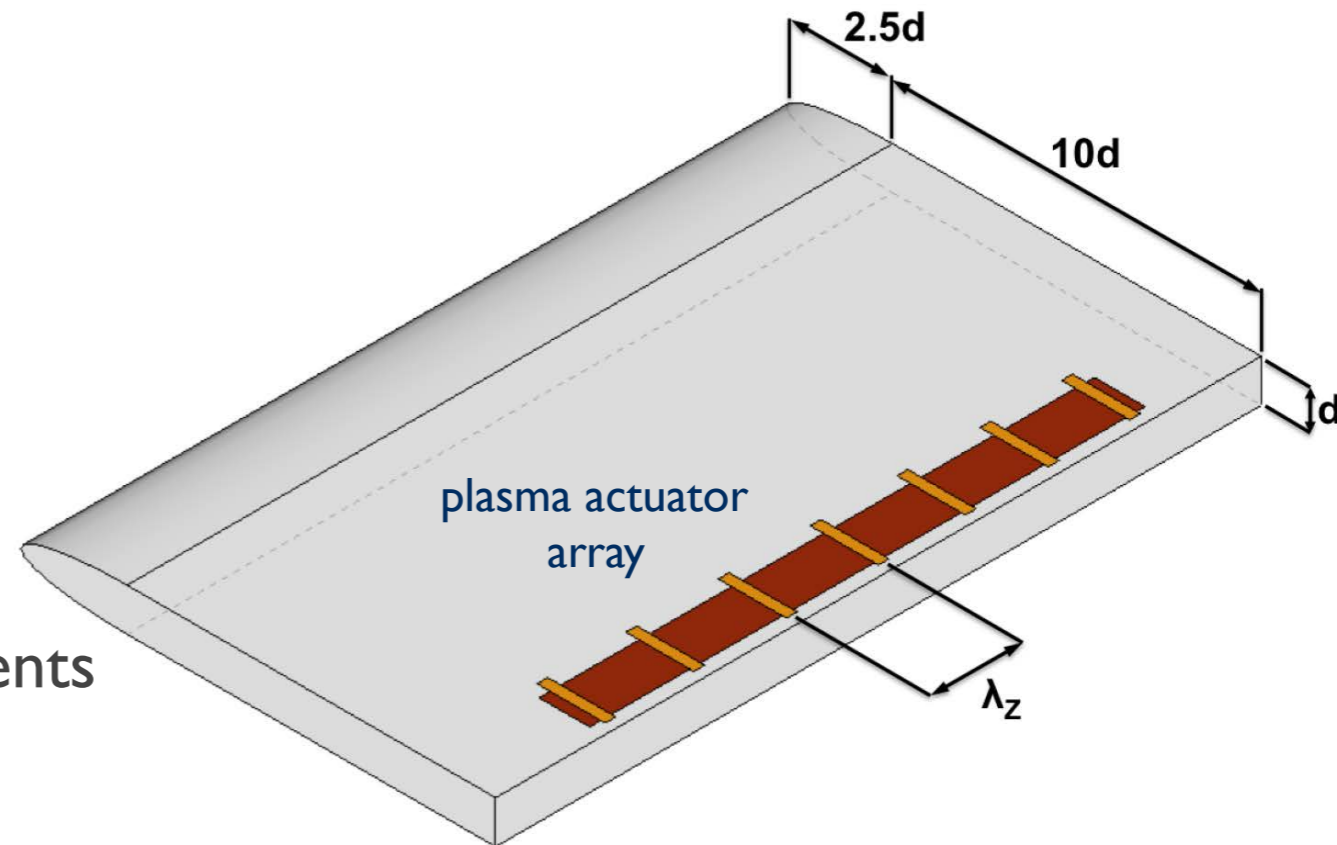
- Considerations for wake modification: drag reduction, modify acoustic signature, reduction of unsteady structural loading
- Pastoor et al. (2008): closed-loop, spanwise uniform actuation
  - Disrupted alternating shedding to increase  $L_f$  and decrease drag by 15%
  - Actuation on one side only and over half span of model



# Distributed Forcing via Plasma

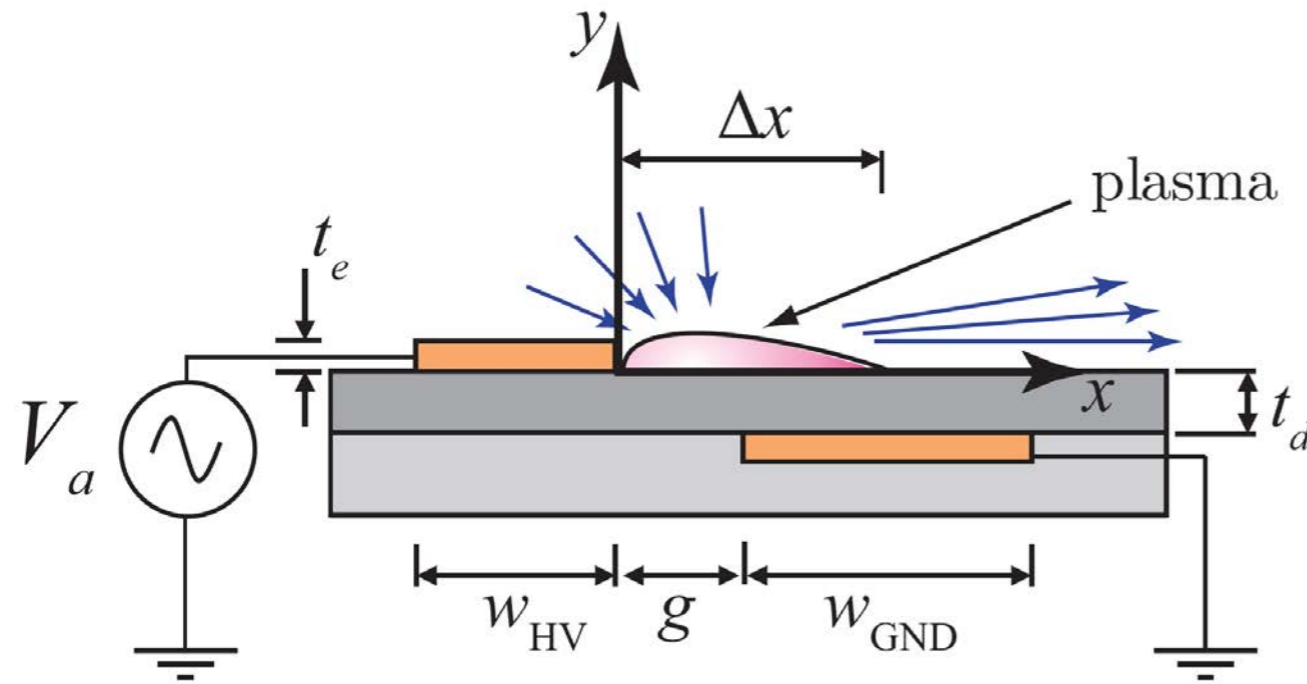


- Focus on **simplified, canonical geometry**
- Elliptical leading edge, flat plate section of thickness  $d$  with blunt trailing edge
- Chord =  $12.5d$
- Additional model features:
  - Active flow control via **plasma actuators**
  - Time-resolved base pressure measurements via **microphone array** extension



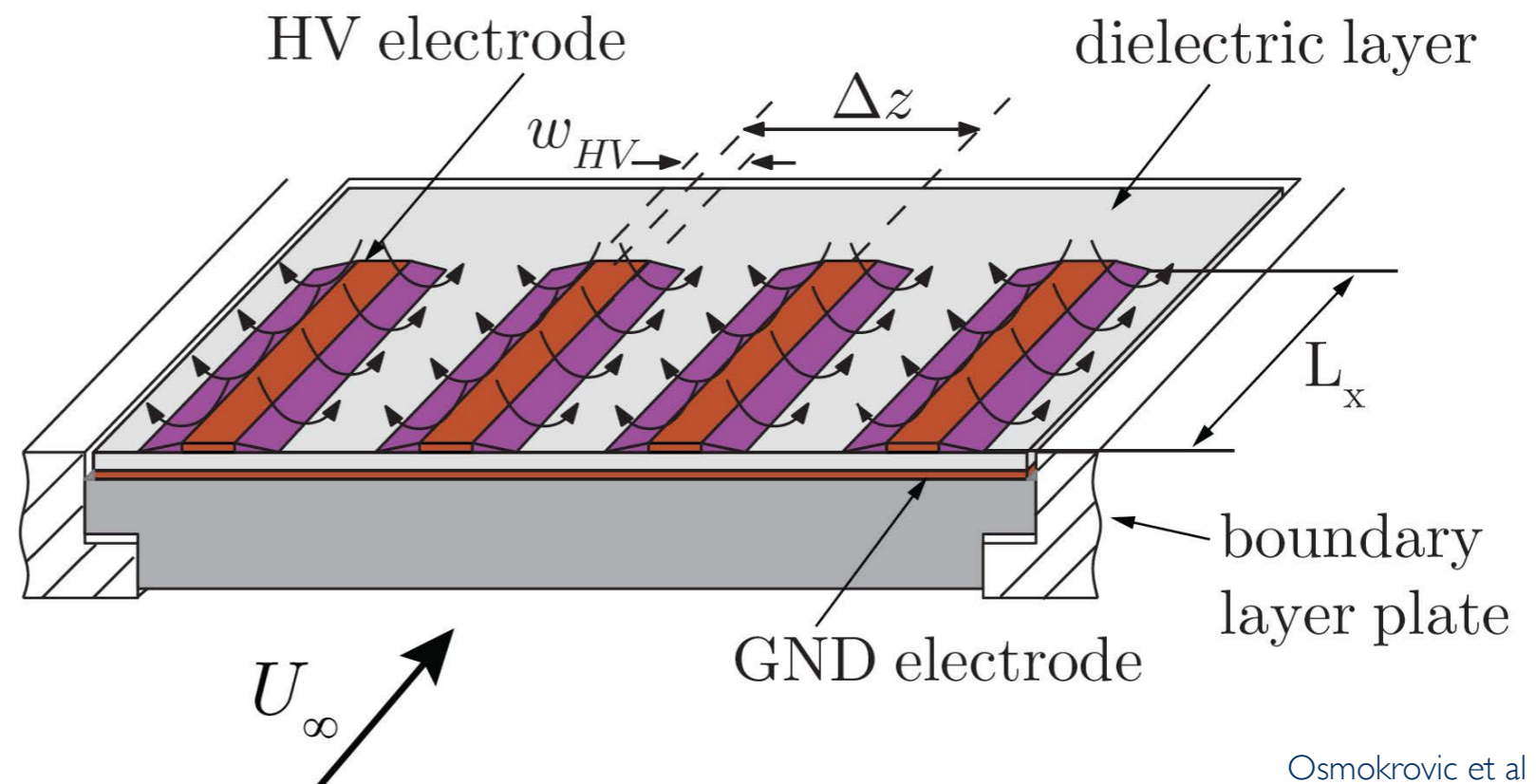


# Plasma Actuators



- Electrical input → fluidic output
- Single actuator produces wall jet

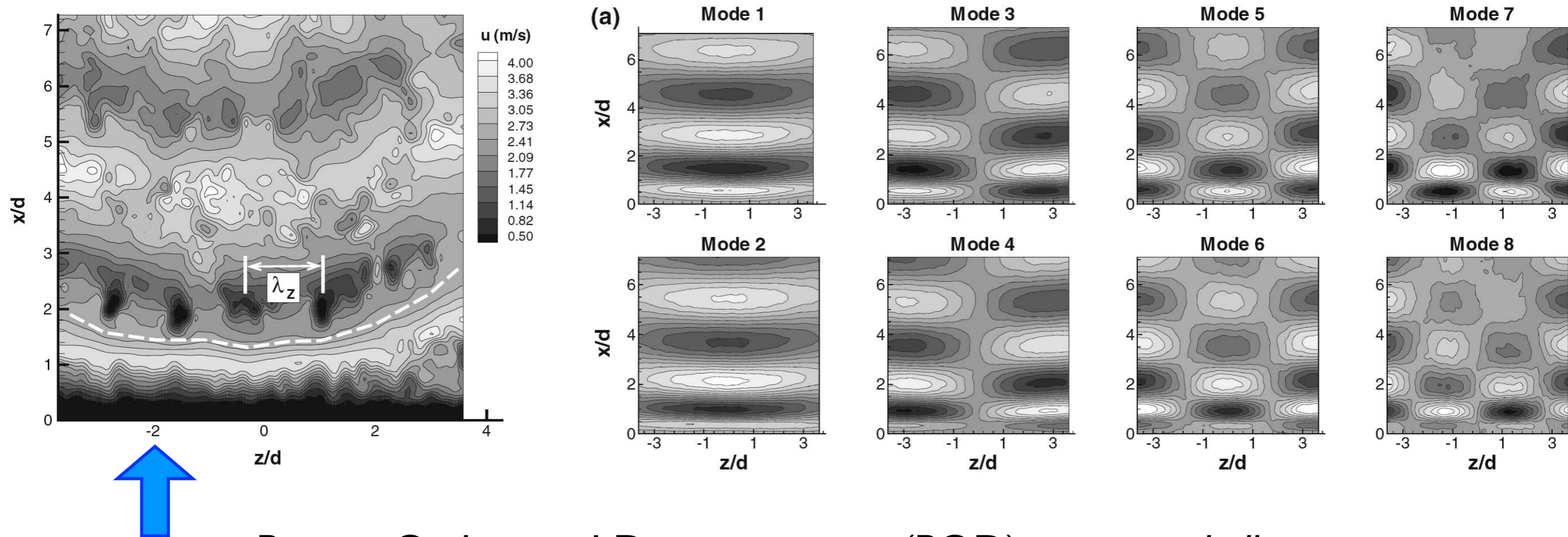
- Actuator array introduces alternating streak disturbance



# POD for Structure Identification



- Naghib-Lahouti et al. (2012, 2014) used POD model to extract dominant structures in wake



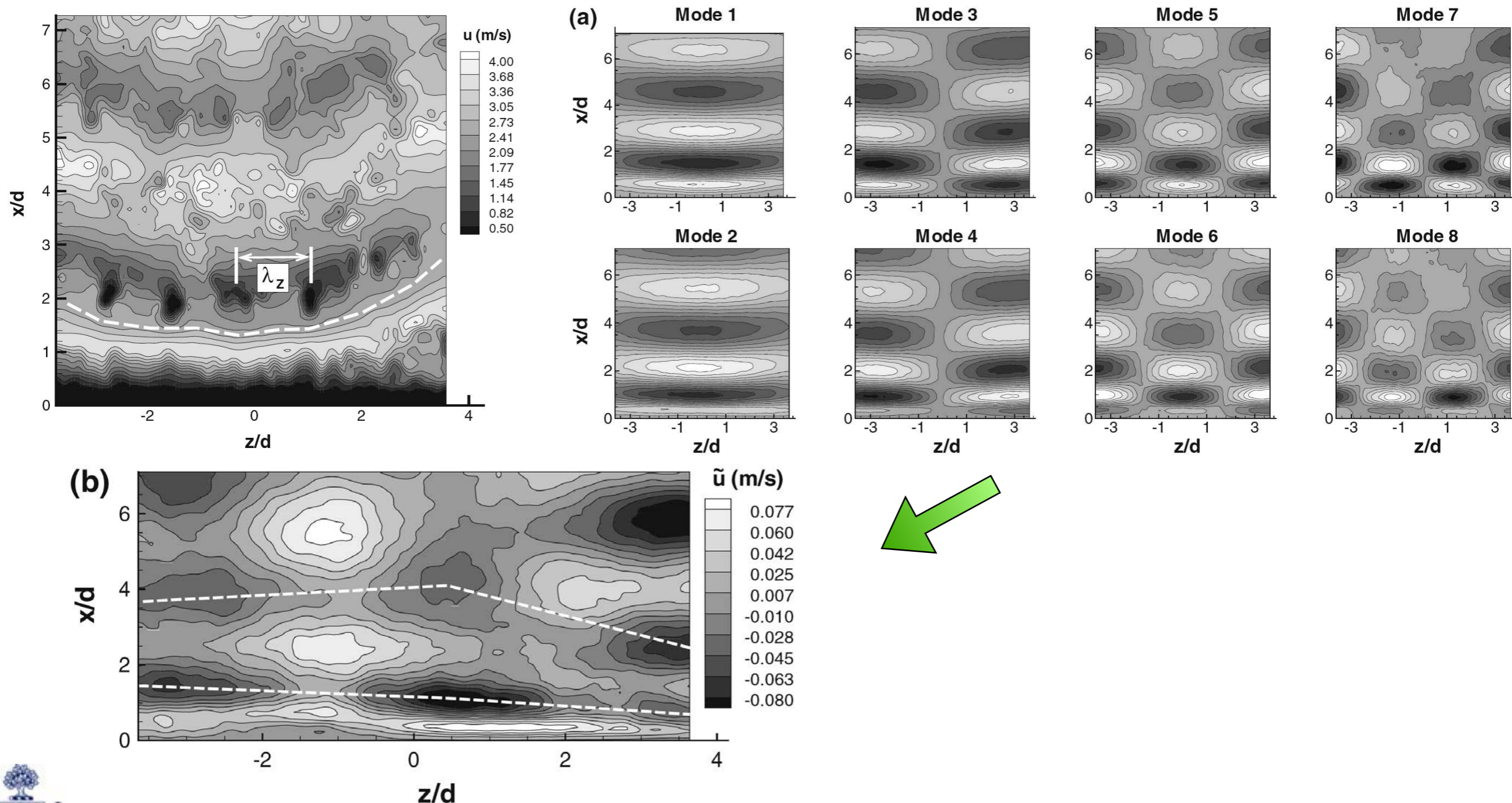
Proper Orthogonal Decomposition (POD) - *in a nutshell*

$$\tilde{u}_i(x_i, t) = \sum_{n=1}^M \underbrace{a^n(t)}_{\text{time varying coefficient}} \underbrace{\phi^n(x_i)}_{\text{orthogonal eigenfunctions (modes)}}$$

# POD for Structure Identification



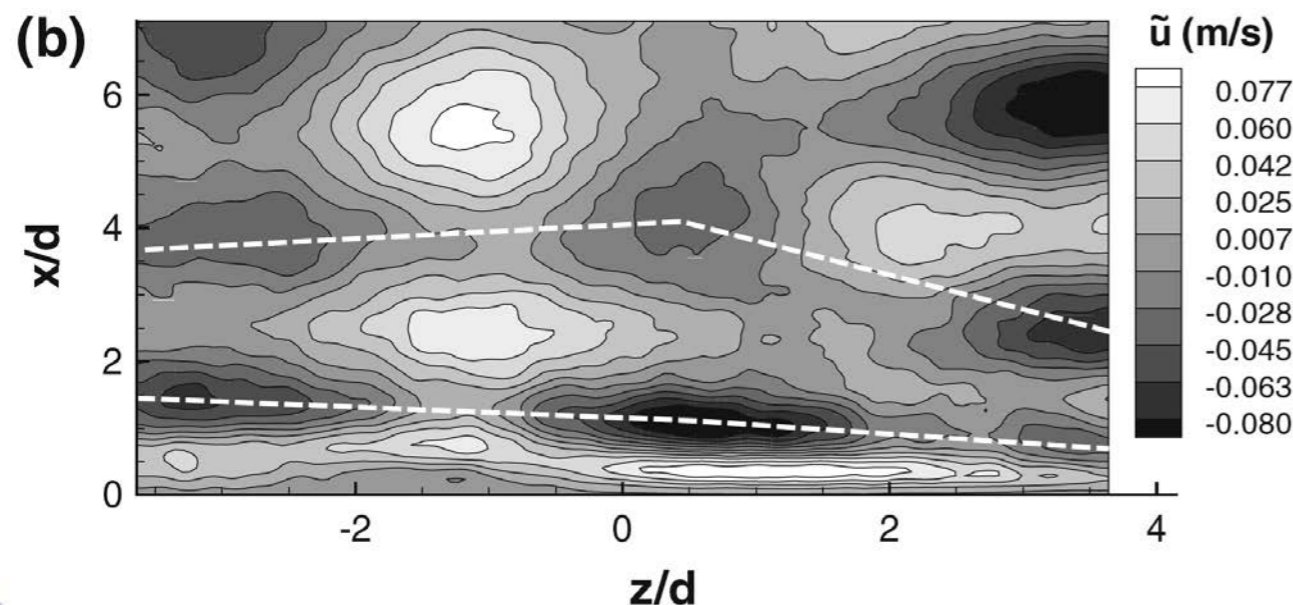
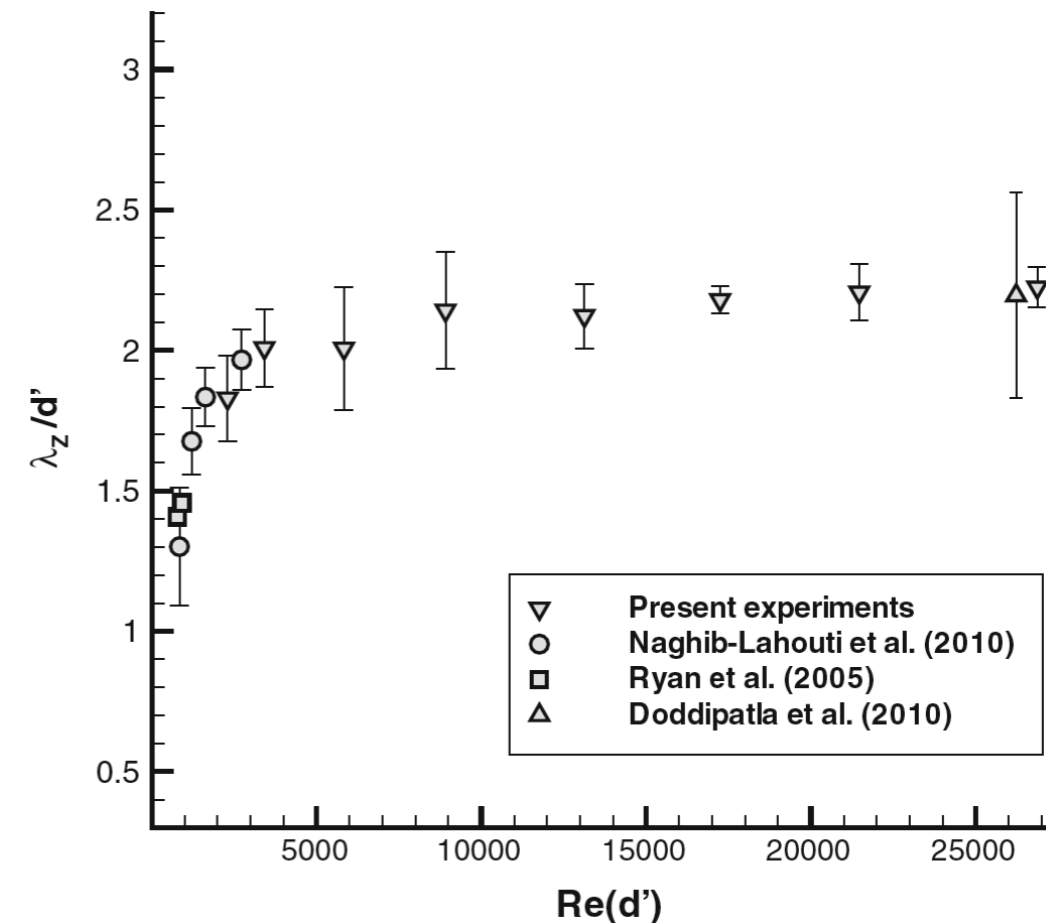
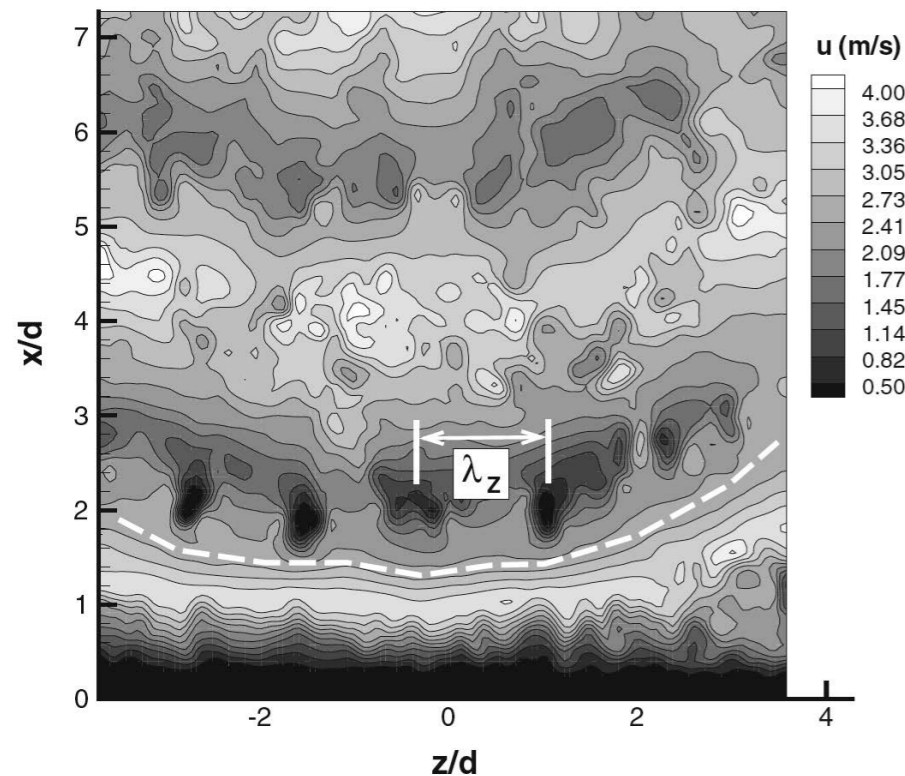
- Naghieb-Lahouti et al. (2012, 2014) used POD model to extract dominant structures in wake



# POD for Structure Identification



- Naghib-Lahouti et al. (2012, 2014) used POD model to extract dominant structures in wake

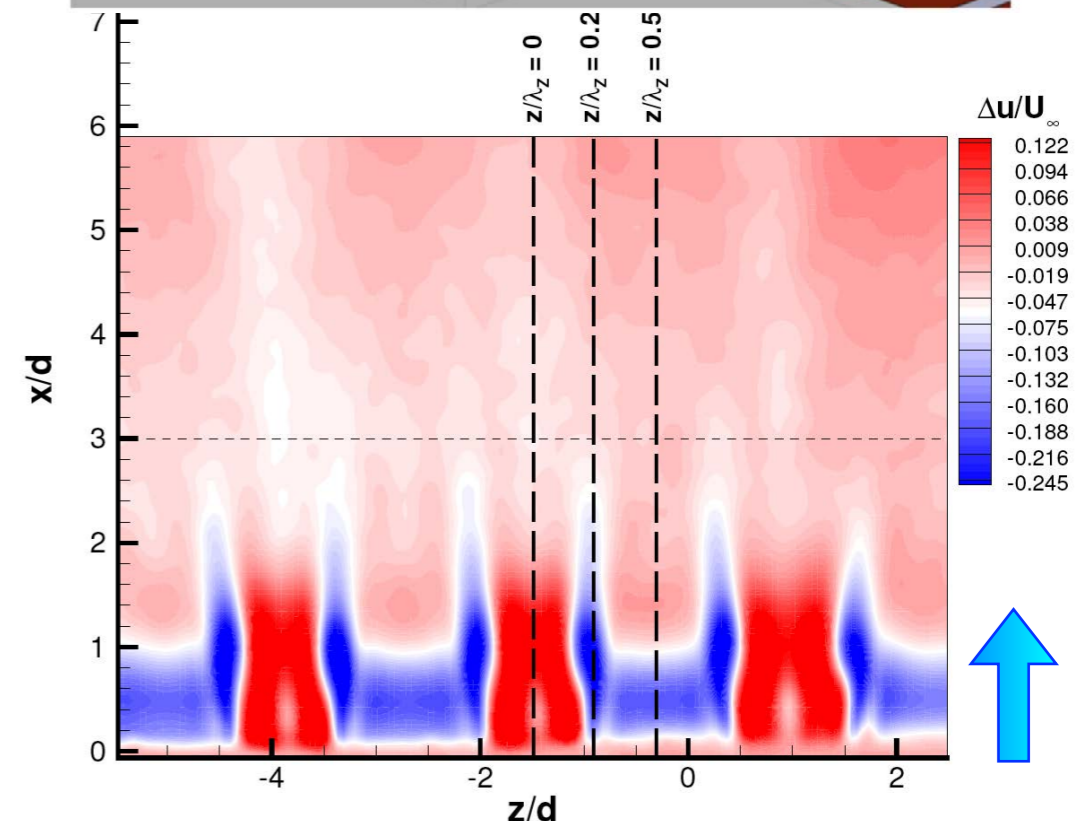
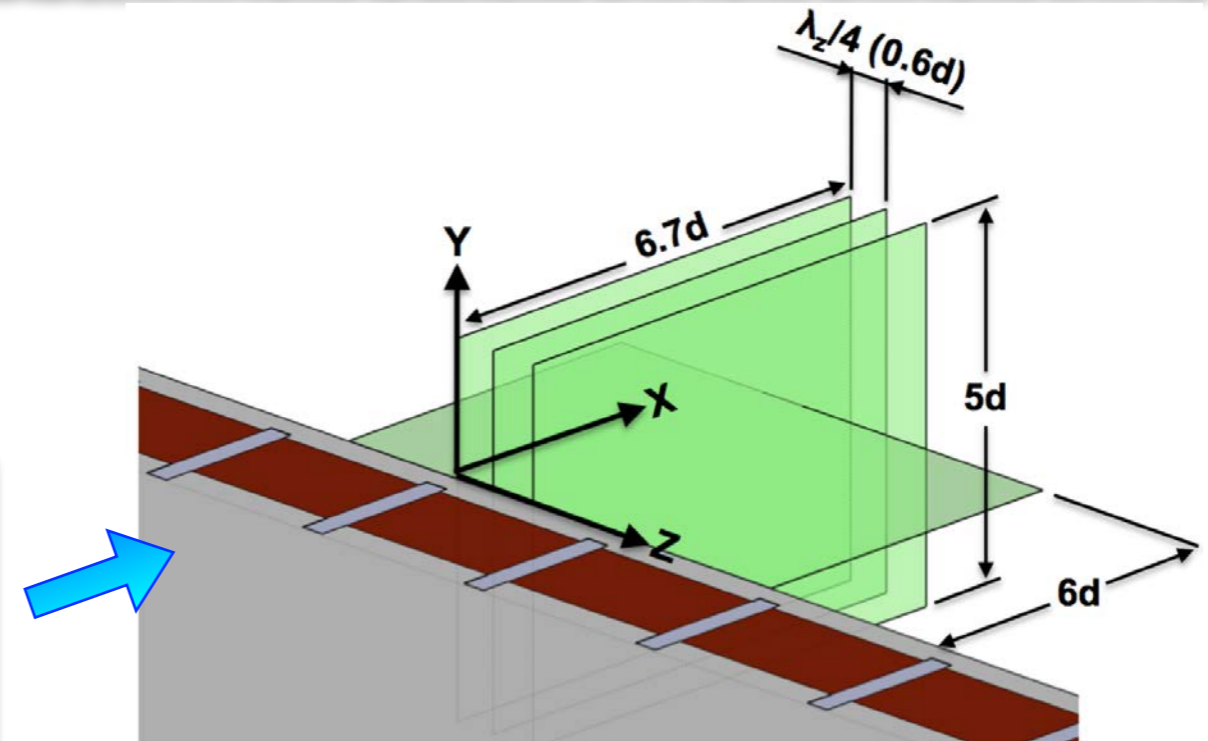
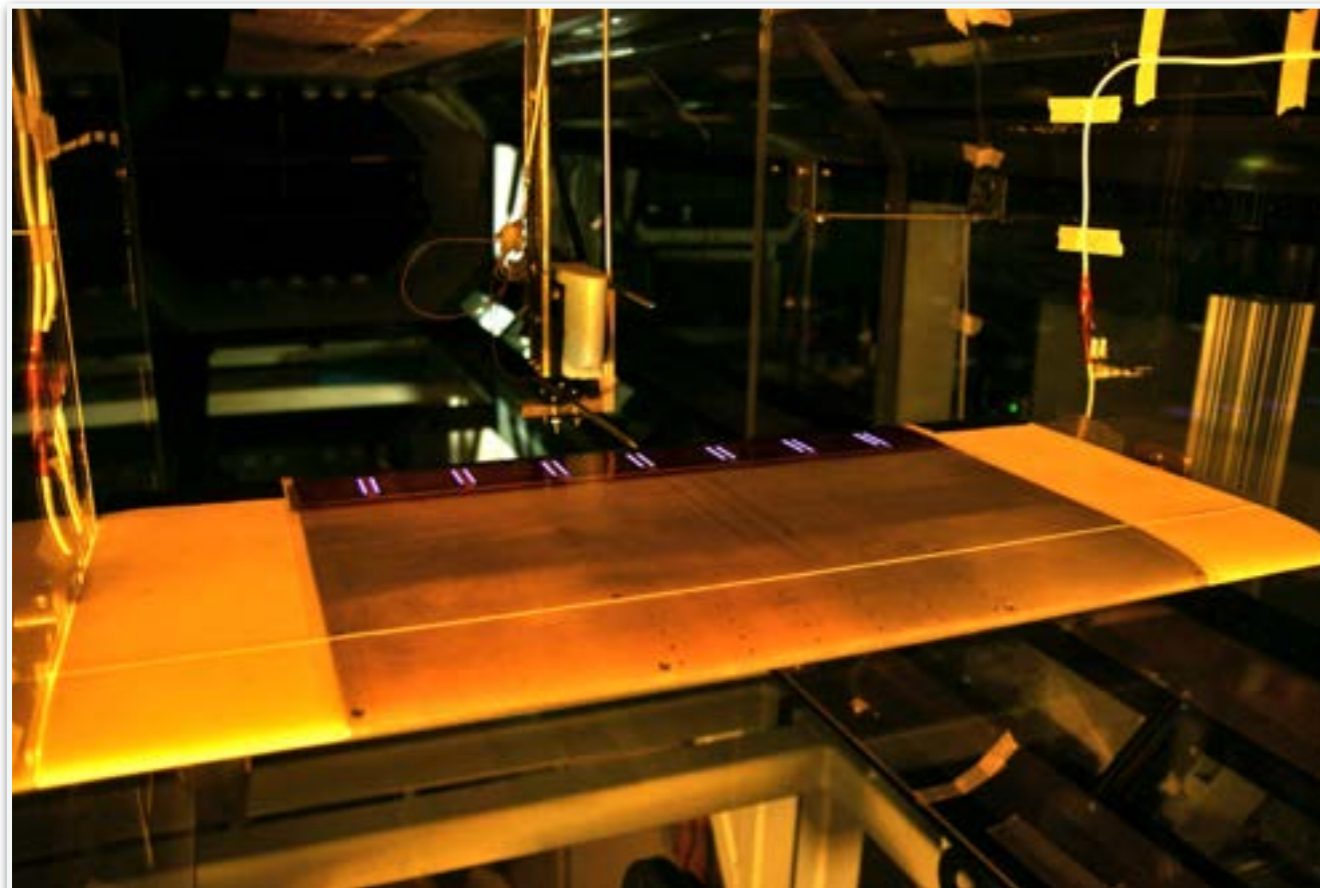


- Instability wavelength,  $\lambda_z$ , relatively constant with  $Re$ 
  - Implications for distributed forcing in wake control strategy

# Plasma Actuator Control

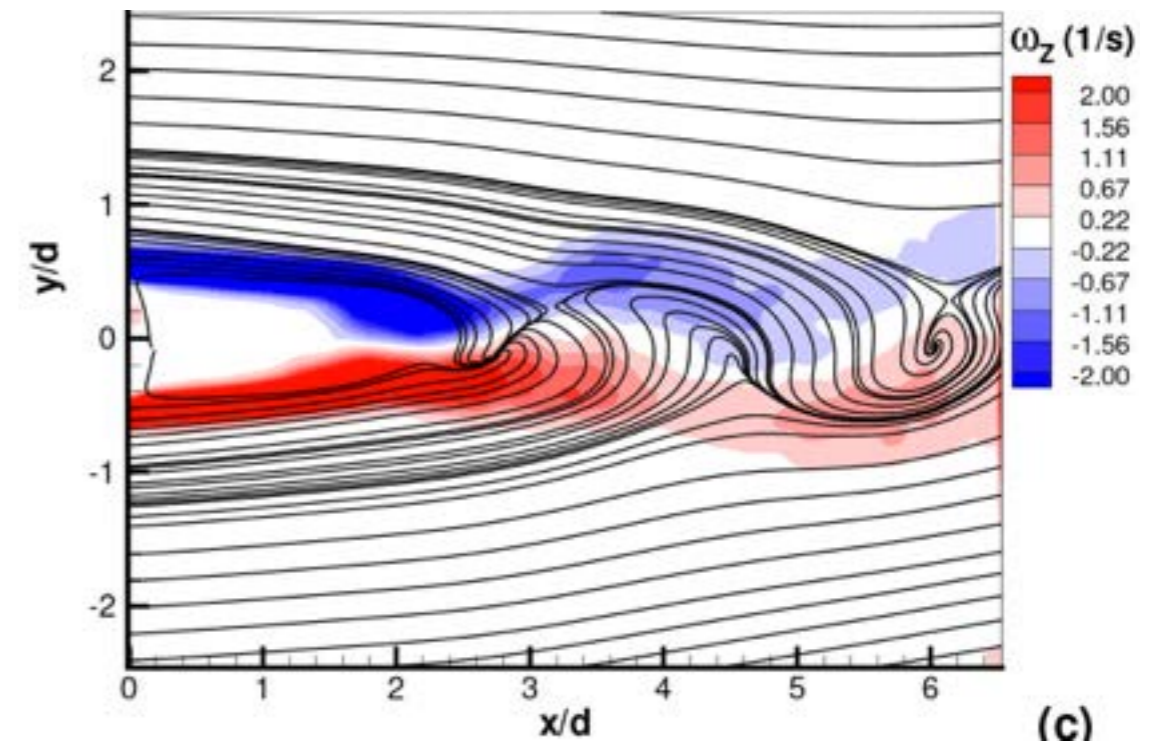
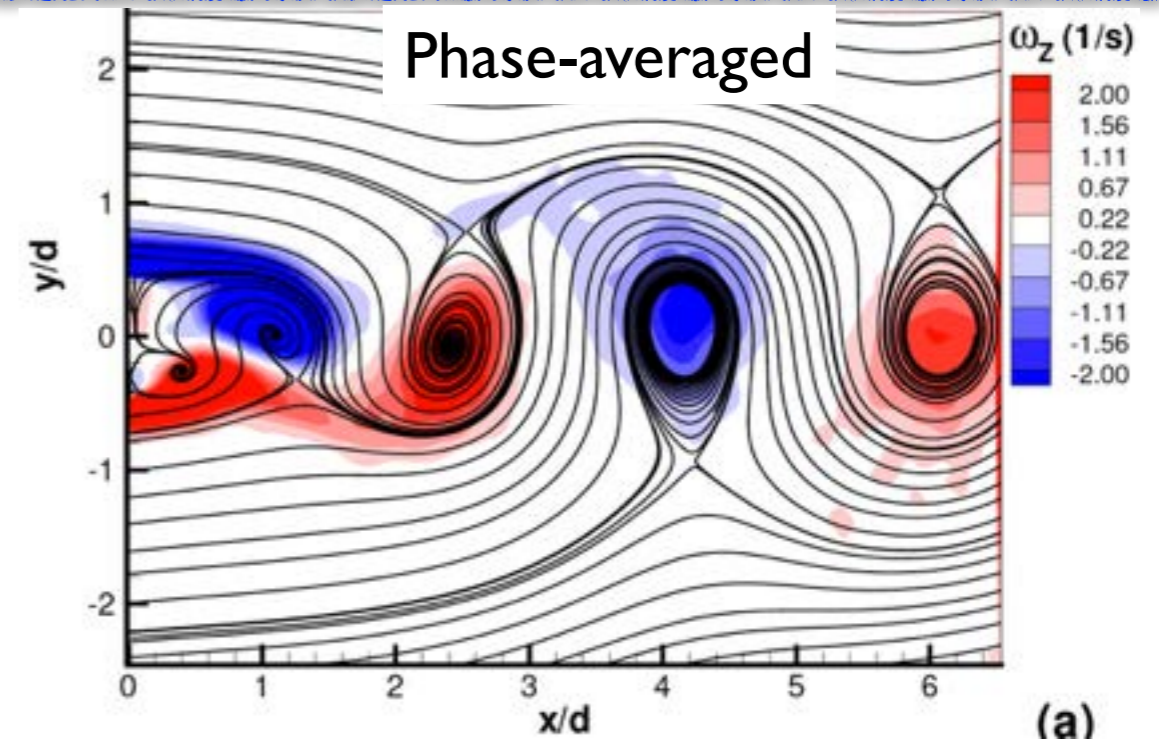
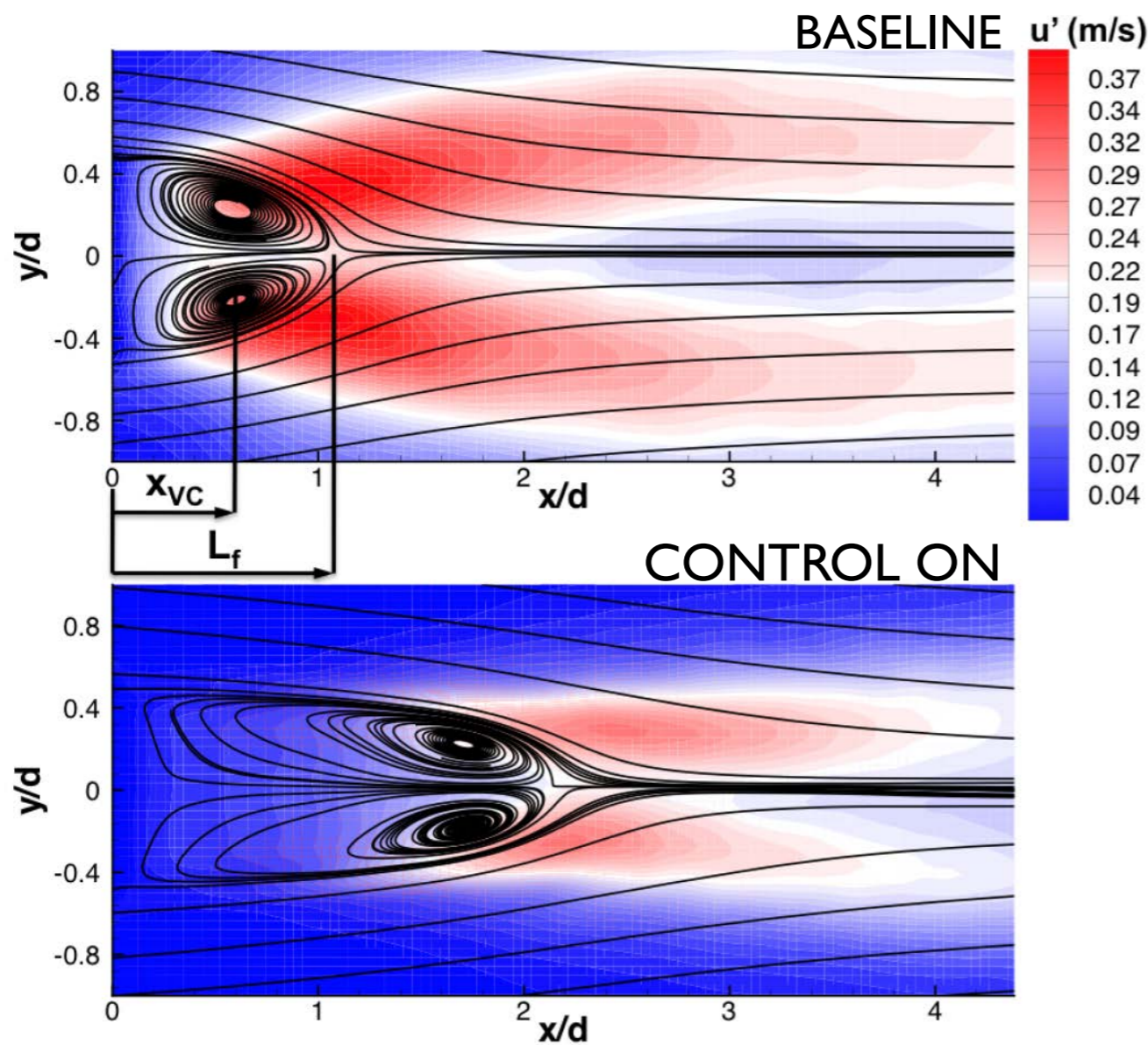


- $Re_d = 2,000, 3,000, \& 5,000$
- Energy input is varied through excitation parameters of actuators



Naghieb-Lahouti, Lavoie & Hangan, Phys. Fluids (2015)

# Plasma Actuator Control



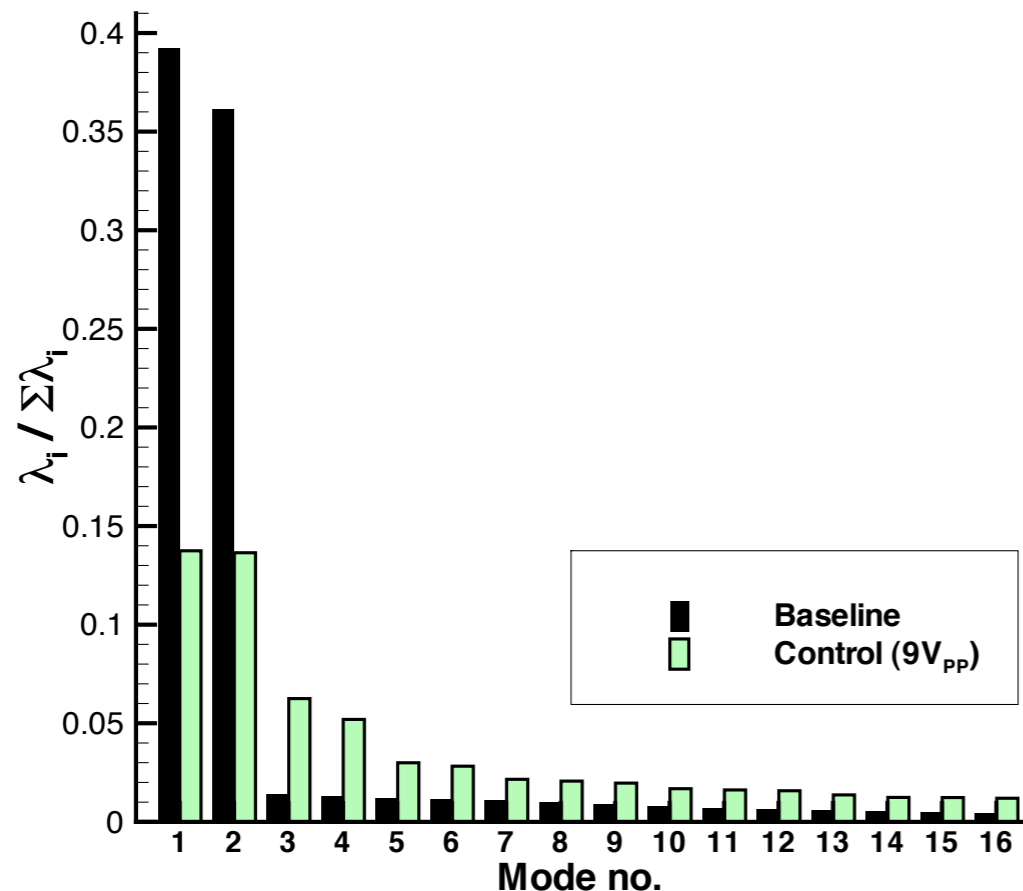
- Formation length extended
- Vortex formation decimated

Naghieb-Lahouti, Lavoie & Hangan, Phys. Fluids (2015)

# Plasma Actuator Control

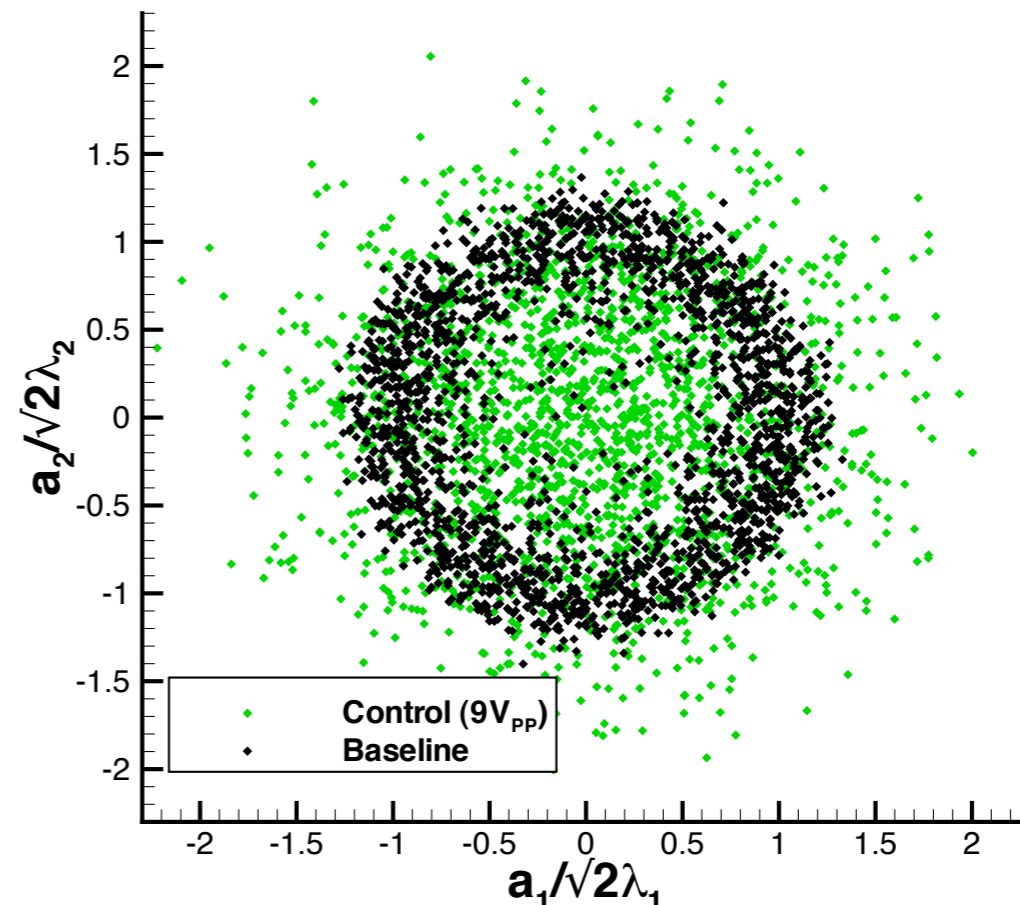


Can be seen from the change in the POD modes



Relative energy of POD modes in (x-y) plane

Re(d) = 2000, z/λ = 0



Phase plot - (x-y) plane

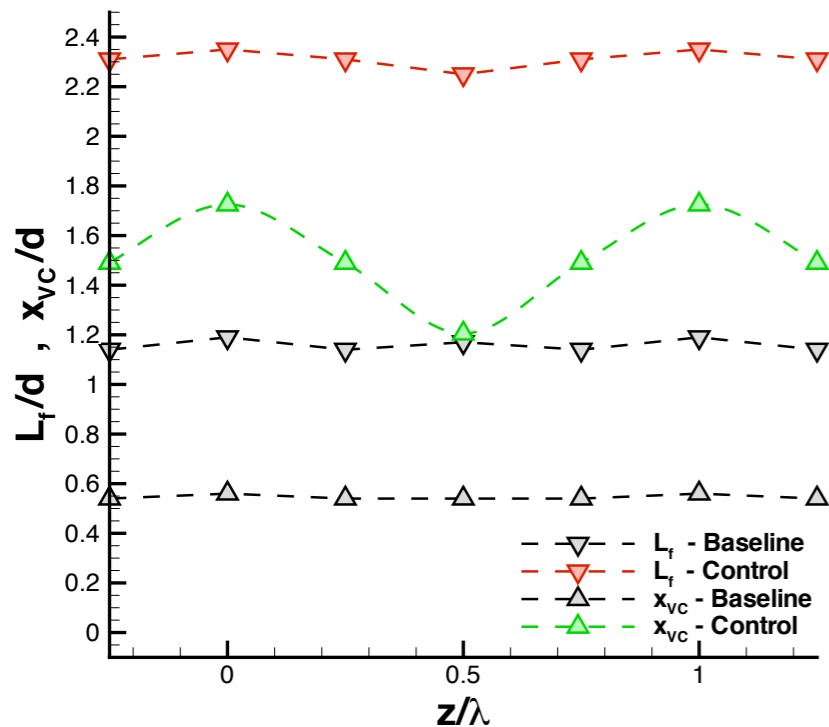
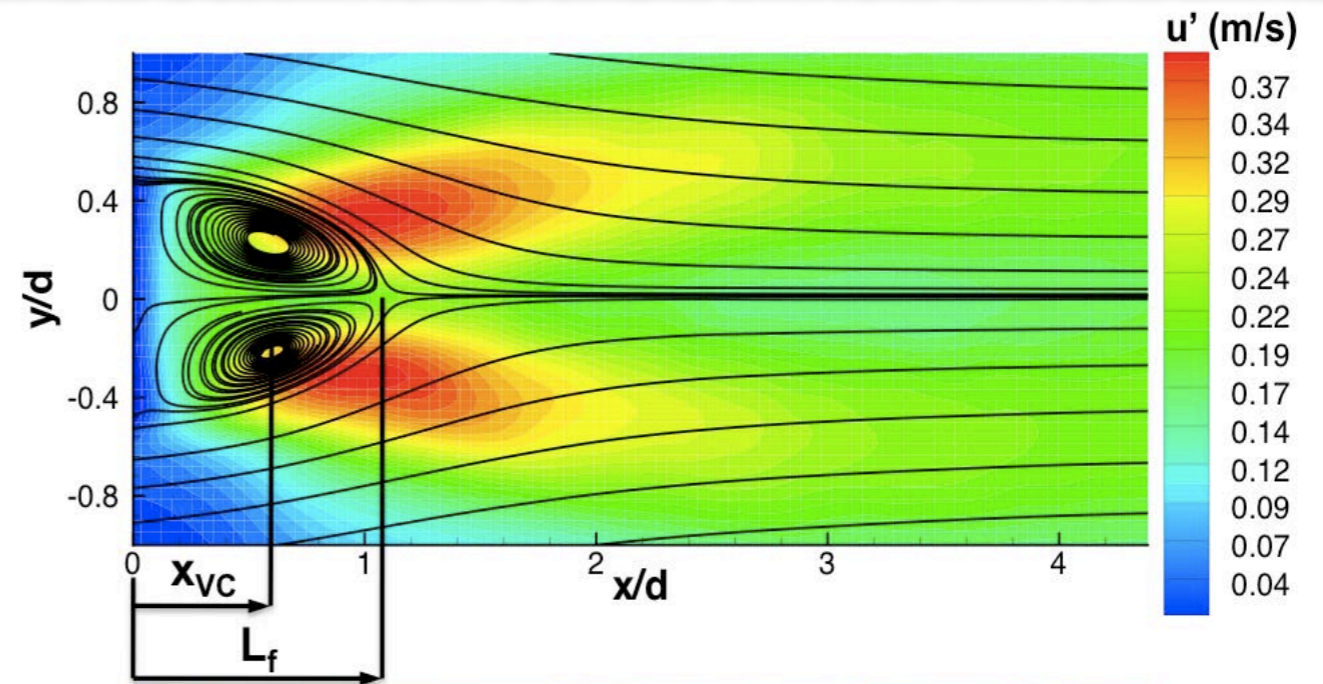
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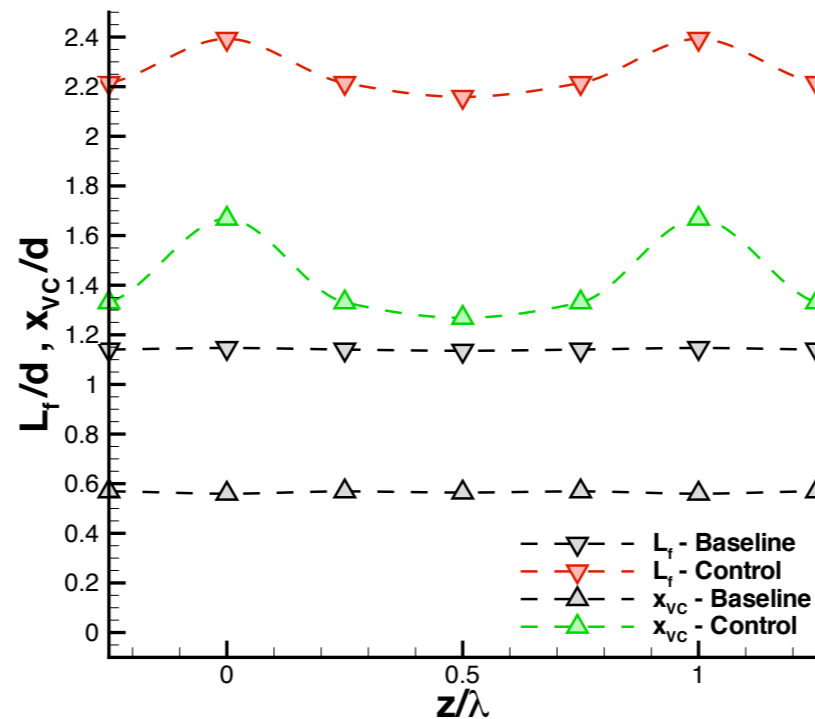
# Plasma Actuator Control



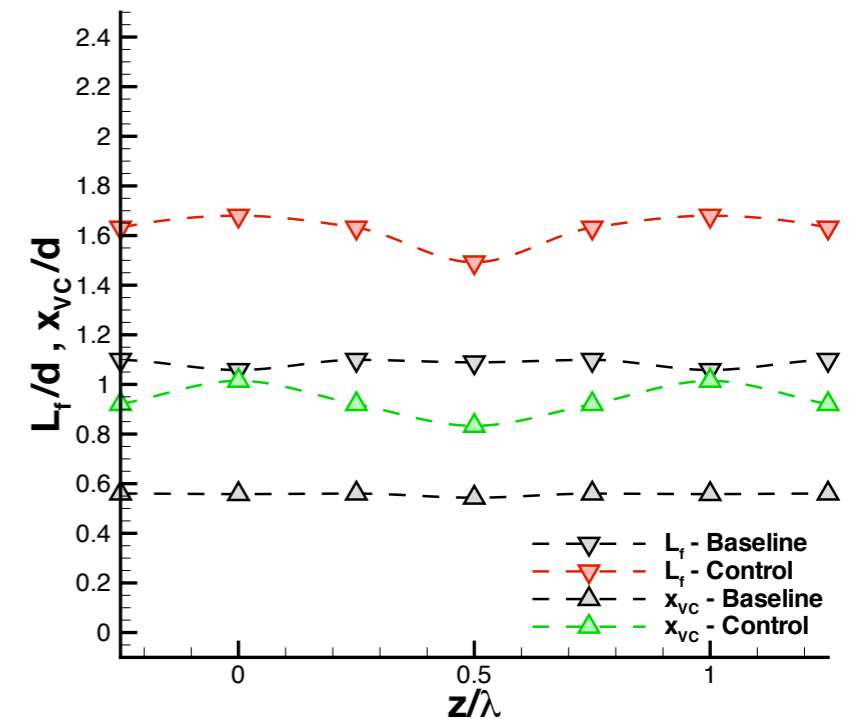
Effect of forcing extend over the full span.



**Re(d) = 2000**  
**9kV<sub>pp</sub>**



**Re(d) = 3000**  
**10kV<sub>pp</sub>**



**Re(d) = 5000**  
**11kV<sub>pp</sub>**

Naghib-Lahouti, Lavoie & Hangan, Phys. Fluids (2015)

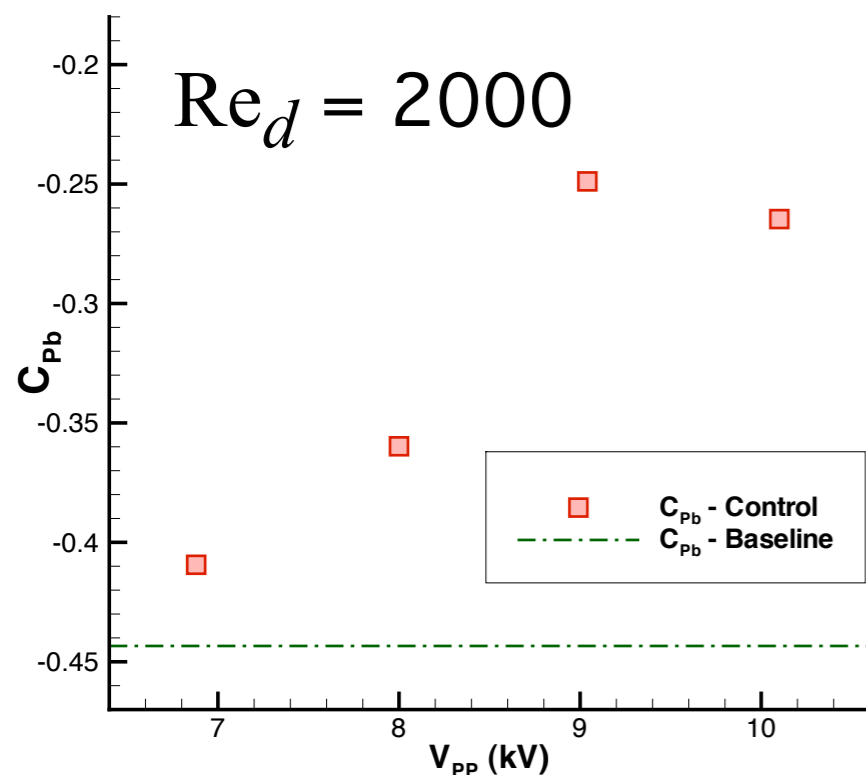
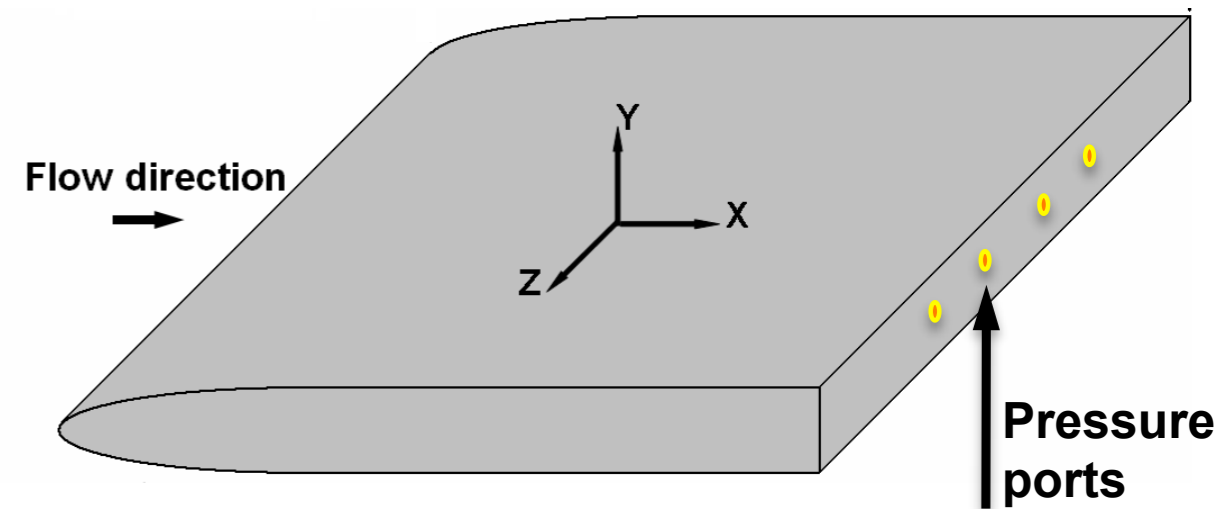


# Plasma Actuator Control

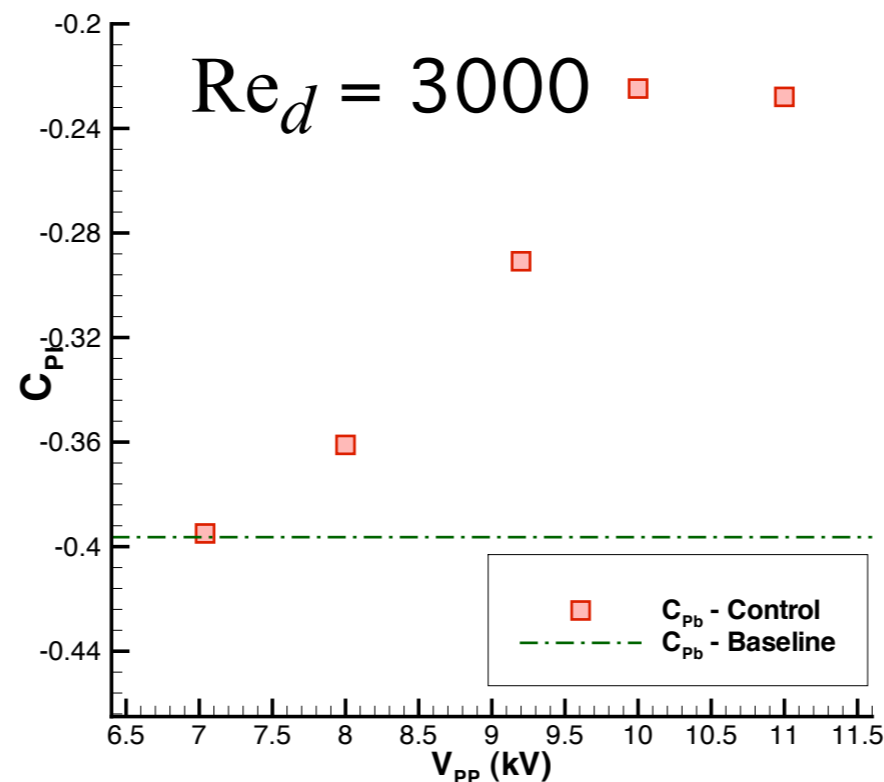


Base pressure coefficient based on the average pressure measured through 4 ports across the span.

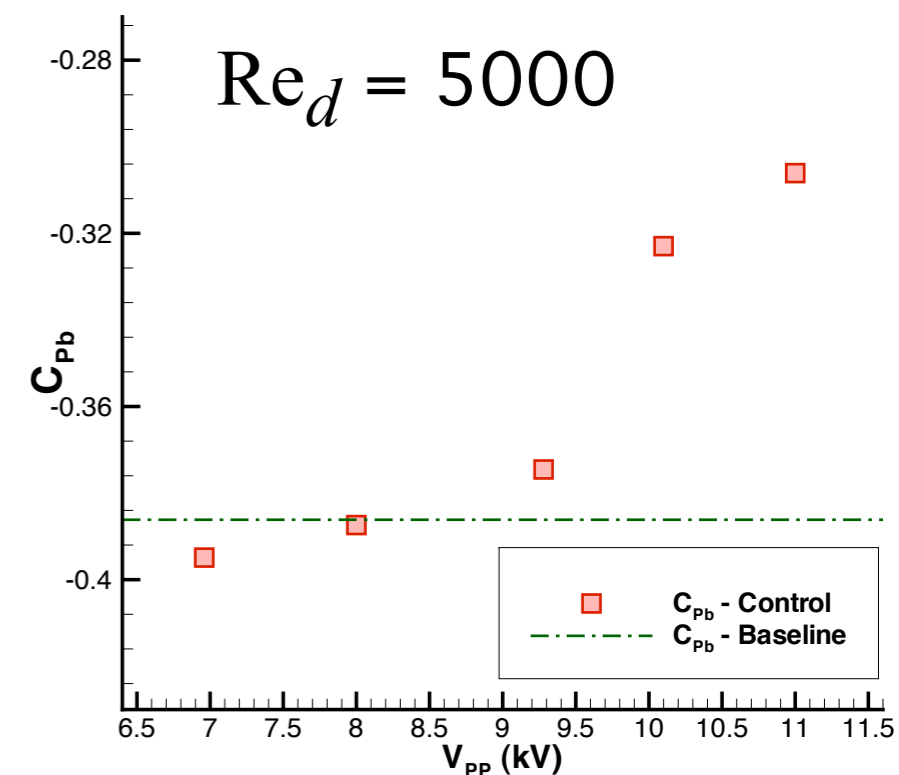
Significant recovery of base pressure, consistent with the behaviour of  $L_f$  and  $x_{VC}$



44% recovery



43% recovery



21% recovery

# Plasma Actuator Control

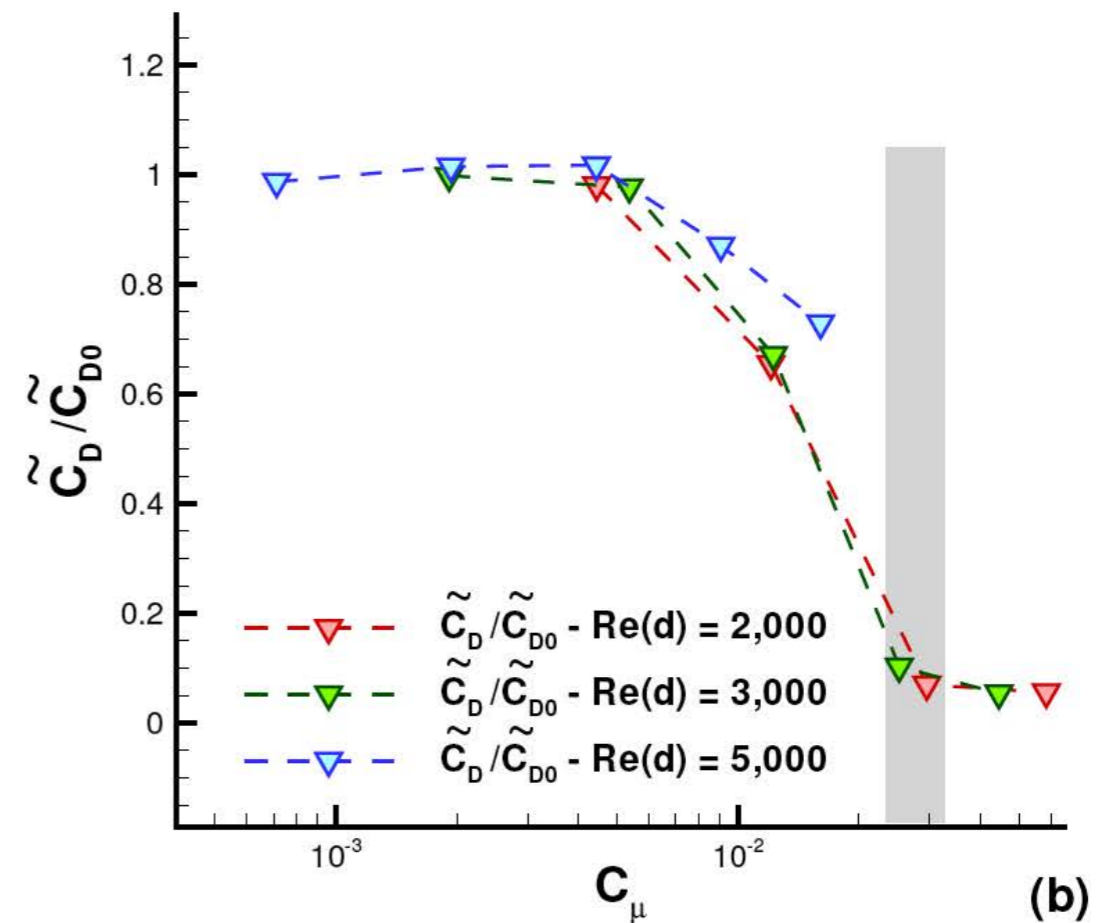
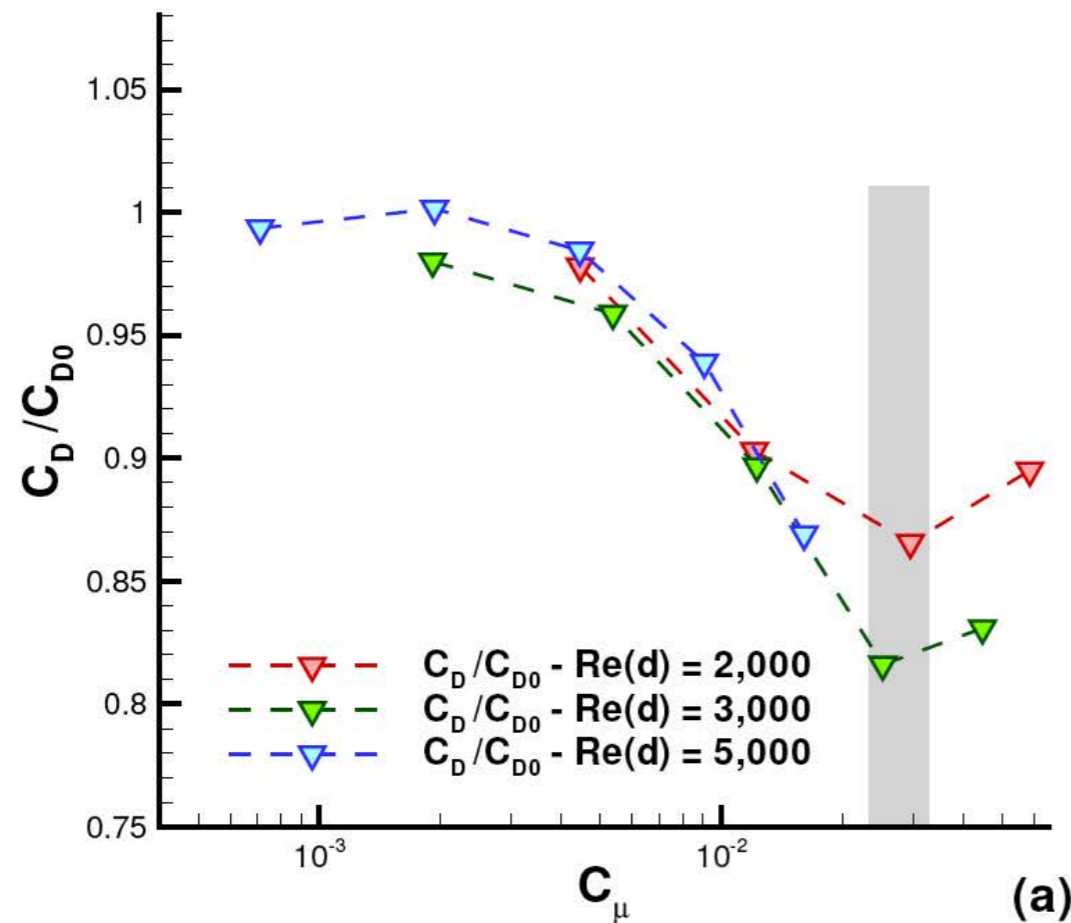


Drag estimate based on flow measurements

(van Oudheusden et al., 2007; Bohl & Koochesfahani, 2009)

$\tilde{C}_D$

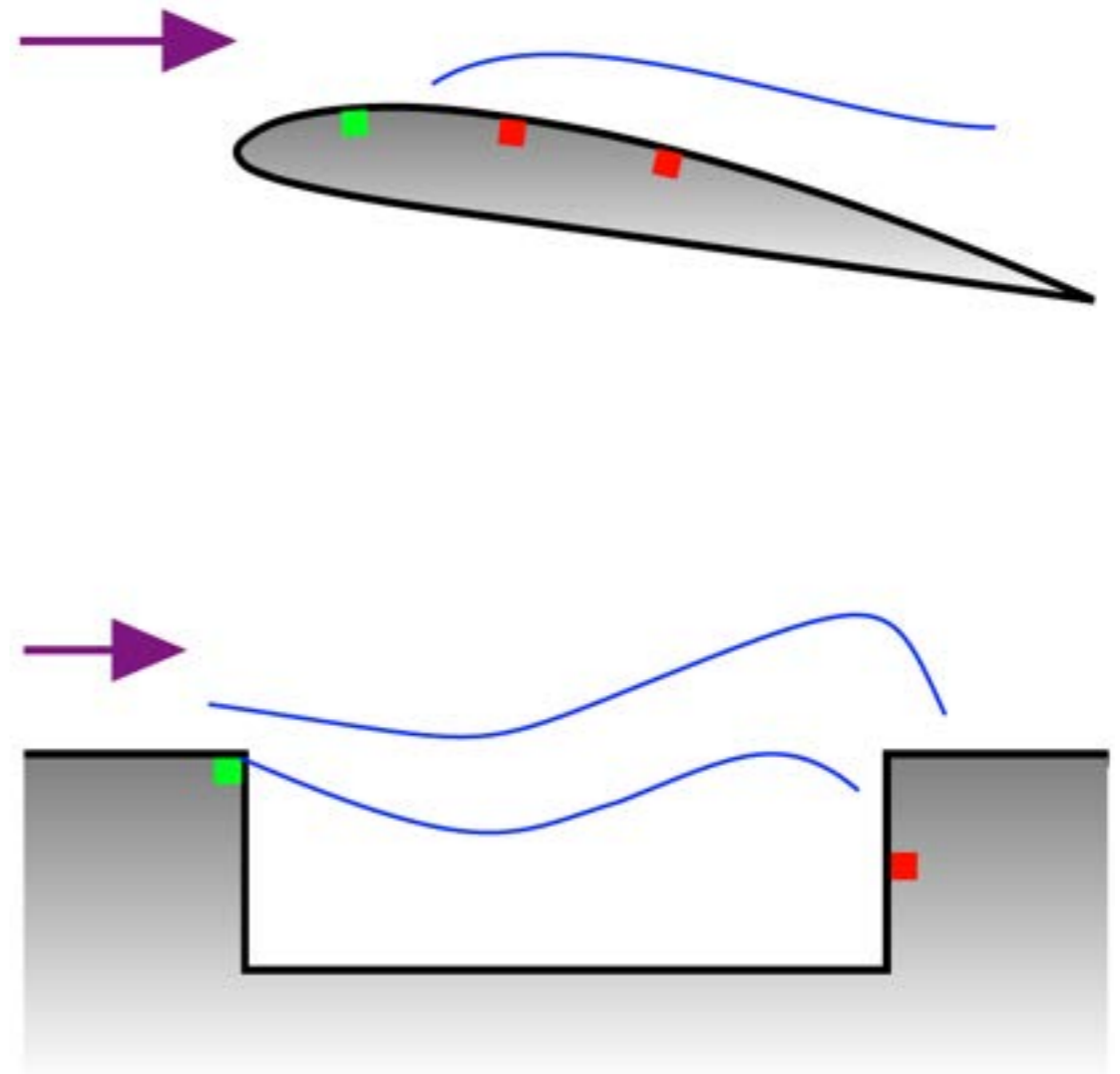
$$C_D = \frac{2}{A} \int \frac{\bar{u}}{U_\infty} \left(1 - \frac{\bar{u}}{U_\infty}\right) dy - \frac{2}{A} \int \left(\frac{u'}{U_\infty}\right)^2 dy + \frac{2}{A} \int \left(\frac{v'}{U_\infty}\right)^2 dy$$



$$C_\mu = \frac{\text{momentum flux induced by actuators}}{\text{freestream momentum flux through base area of body}}$$

- Estimation in complex flows: coherent structures analysis and low-dimensional modelling, closed-loop control
- Two broad categories (Cattafesta et al., 2008)
  - **static estimator** based on empirical mapping
  - **dynamic estimator** based on physical model

Clark, Naghib-Lahouti & Lavoie, EiT (2014)

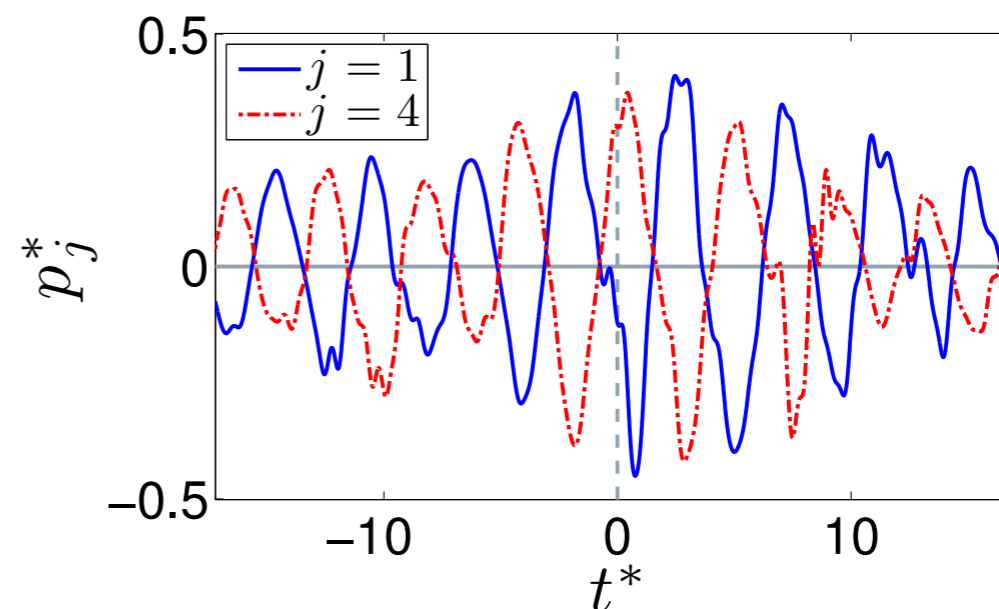
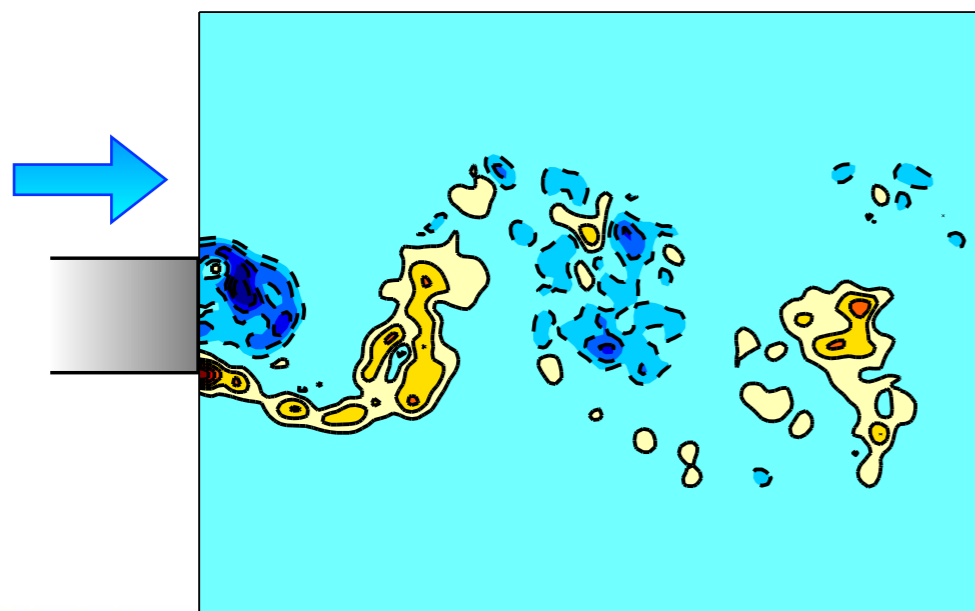
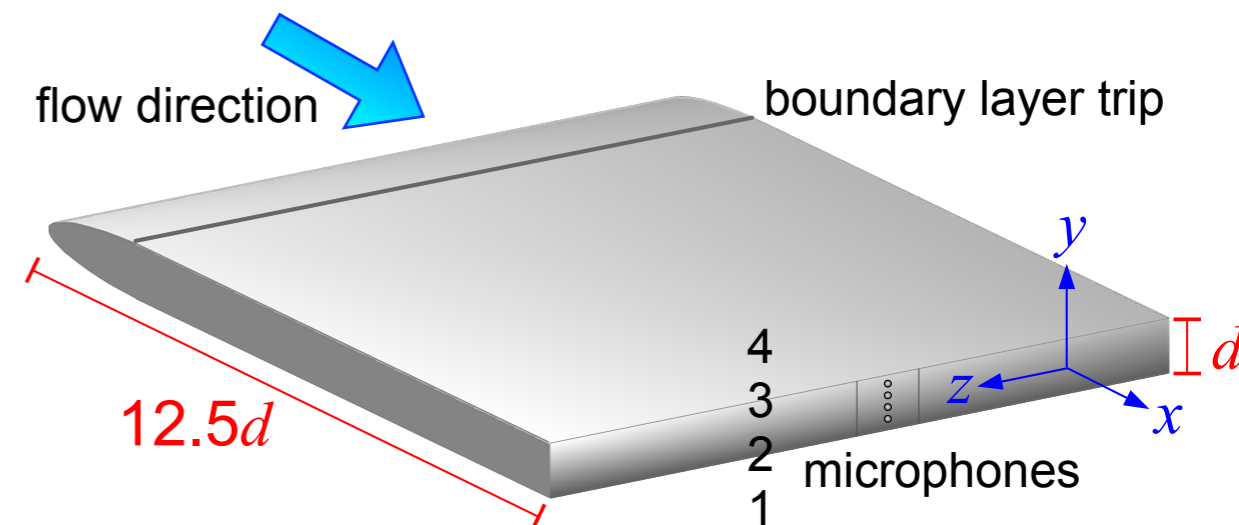


## Objective:

Estimate states of interest from limited surface sensing

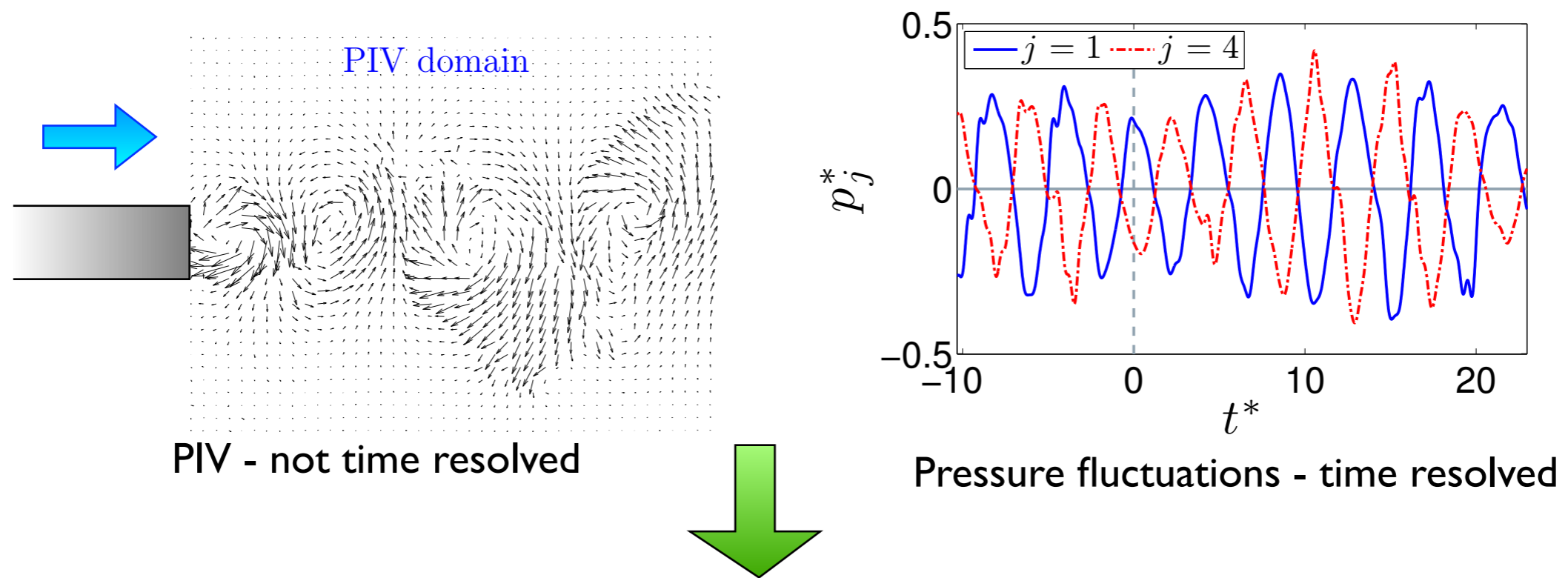
# Estimation for Control

- **Objective:** extend distributed forcing control strategy to closed-loop method
- Actuation depends on state of flow, requires nonintrusive real-time measurements
- Advantages
  - Reduce actuator power consumption
  - Robustness to external disturbances and conditions
- Base pressure fluctuations indicate time evolution of local vortex shedding



# Empirical Estimator

1. Obtain synchronized measurements of velocity and fluctuating pressure in the wake



2. Construct model from flow statistics to estimate velocity from pressure signals only

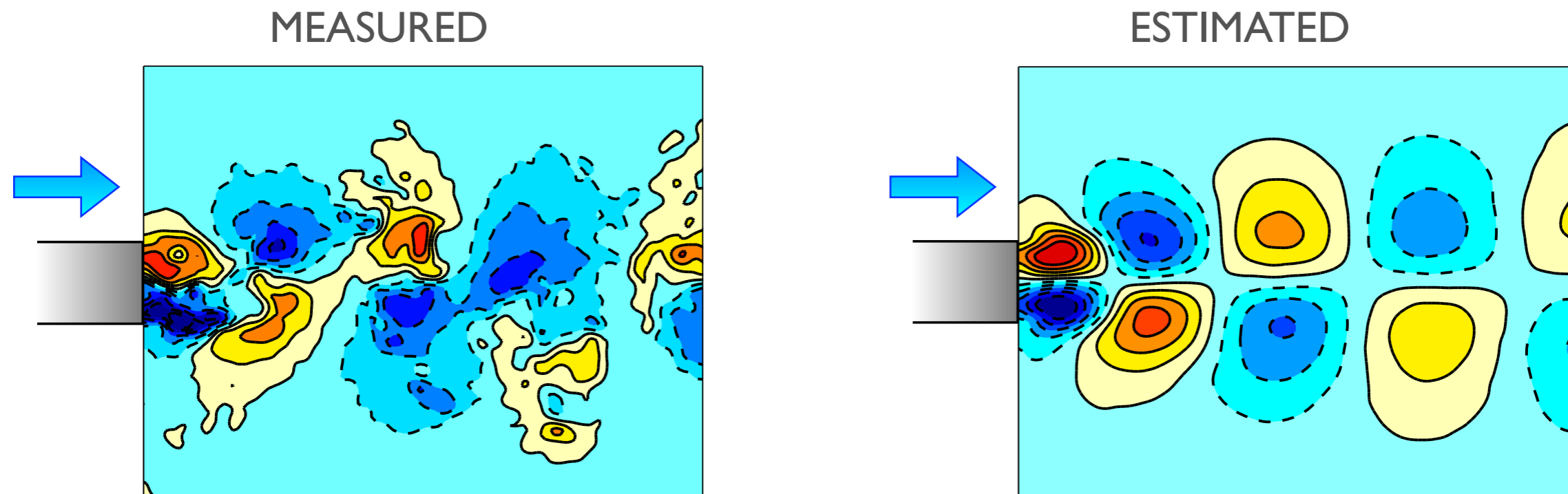
$$p_j(t) \rightarrow \mathcal{F} \rightarrow \tilde{\mathbf{u}}(x, y, t)$$

**Reduced-order model**

# POD for Reduced Order Modelling

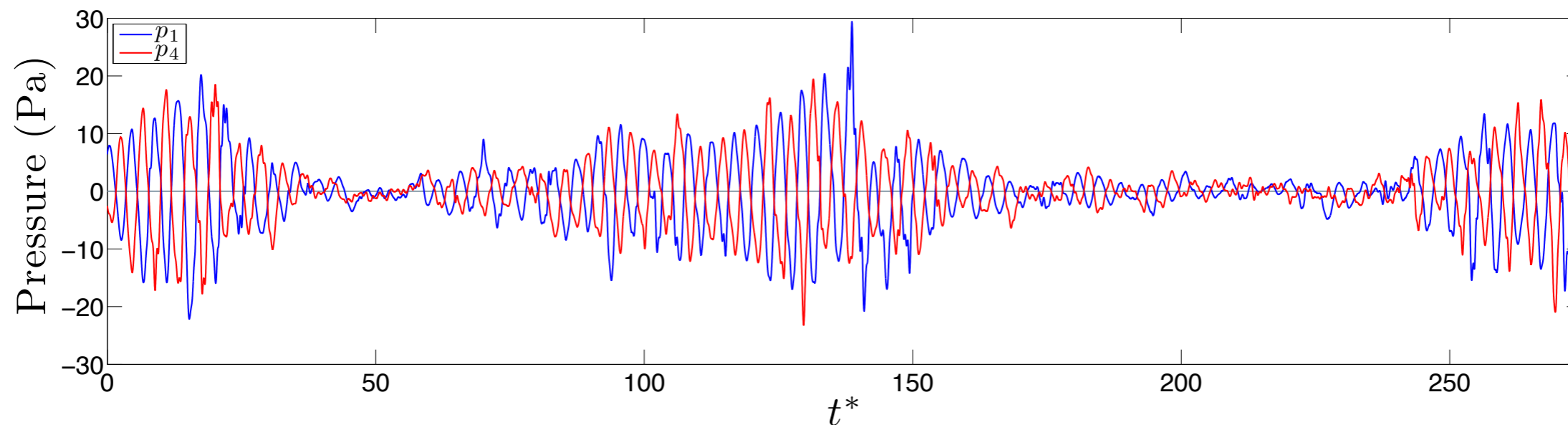
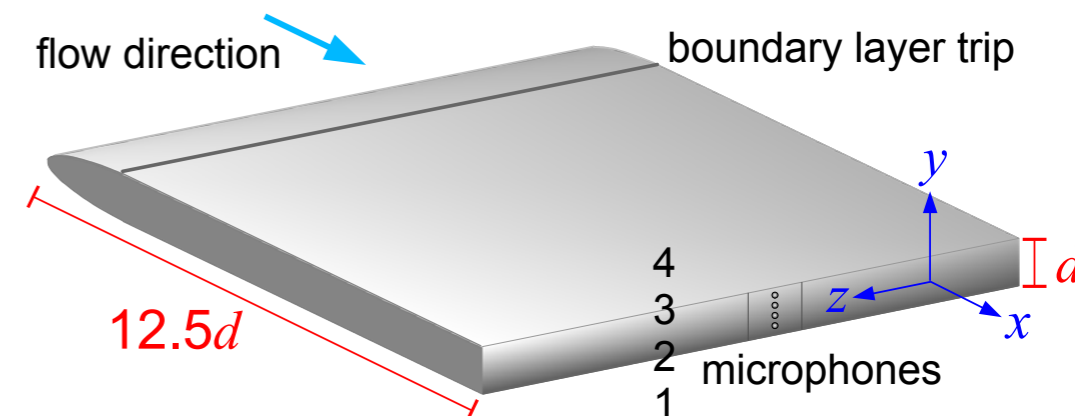


- Usually not interested in knowing the velocity at every point - can use POD modes to reduce the order of the problem and provide a **reduced order model**
- Estimate evolution of the large coherent structures only (POD model)
  - Simple model required for real-time feedback control



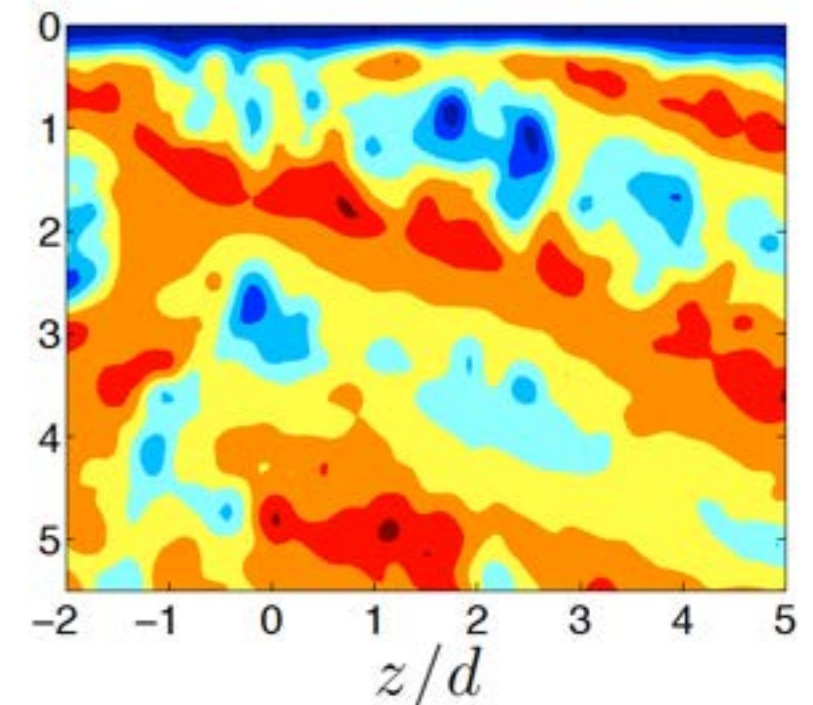
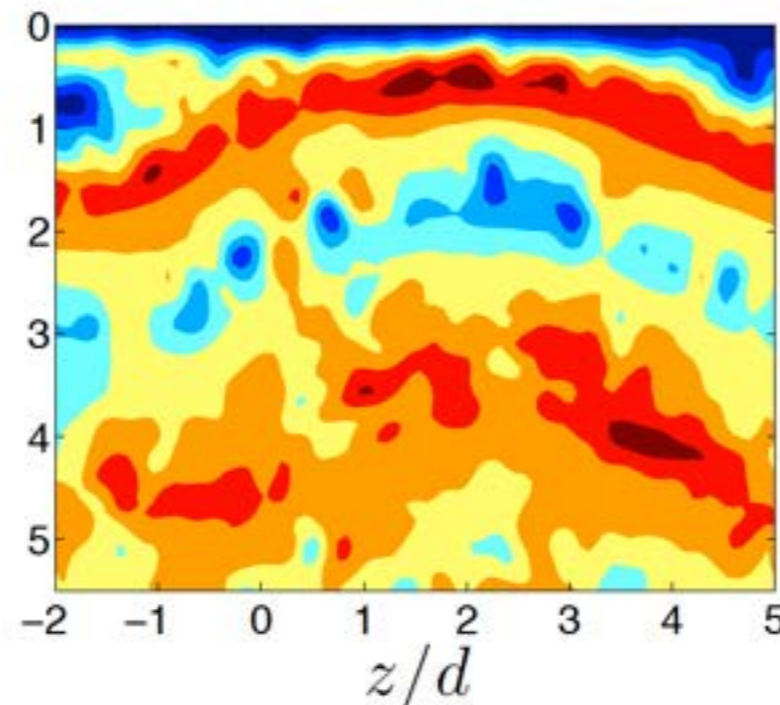
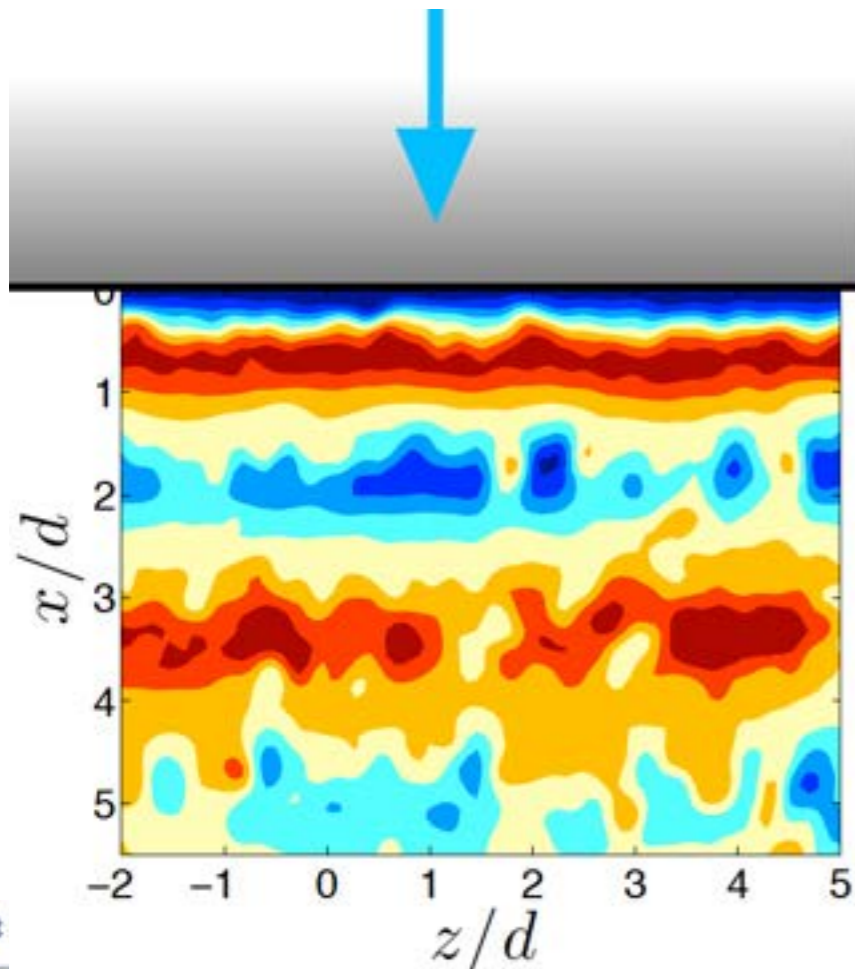
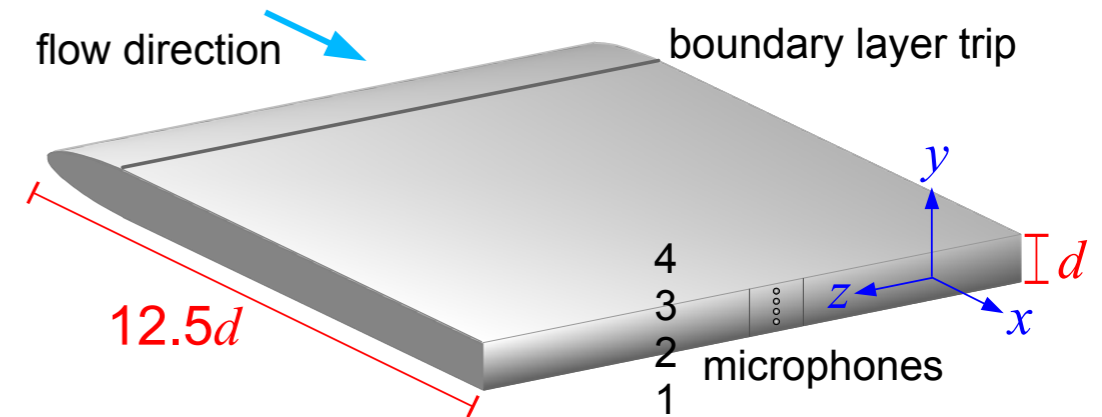
# 3D Complications

- Low-frequency modulation of base pressure fluctuations
- Wu et al. (2005), normal trapezoidal body; Lemkuhl et al. (2013), cylinder: modulation associated with unsteady variation in vortex formation length



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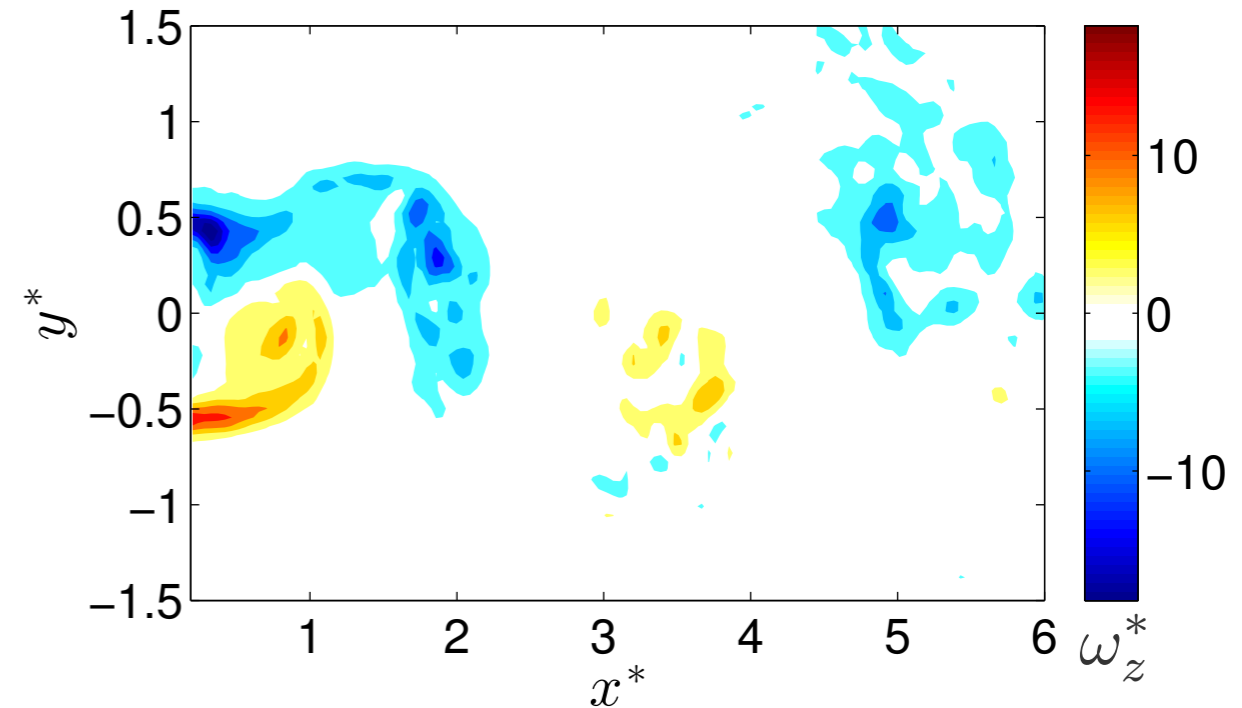
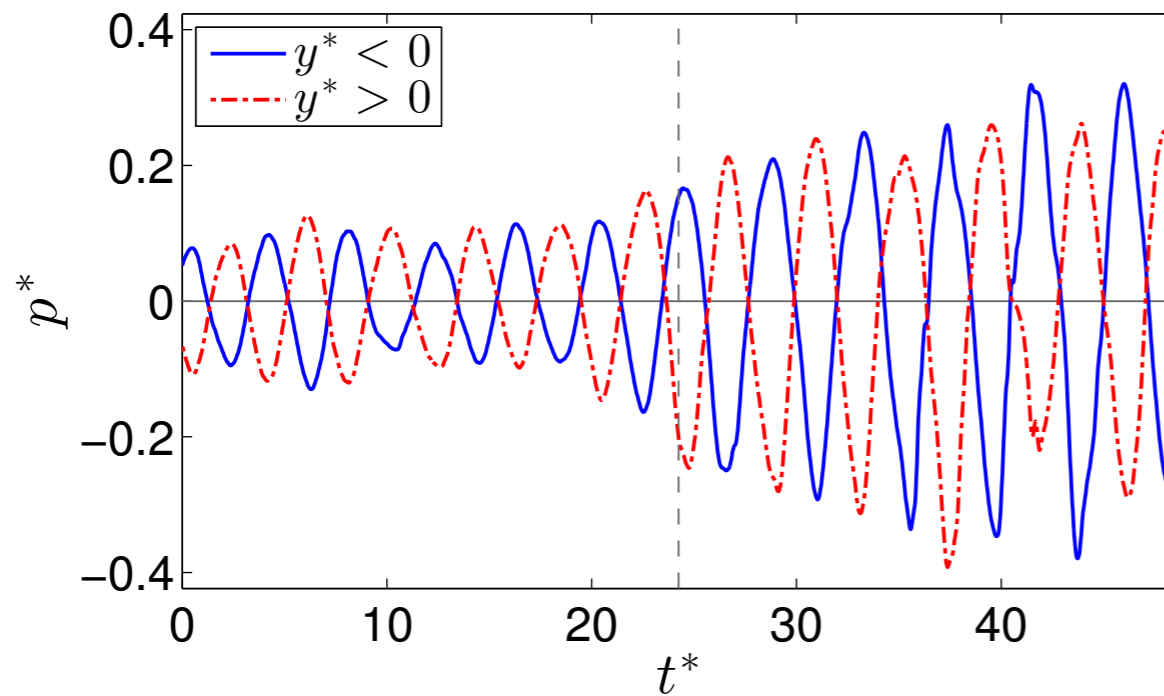




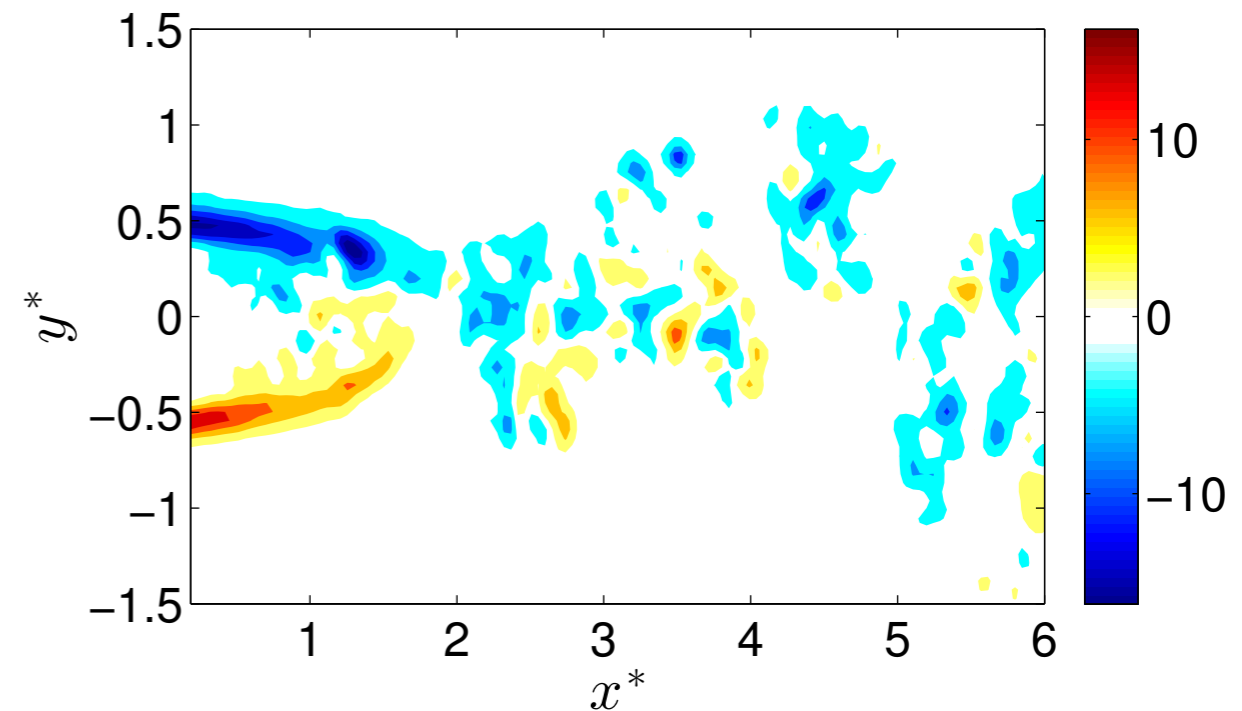
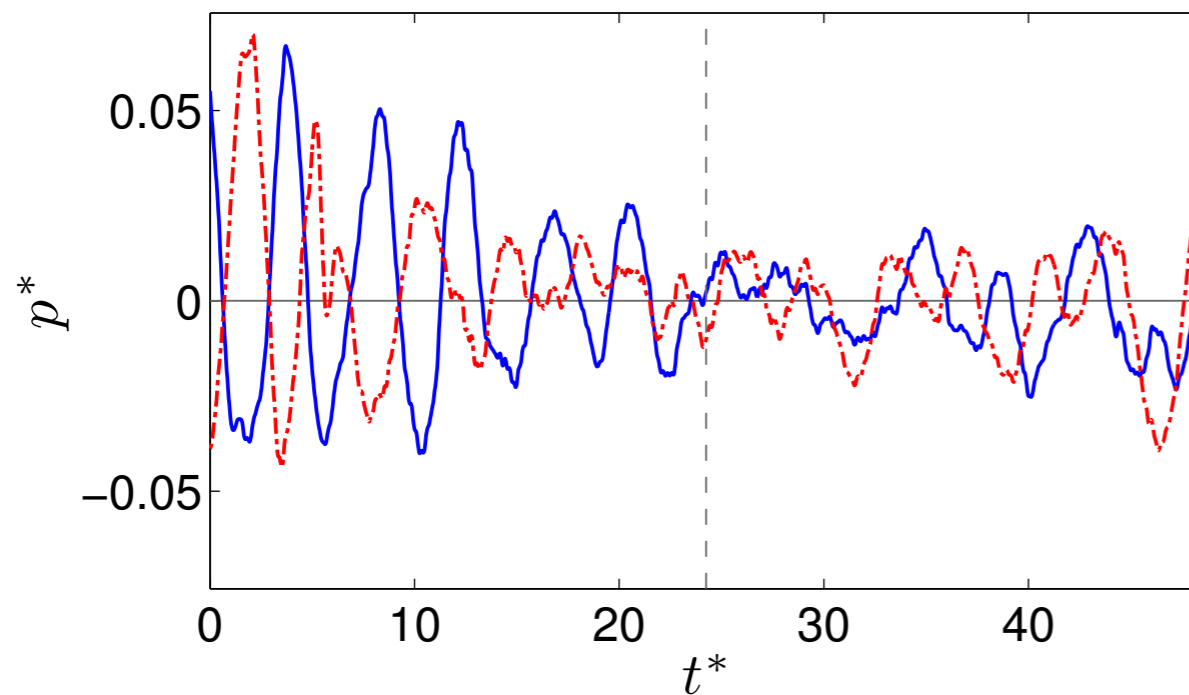
# Variable state of shedding



COHERENT

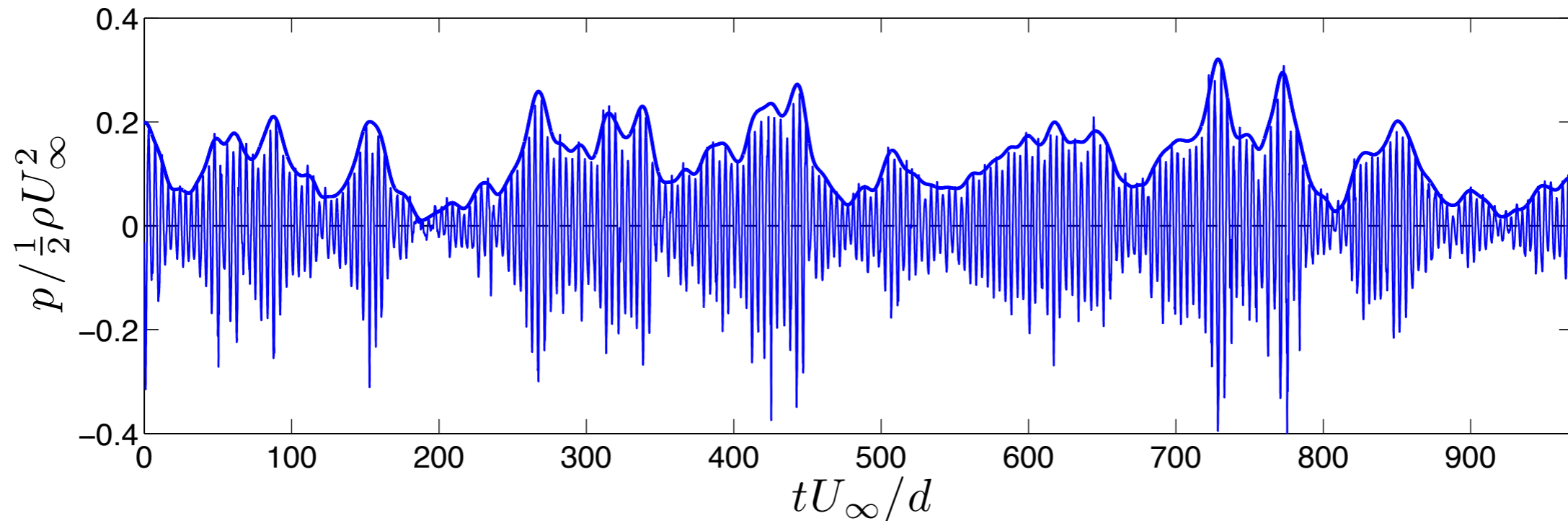


DISORGANIZED



- Phase disorganization typically occurs with low amplitude fluctuations at the same  $z^*$

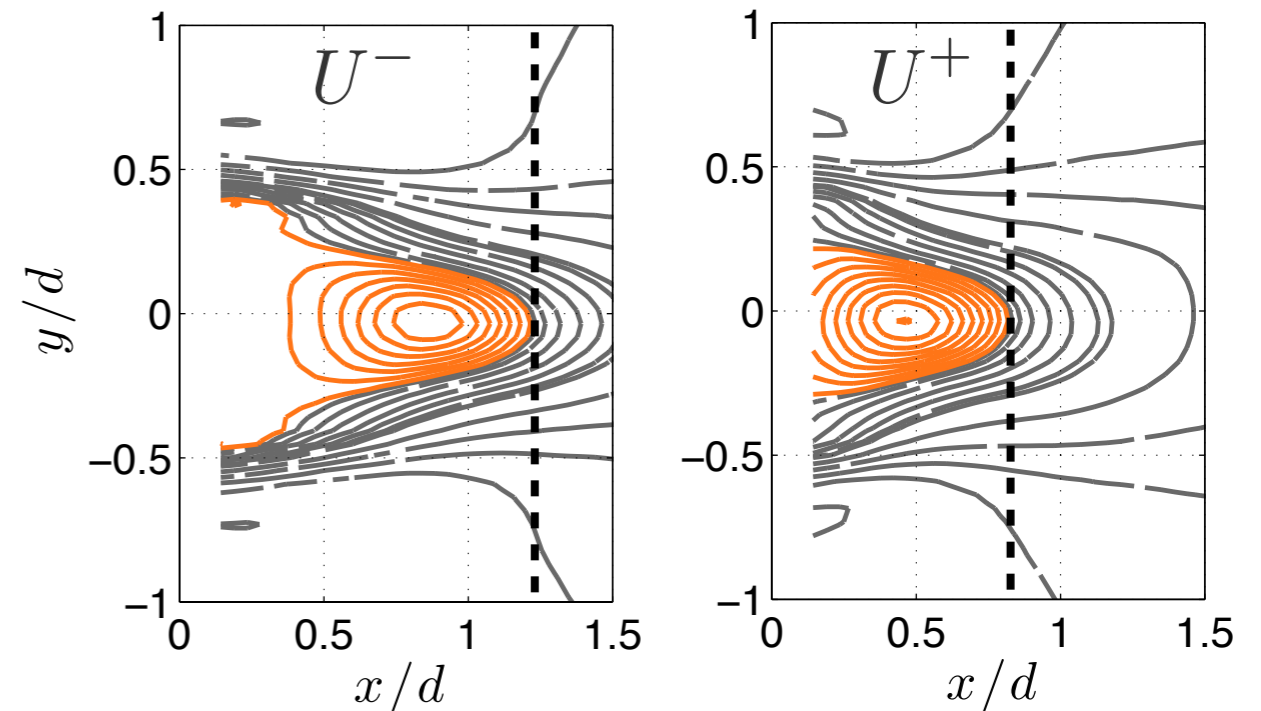
# Amplitude modulation



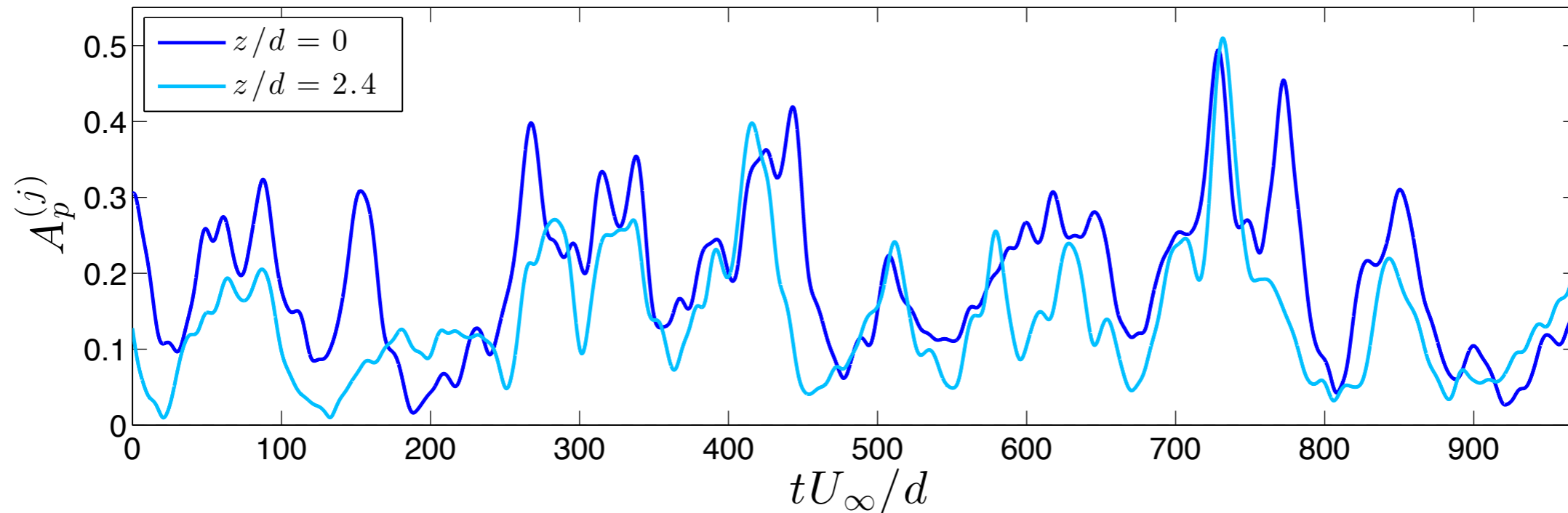
- $A_p(t)$  determined from wavelet transform, using MATLAB software by Torrence & Compo (1998)
- Szepessy (1994), Wu *et al.* (2005), Lehmkuhl *et al.* (2013): modulation at frequency one order of magnitude less than shedding
- Conditional average based on amplitude

$$U^- = \langle u \mid A_p \leq \overline{A_p} - \beta \sigma_A \rangle$$

$$U^+ = \langle u \mid A_p > \overline{A_p} + \beta \sigma_A \rangle$$



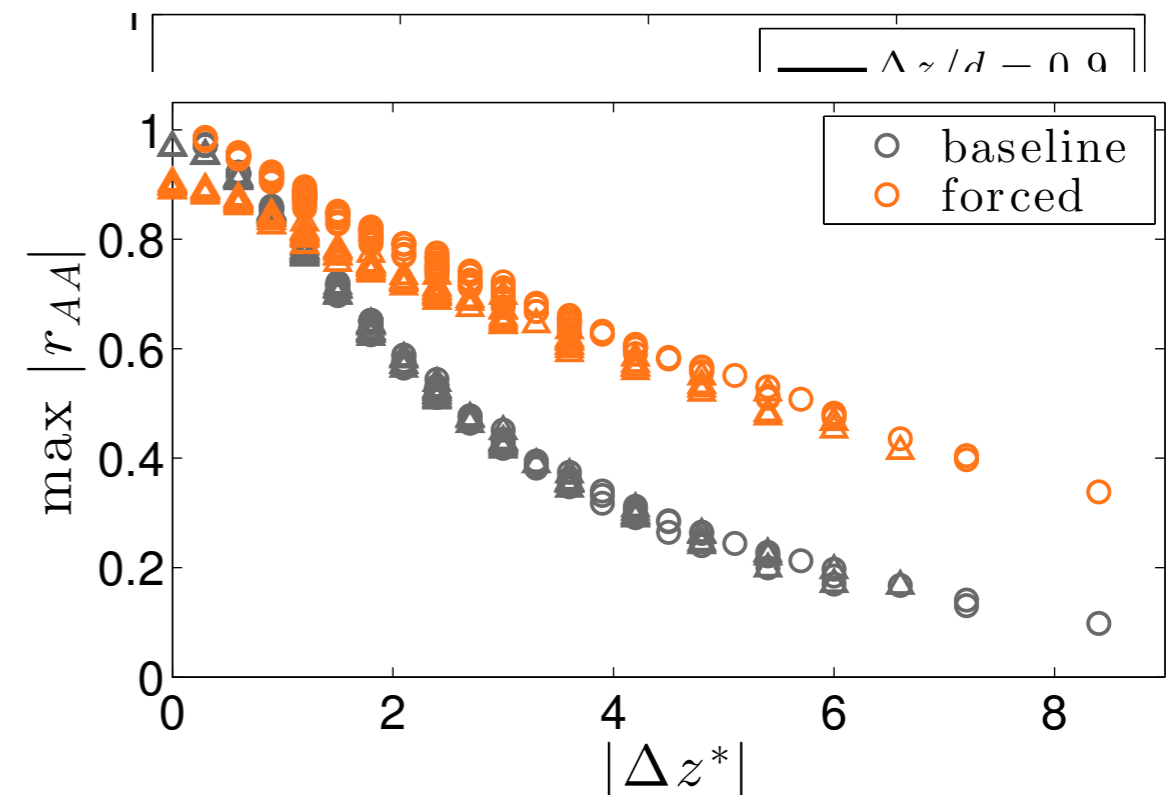
# Spanwise Amplitude Correlation



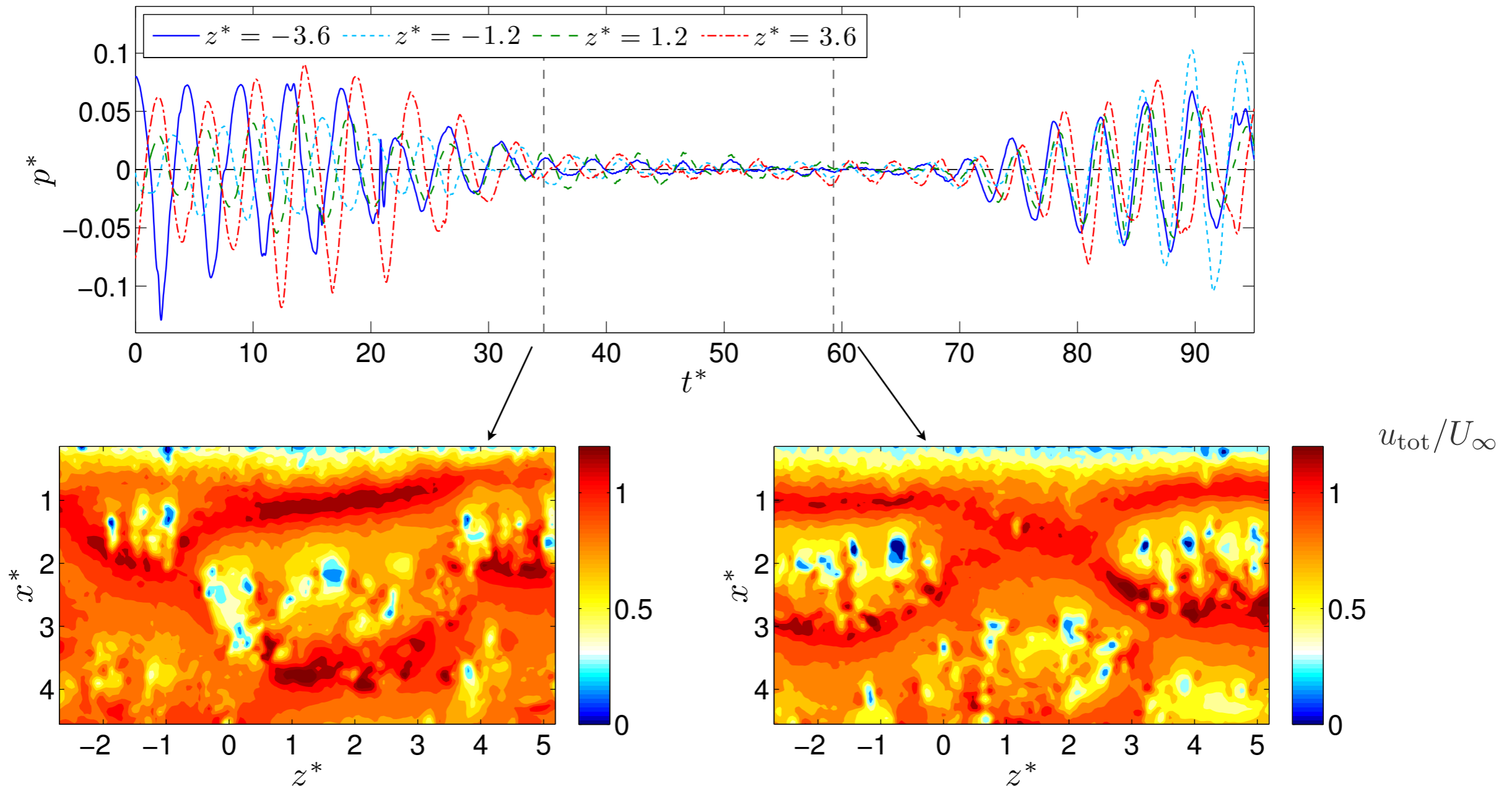
- Correlation of amplitude performed for all combinations of locations in array

$$r_{AA}(i, j, \tau) = \frac{\langle A_p^{(i)}(t) A_p^{(j)}(t + \tau) \rangle}{\sigma_p^{(i)} \sigma_p^{(j)}}$$

- Spanwise coherence of low-frequency modulation increased by forcing

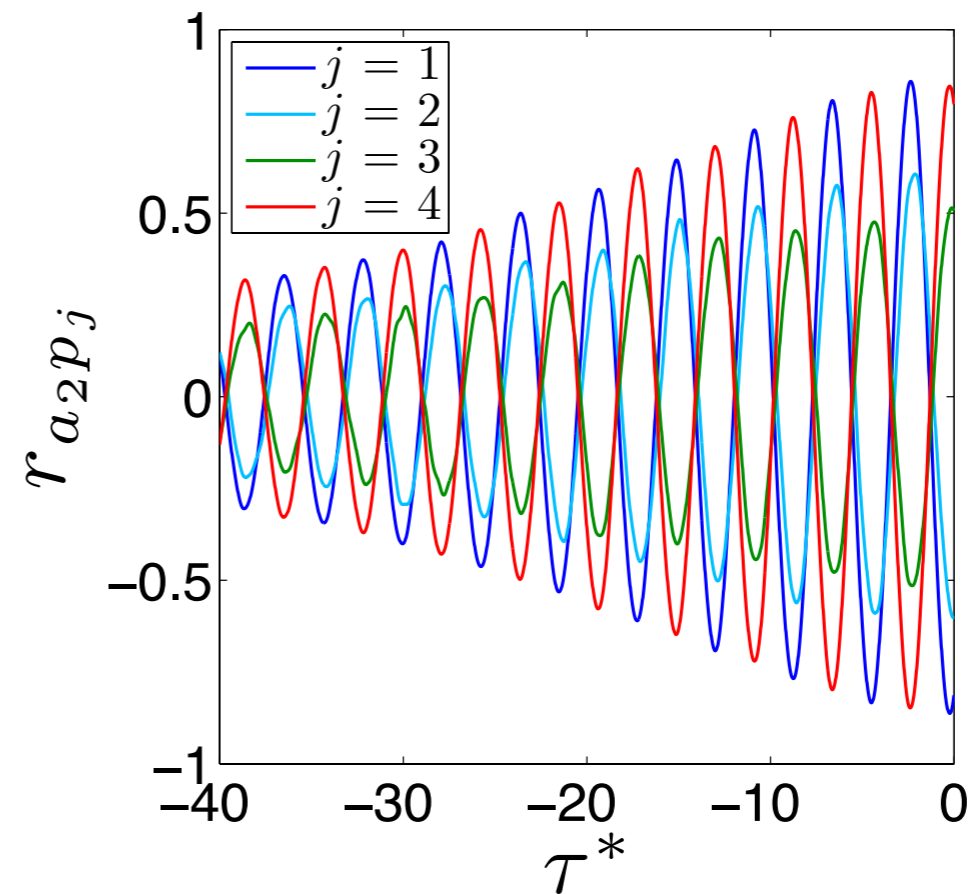
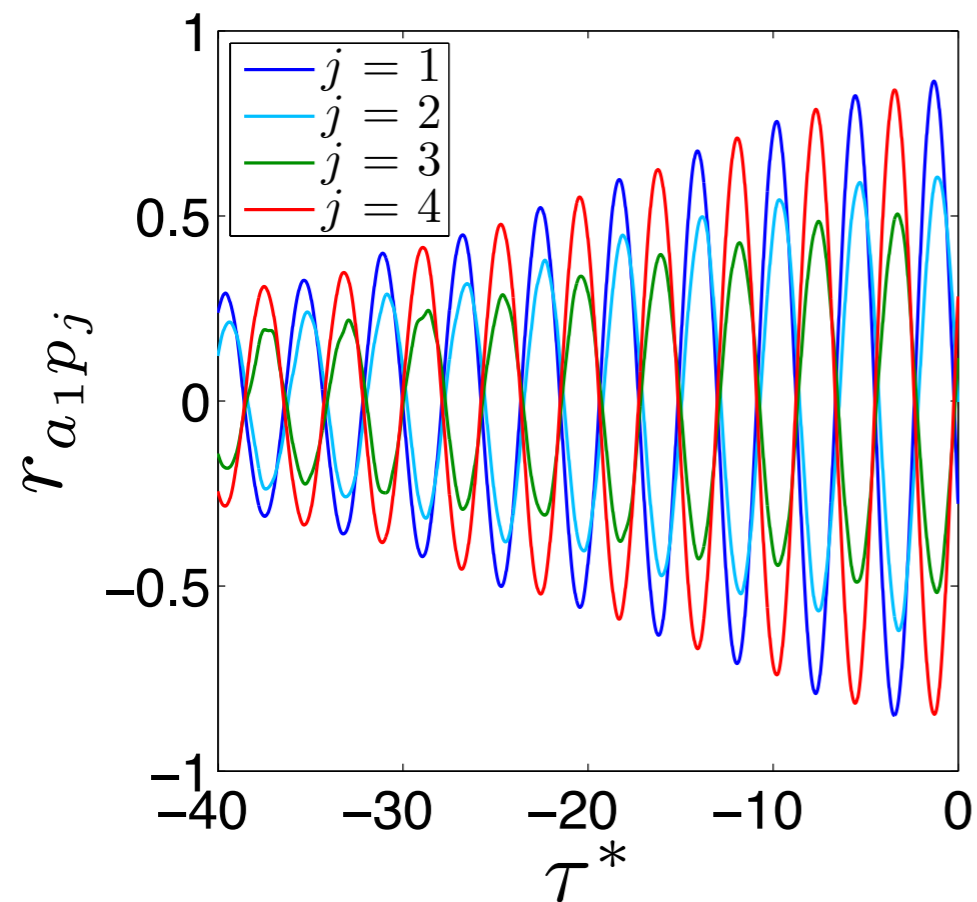


# Modulation and vortex dislocations



- Complex vortex distortions and interactions occur with globally reduced amplitude
- Najjar & Balachandar (1998): switching between regimes in wake of normal flat plate, primary vortices torn apart and secondary vortices disorganized in low drag state

# Stochastic Estimator



$$r_{a_i p_j}(\tau) = \frac{\langle a_i(t) p_j(t + \tau) \rangle}{\sqrt{\langle a_i^2 \rangle \langle p_j^2 \rangle}}$$

$$\tau^* = \tau U_\infty / d$$

- Maximum  $r_{a_i p_j} \simeq 0.85$  for  $j = 1, 4$
- $r_{a_i p_j} < 0.3$  for  $i > 2$  : higher modes not observable with current sensing strategy
- High linear correlation coefficient  $\rightarrow$  linear model: multi-time-delay LSE, used in similar BTE geometry by Durgesh & Naughton (2010) and Tu et al. (2013), curvilinear cavity by Lasagna et al. (2013)



# Model Construction



- Model construction

$$\hat{a}_i(t) = \beta_{ijk} p_j(t + \tau_k), \quad \tau < 0$$

$i = 1, 2$  as motivated by correlation

- Varying parameters in model construction:

$m_1, m_2, \Delta\tau^*$ , and number of sensors used

## Performance Quantification

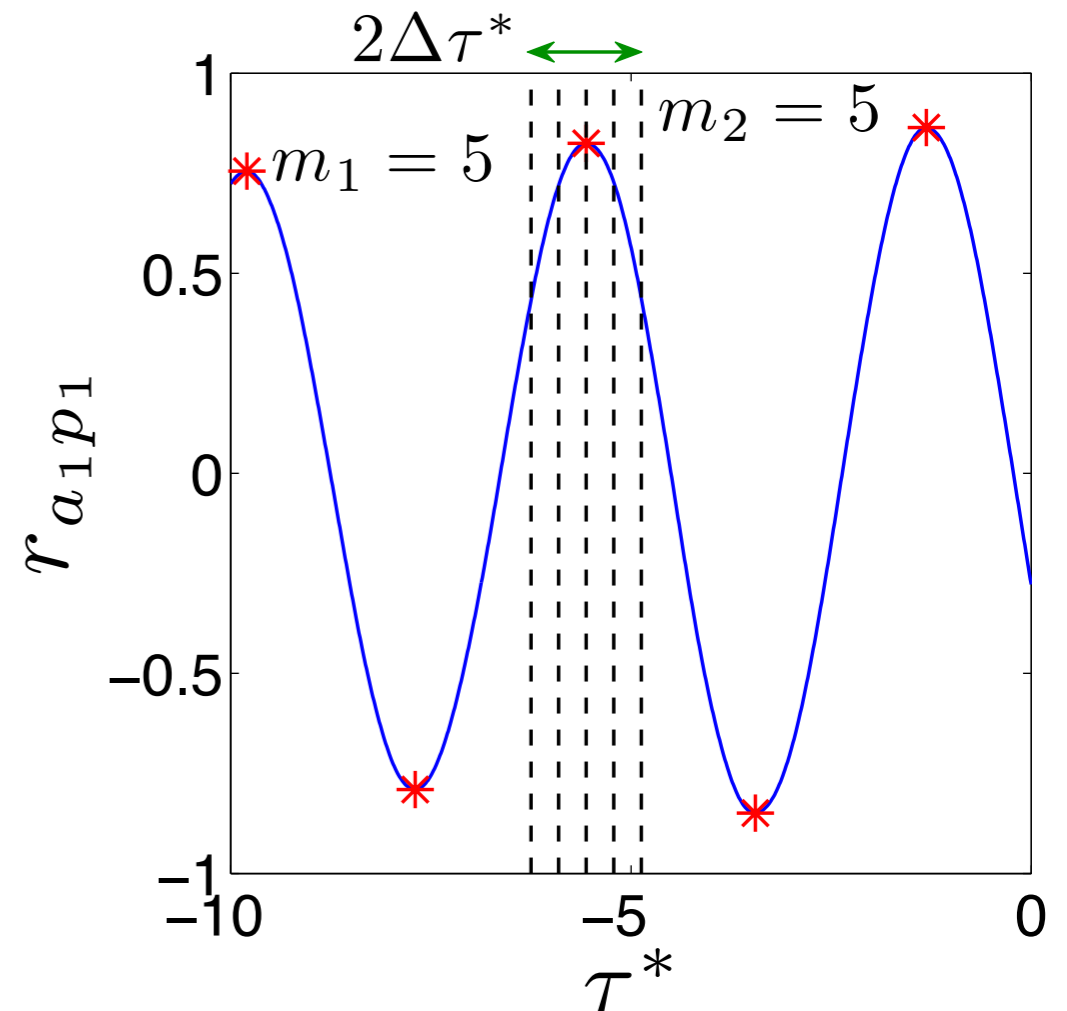
- Fraction of variance residual

$$e_v = \frac{\langle (\hat{a}_1 - a_1)^2 + (\hat{a}_2 - a_2)^2 \rangle}{\langle a_1^2 + a_2^2 \rangle}$$

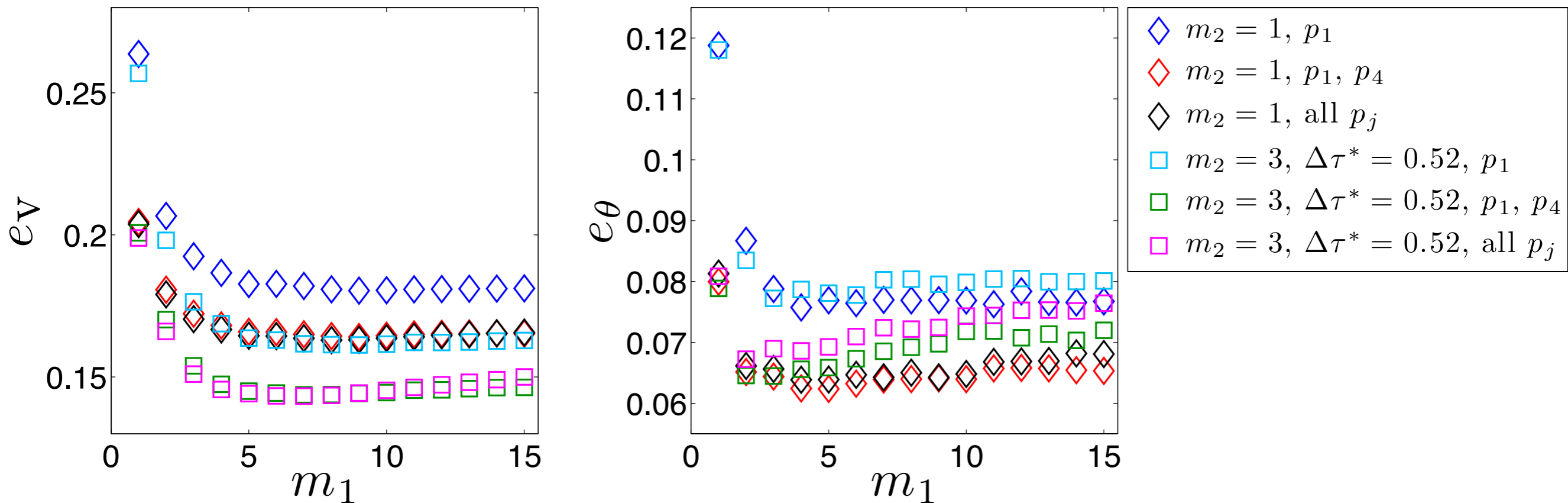
- Phase  $e_\theta = \langle |\hat{\theta} - \theta| \rangle / \pi$ ,  $\theta = \tan^{-1} \left( \frac{a_2 / \sqrt{2\lambda_2}}{a_1 / \sqrt{2\lambda_1}} \right)$ ,  $\theta \in [-\pi, \pi]$

- Variance Inflation Factor (VIF) of regression coefficients quantifies collinearity of predictors:

use average value 
$$\overline{VIF} = \sum_{j=1}^n 1 / (1 - R_j^2)$$



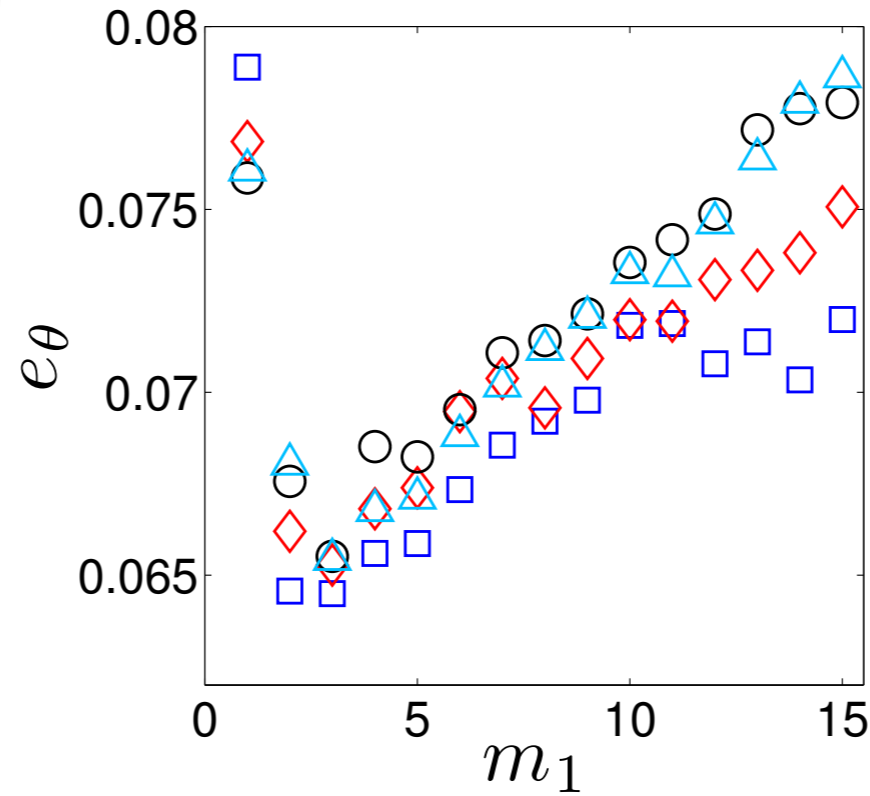
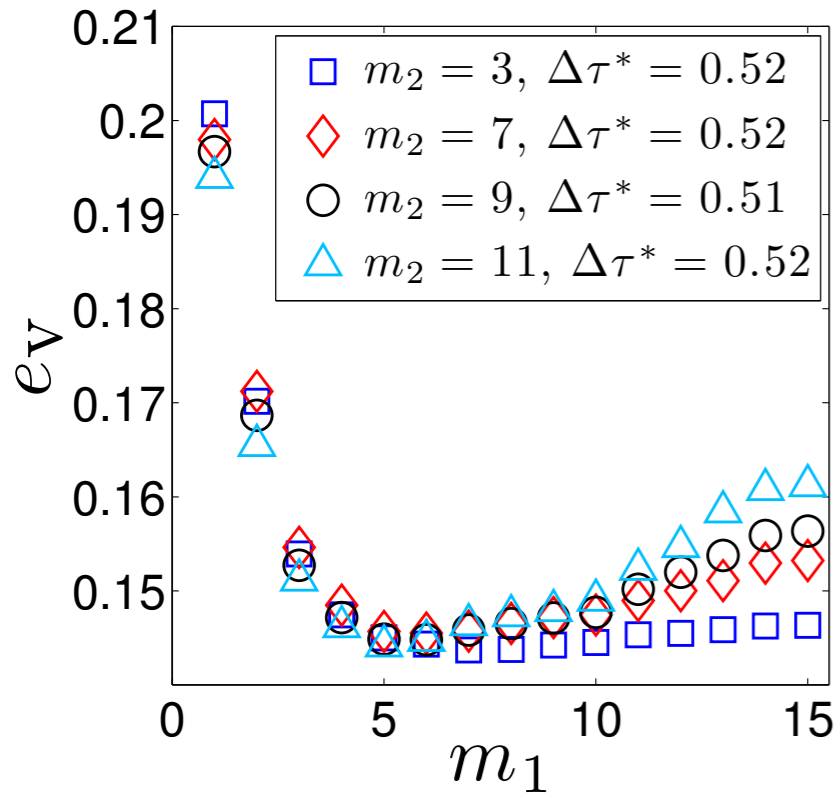
# Correlation and Variance Inflation



- Optimum  $m_1$  around 5
- Significant improvements when using both  $p_1$  and  $p_4$
- $p_2$  and  $p_3$  are not useful

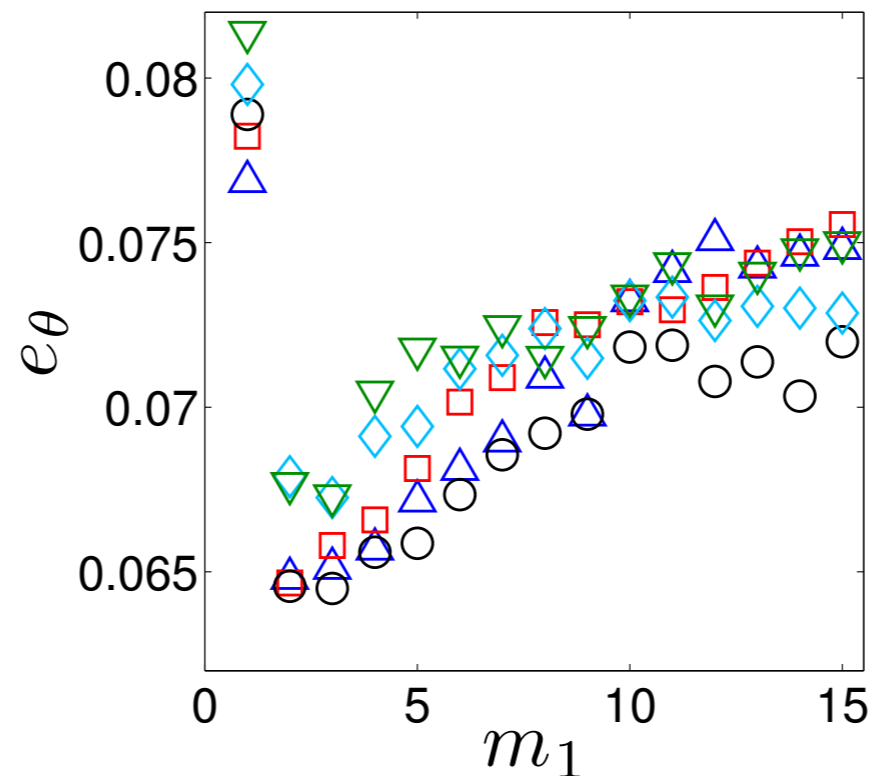
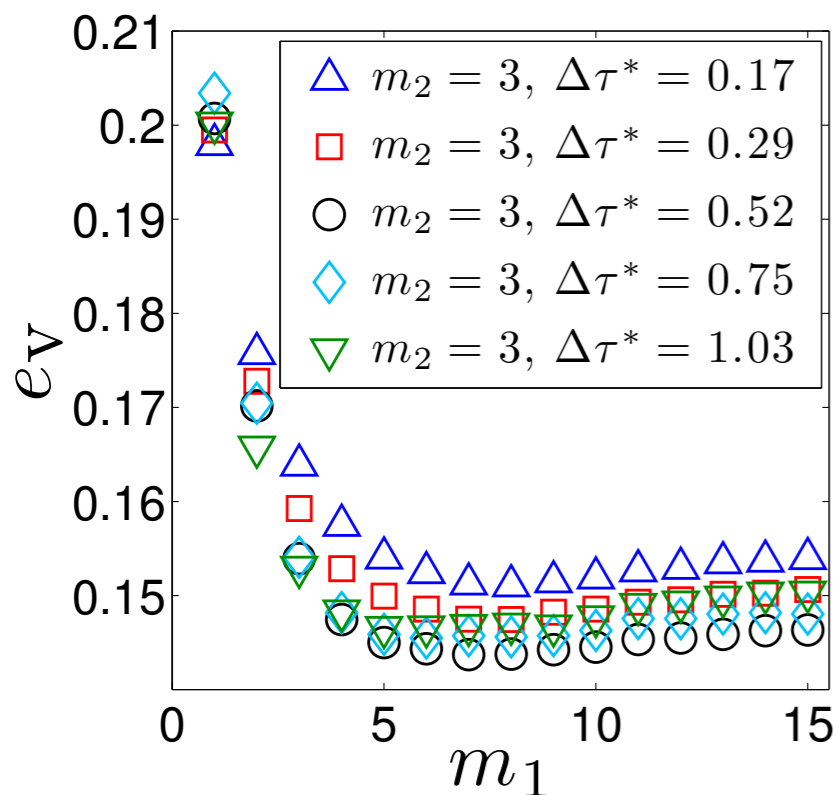


# Correlation and Variance Inflation



Varying  $m_2$ :

$$\overline{VIF} \simeq 12 \text{ to } 10^9$$

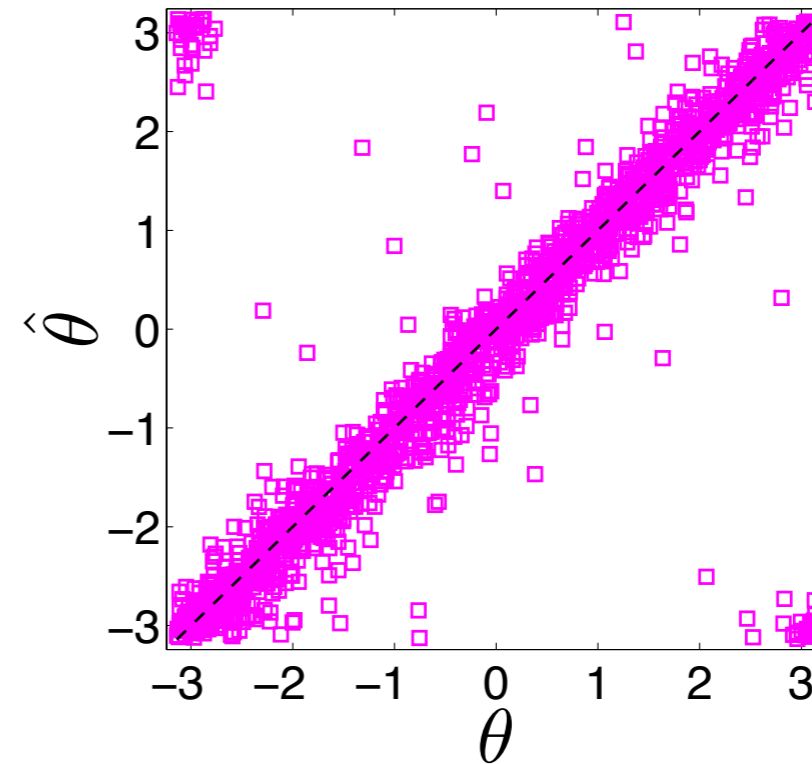
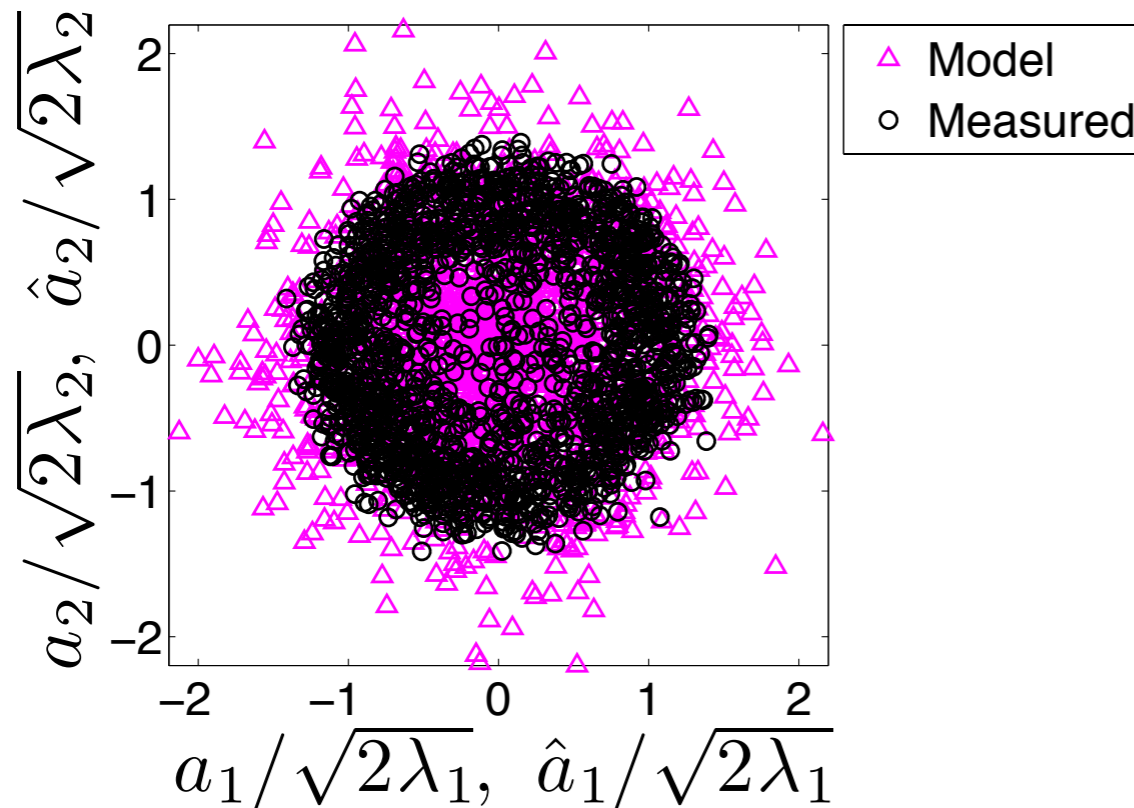


Varying  $\Delta\tau^*$ :

correlation of some predictors decreasing

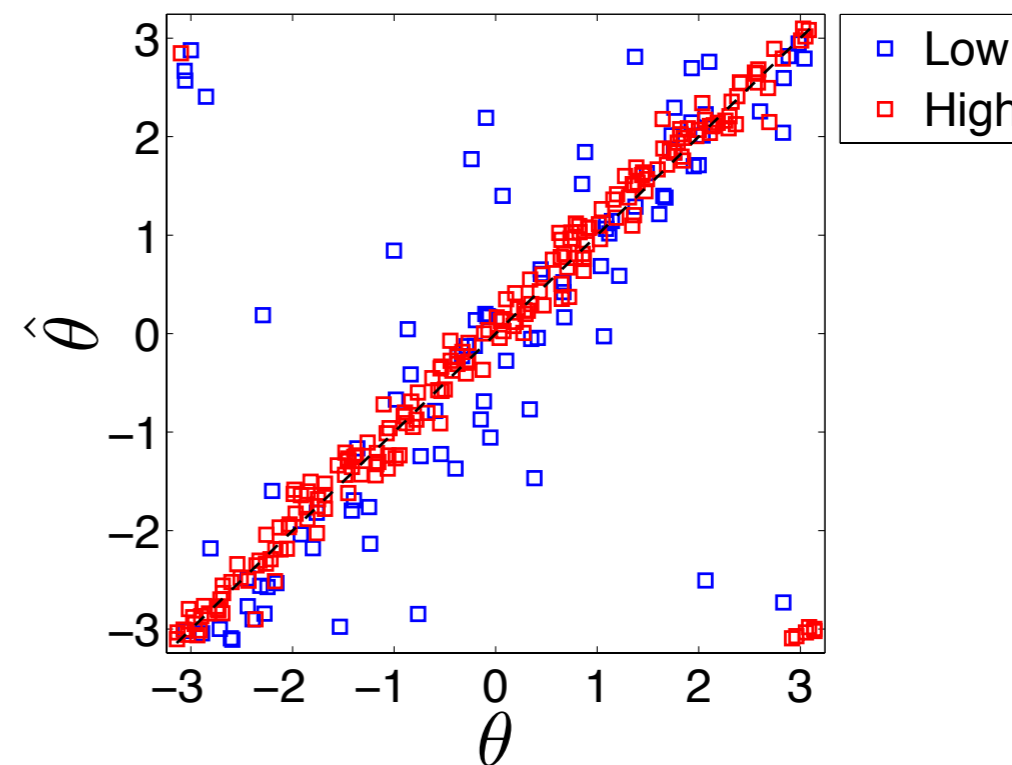
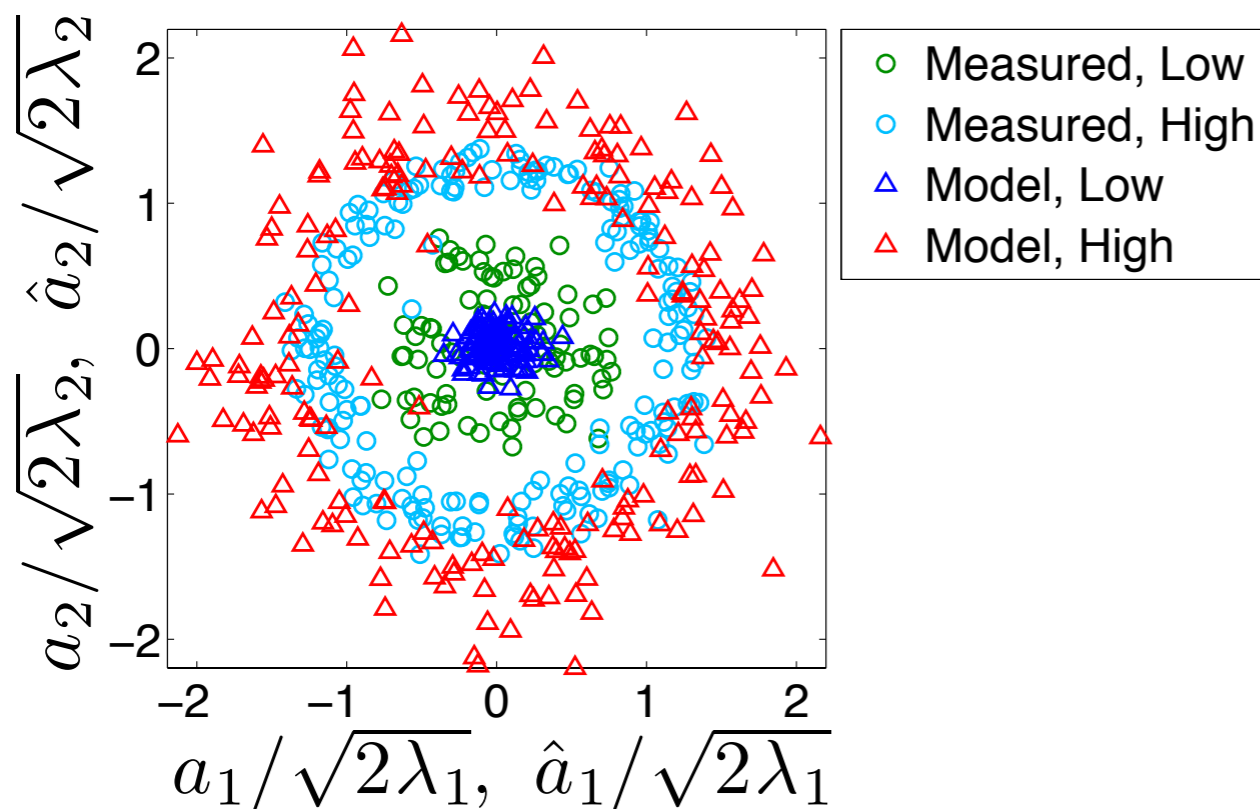


# Validation Set: All Observation



- Model parameters  $m_1 = 5$ ,  $m_2 = 3$ ,  $\Delta\tau^* = 0.52$
- $e_v = 0.145$ ,  $e_\theta = 0.07$
- Linear regression model uses only correlations  
e.g.,  $\langle p_i p_j \rangle$ ,  $\langle a_i p_j \rangle$
- Amplitude modulation of pressure results in model prediction outside of physical variation of  $a_1$ ,  $a_2$

# Validation Set: Categorized Observation

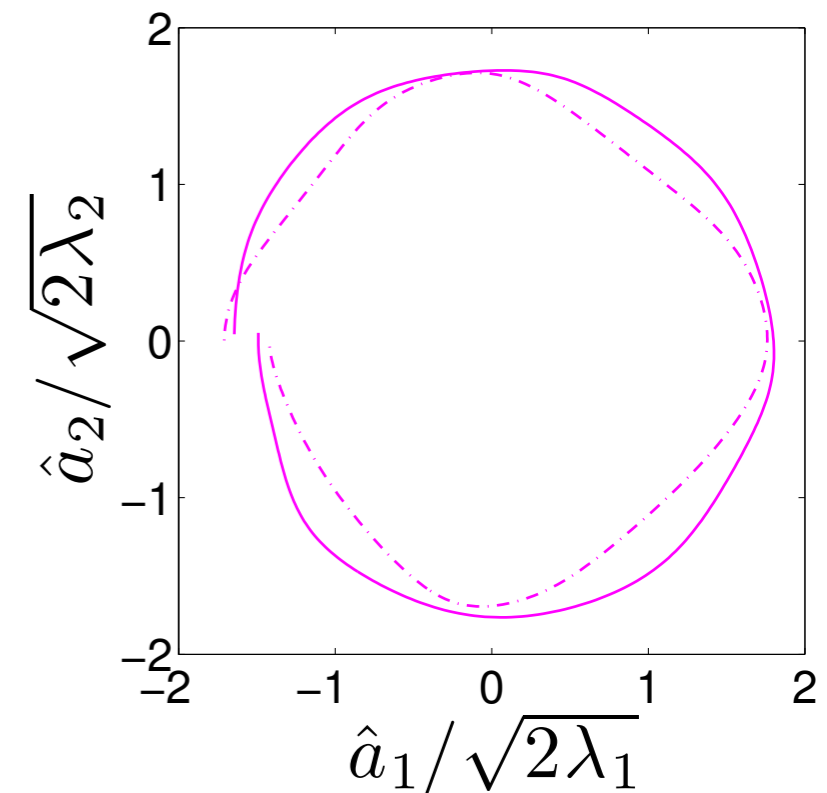
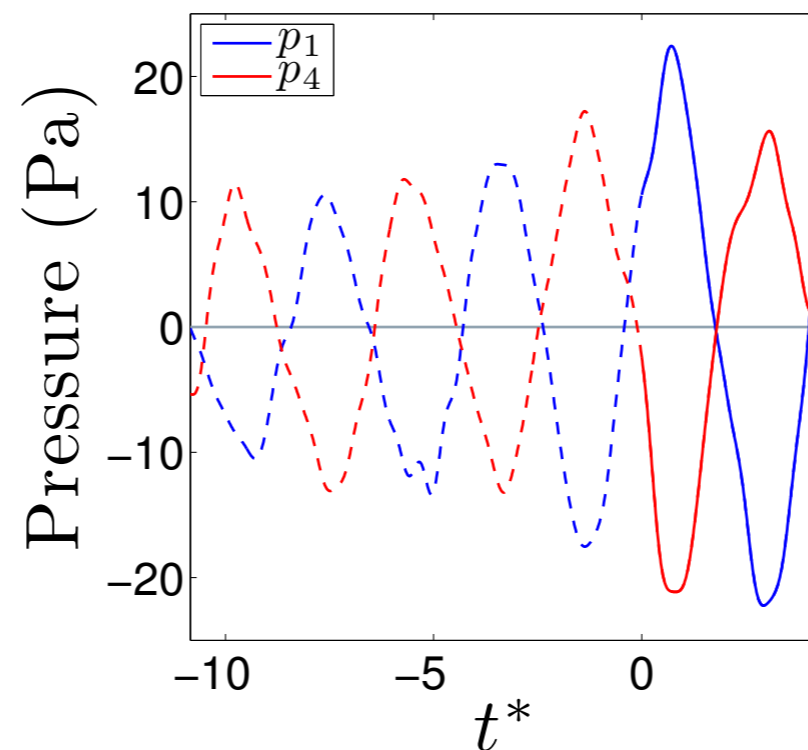
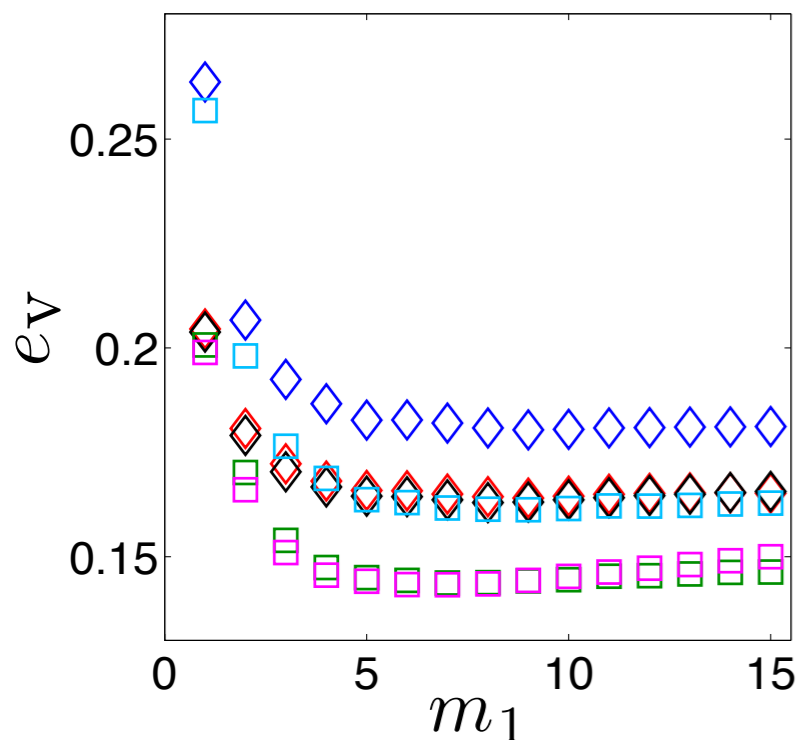


- Separated snapshots at times of low- and high-amplitudes pressure fluctuations
- Lower  $a_1^2 + a_2^2$  related to low pressure, but not linearly
- Low:  $e_v = 0.56$ ,  $e_\theta = 0.18$
- High:  $e_v = 0.15$ ,  $e_\theta = 0.04$
- Degradation of phase relationship typical of low amplitude pressure state

# Time-resolved Estimation



- What physical differences result from parameter changes?
- Model I:  $m_1 = 5, m_2 = 1 \rightarrow e_v = 0.166$
- Model II:  $m_1 = 5, m_2 = 5 \rightarrow e_v = 0.141$



- Lower minimum  $e_v$  associated with more round trajectory in phase-space, believed to match physical evolution more closely
- $m_2 = 1$  (predictors at peak correlation only) shows that pressure fluctuations are not purely sinusoidal. Curvature modified by  $m_2 > 1$

# Effect of Variance Inflation



- High VIF becomes problematic if input to model is noisier than the training data

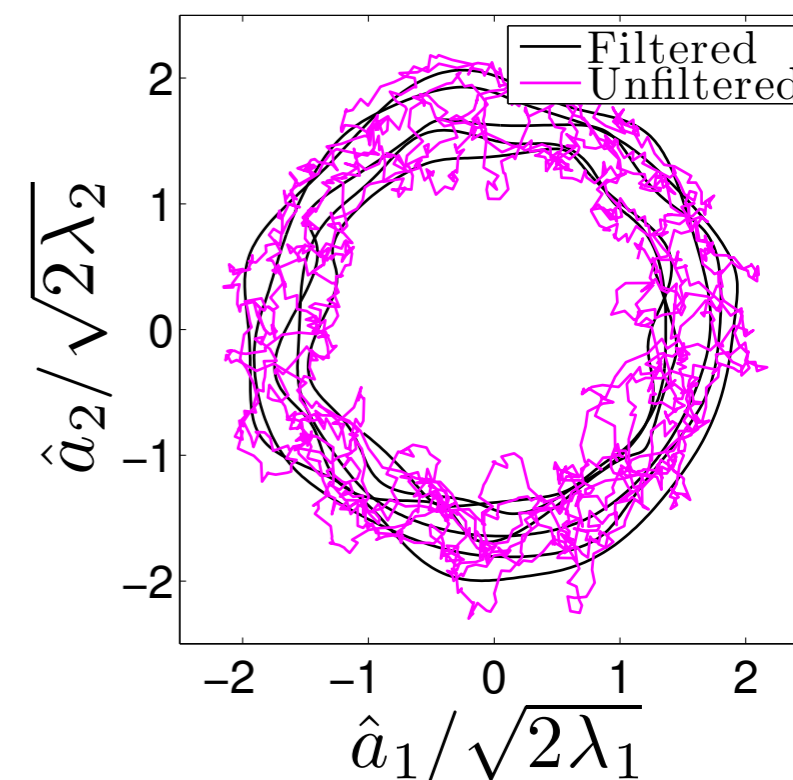
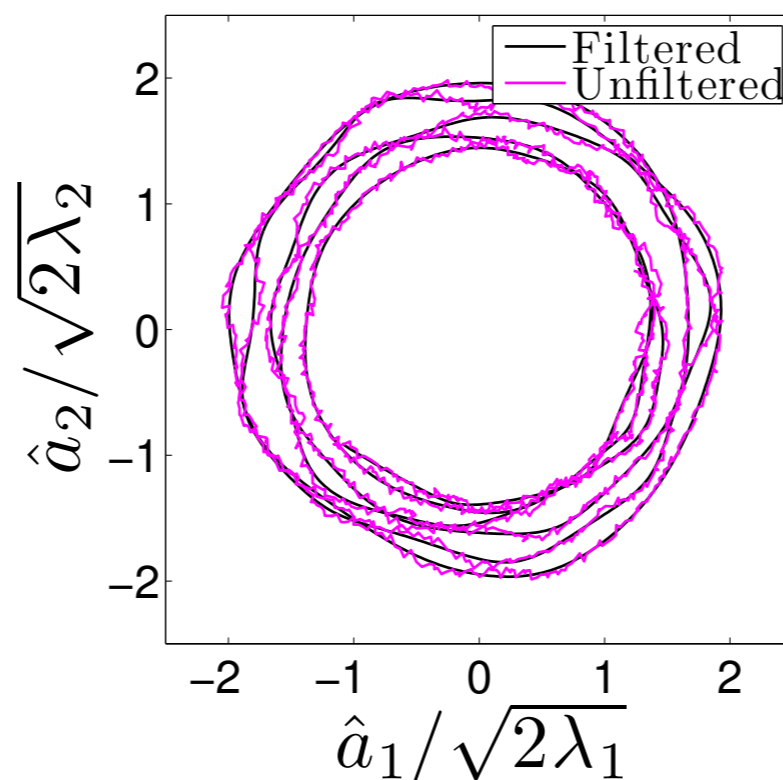
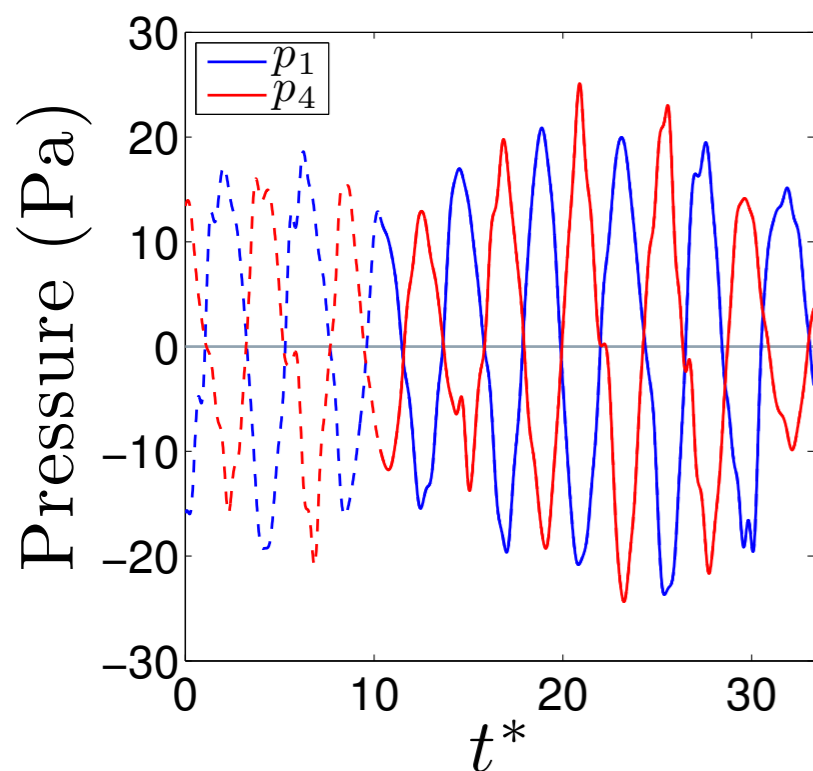
- Model I:

$$m_1 = 5, m_2 = 3, \Delta\tau^* = 0.52 \rightarrow \overline{VIF} \simeq 12, |\hat{\beta}_j| \simeq 1$$

- Model II:

$$m_1 = 5, m_2 = 7, \Delta\tau^* = 0.52 \rightarrow \overline{VIF} \simeq 900, |\hat{\beta}_j| \simeq 10$$

- Estimate from unfiltered pressure signals in test set



# Summary and Perspective



- Active **control of flow separation** can have some significant benefits to improve aircraft performance.
- Commercial aviation: primarily **indirect benefits to be gained**.
- Approach to separation control fundamentally different for streamlined bodies (e.g., airfoil) and bluff bodies (e.g., landing gear).
- **Streamlined bodies**: unsteady forcing at high-frequency is promising
  - use of 3D forcing and/or closed-loop control could further improve performance and efficiency
- **Bluff bodies**: inhibiting movement transfer in the wake reduces unsteadiness and pressure drag
- Beginning to see application in industry - with much more potential to be reached

